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## **Master Degree Thesis**



# EFFECT OF ARGON INJECTION ON FLUID FLOW AND HEAT TRANSFER IN THE CONTINUOUS SLAB CASTING MOLD

BY

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#### THESIS

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#### Abstract

Three-dimensional computational models of turbulent multiphase flow and heat transfer in the molten steel in a continuous casting nozzle and mold region are developed and analyzed using the commercial package CFX. To increase the accuracy of the inlet conditions for the mold simulation, a nozzle simulation is first done and then its results at the port are input as the inlet conditions in the mold simulation. Experiments are conducted on water models to verify both the single phase and multiphase turbulent flow models both qualitatively and quantitatively. The simulations consistently match measurements in both the water model and steel caster, although flow in the 0.4-scale water model is sometimes different than in the steel caster when gas is present. Thermocouple measurements are conducted in the liquid pool of the steel caster to verify the heat transfer model. The computed temperatures match the measurements. Parametric studies are performed with the multiphase flow model to quantify the important effects of gas fraction, slab width, SEN submergence depth and bubble size on the flow pattern. It is shown that there is more single roll tendency with higher gas fraction, wider slab, smaller SEN submergence depth and smaller bubble size. Single roll flow pattern is more likely than double roll to cause defects such as slivers and pencil pipes due to higher predicted level fluctuation at the top surface and larger downward velocity along the solidified shell. Flow pattern transition might also be detrimental. For a given slab width and SEN submergence depth, there is a critical gas fraction where the flow pattern changes from double roll to single roll. The critical gas fraction appears not to change much with the casting speed. For a given casting speed, the critical gas fraction increases with decreasing the slab width, increasing the SEN submergence depth, and decreasing bubble size.

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List of Tablesix		
List of Figuresx		
Nomenclaturexiv		
Section 1. Introduction1		
Section 2. Model Formulation		
2.1 Introduction		
2.2 Single Phase Model		
2.3 Multiphase Flow Model8		
2.3.1 Governing Equation in Multiphase Flow		
2.3.2 Inter-phase Drag Model for Multiphase Flow9		
2.3.3 Buoyancy Term in Multiphase Momentum Equations11		
2.3.4 Source Term in Multiphase Model11		
2.4 Multi-Size-Group Model		
2.4.1 Break-up Model15		
2.4.2 Coalescence Model16		
2.5 Heat Transfer Model		
2.6 Boundary Conditions19		
2.6.1 Inlet Boundary Condition19		
2.6.2 Outlet Boundary Condition20		
2.6.3 Symmetry Boundary Condition21		
2.6.4 Wall Boundary Condition21		

## **Table of Contents**

2.7 Solution Method	23
Section 3. Validation of Fluid Flow Models	25
3.1 Introduction	25
3.2 Comparison of Simulations and PIV Measurements	25
3.2.1 Scaled Water Model Experimental Description	25
3.2.2 Numerical Simulation of 0.4 Scaled Water Model	26
3.2.3 Proportional Scaling and Froude Scaling	27
3.2.4 Comparison of Numerical Simulation Results and PIV	
Measurements on 0.4 Scaled Water Model	29
3.3 Numerical Simulation of Flow in Real Continuous Caster	30
3.4 Conclusions	32
Section 4. Validation of Heat Transfer Model	45
4.1 Introduction	45
4.2 Experiments	46
4.3 Numerical Simulation of Heat Transfer	47
4.4 Modeling and Experimental Results	48
4.5 Conclusions	49
Section 5. Effect of Gas Bubble Size on Fluid Flow in Continuous Casting Mold	64
5.1 Introduction	64
5.2 Model Description	65
5.3 Results and Analysis	67
5.3.1 Effect of Bubble Size on Flow Pattern	67
5.3.2 Effect of Bubble Size on Gas Entrapment	67

5.3.3 Effect of Bubble Size on Inclusion Formation	69
5.3.3.1 Effect of Bubble Size on Level Fluctuations	70
5.3.3.2 Effect of Bubble Size on Downward Velocity	
Near Solidifying Shell	71
5.4 Conclusions	72
Section 6. Effect of Argon Flow Rate on Flow Patterns Switching	
in Continuous Slab Casting	83
6.1 Introduction	83
6.2 Definition of the Mean Bubble Size in This Study	84
6.3 Bubble Size and Its Distribution in the 0.4-Scale Water Model	85
6.4 Validation of MUSIG Model by 0.4-Scale Water Model	86
6.5 Bubble Size and Its Distribution in the Steel Casting Nozzle	87
6.6 Simulation of Flow in the Nozzle	89
6.7 Simulation of Flow in the Continuous Caster	90
6.8Results and Analysis	91
6.8.1 Influence of Casting Speed on Flow in the Continuous Caster	91
6.8.2 Influence of Gas Fraction on the Flow Pattern	92
6.8.3 Influence of Slab Width on the Flow Pattern	93
6.8.4 Influence of Submergence Depth on the Flow Pattern	94
6.8.5 Difference between Flow in Water Model and Steel Caster	94
6.9 Conclusions	95
Section 7. Conclusions	31
7.1 Validation of Fluid Flow Models in Continuous Casting1	31

7.2 V	Validation of Heat Transfer Model	132
7.3 E	Effects of Gas Injection on Fluid Flow in the Steel Caster	132
7.4 E	ffects of Slab Width and Nozzle Submergence Depth on	
th	e Fluid Flow in the Steel Caster	134
Appendix A	A Sample Command File for Heat Transfer Model	136
Appendix B	A Sample Command File for Multiphase MUSIG Model	
	Using Output of Nozzle Simulation as Input Conditions	138
Appendix C	Subroutine for Heat Transfer Model	142
Appendix D	Subroutine for Multiphase Flow with MUSIG Model	147
Appendix E	Subroutine for Nozzle Simulation	157
Appendix F	Subroutine for Multiphase MUSIG Model	
	Using Output of Nozzle Simulation as Input Conditions	170
Reference		184

## List of Tables

Table 3.1	Boundary Conditions in Numerical Simulation
Table 3.2	Parameter Used in Water Model Simulation
Table 3.3	Dimensionless Position of Circulation Eyes and Impingement Points34
Table 4.1	Parameters of Mold and Nozzle50
Table 4.2	Experimental Parameters51
Table 4.3	Parameters Used in Heat Transfer Simulation52
Table 5.1	Parameters Used in the Multiphase Flow Simulation74
Table 6.1	Nozzle Parameters (Steel Simulation)97
Table 6.2	Measured Bubble Size Distribution in LTV 0.4-scale Water Model
	for Case A (23.2 mm/s (55ipm) + 13 SLPM/11% hot gas)98
Table 6.3	Measured Bubble Size Distribution in LTV 0.4-scale Water Model
	for Case B (14.8 mm/s (35ipm) + 6.3 SLPM/8.5% hot gas)98
Table 6.4	Parameters for 0.4-scale Water Model for Case A99
Table 6.5	Parameter for 0.4-scale Water Model for Case B100
Table 6.6	Measured Bubble Size Distribution for Case A in Double-Needle
	Experiments101
Table 6.7	Measured Bubble Size Distribution for Case B in Double-Needle
	Experiments101
Table 6.8	Parameters in the Steel Caster Simulation102

## List of Figures

Figure 1.1	Schematic of tundish, nozzle and mold in continuous casting5
Figure 3.1	Simulation domain for three molds35
Figure 3.2	Convergence history of numerical simulations of 0.4 proportionally
	scaled water model, 0.4 Froude scaled water model, full size water
	model and full size real continuous casting
Figure 3.3	Comparison of computed flow patterns between 0.4 Froude scaled,
	0.4 proportionally scaled and full size water model
Figure 3.4	Comparison of speed along jet and centerline between 0.4 Froude
	scaled, 0.4 proportionally scaled and full size water model
Figure 3.5	Comparison of speed along horizontal lines between 0.4 Froude
	scaled, 0.4 proportionally scaled and full size water model
Figure 3.6	Comparison of flow patterns between
	PIV measurements and simulation of 0.4-scale water model40
Figure 3.7	Comparison of velocities along jet between simulation and
	PIV measurements of 0.4-scale water model and real caster41
Figure 3.8	Comparison of velocities along centerline between simulation and
	PIV measurements of 0.4-scale water model and real caster42
Figure 3.9	Comparison of flow patterns between simulations of continuous
	casting and 0.4 proportionally scaled water model43
Figure 3.10	A schematic of relationship between numerical simulation,
	scaled water model and real caster44
Figure 4.1	Equipment used in measuring temperature in the liquid pool53

Figure 4.2	Temperature slope of three different material layers54
Figure 4.3	Position of measuring points54
Figure 4.4	Residuals of heat transfer modeling55
Figure 4.5	Predicted velocity on center plane between wide faces
Figure 4.6	Predicted temperature on center plane between wide faces57
Figure 4.7	Measurement and modeling at 50 mm from SEN58
Figure 4.8	Measurement and modeling at 125 mm from SEN59
Figure 4.9	Measurement and modeling at midway between SEN and narrow face60
Figure 4.10	Measurement and modeling at 125 mm from narrow face61
Figure 4.11	Measurement and modeling at 50 mm from narrow face62
Figure 4.12	Heat flux along three typical lines in the caster63
Figure 5.1	Schematic of simulation domain
	with boundary conditions and typical meshes75
Figure 5.2	Effects of bubble size on fluid flow pattern in continuous casting mold.76
Figure 5.3	Effect of bubble size on gas penetration depth78
Figure 5.4	Gas fraction contours for different sized bubbles78
Figure 5.5	Gas percentage contours for different bubbles79
Figure 5.6	Overlap of the contours of gas fraction and vertical velocity80
Figure 5.7	Schematic of the first mechanism of flux entrapment81
Figure 5.8	Schematic of the second mechanism of flux entrapment81
Figure 5.9	Effect of bubble size on the maximum kinetic energy and
	level fluctuation on the top surface82
Figure 5.10	Effect of bubble size on vertical velocity around top surface perimeter82

Figure 6.1	Bubble size in the water model (23.2 mm/s + 13 SLPM/ 11% hot gas)103
Figure 6.2	Bubble size in the water model
	(14.8 mm/s + 6.3 SLPM/ 8.5% hot gas)104
Figure 6.3	Bubble size distribution in the mold
	(measurements in 0.4-scale water model)105
Figure 6.4	Comparison of velocity at centerplane between PIV measurements,
	simulation and eyeview for case A106
Figure 6.5	Comparison of simulation and eyeviews while adjusting liquid level
	with 15% increase in liquid flow rate for case A107
Figure 6.6	Comparison of velocity at centerplan between PIV measurements,
	simulation and eyeview for case B108
Figure 6.7	Comparison of simulation and eyeviews while adjusting liquid level
	with 15% increase in liquid flow rate for case B109
Figure 6.8	Schematic of water model of bubble formation110
Figure 6.9	Assumed bubble size distribution in the steel caster nozzle
	(based on double-needle experimental measurements)111
Figure 6.10	Average diameter of bubbles vs gas flow rate112
Figure 6.11	Geometry of the nozzle113
Figure 6.12	Liquid velocity in the nozzle114
Figure 6.13	Liquid velocity in the nozzle115
Figure 6.14	Correspondence of nozzle simulation output to mold simulation input116
Figure 6.15	Correspondence of nozzle simulation output to mold simulation input117
Figure 6.16	Simulation domain for mold with typical meshes118

Figure 6.17	Computed steel flow pattern with distributed bubble size
	(Wide face, 1.854 m slab, 23.2 mm/s, 11% gas)119
Figure 6.18	Computed steel flow pattern with distributed bubble size
	(Wide face, 1.854 m slab, 14.8 mm/s, 8.5% gas)120
Figure 6.19	Computed steel flow pattern with distributed bubble size
	(Wide face, 1.854 m slab, 14.8 mm/s, 16.4% gas)121
Figure 6.20	Computed steel flow pattern with distributed bubble size
	(Narrow face, 1.854 m slab, 23.2 mm/s, 11% gas)122
Figure 6.21	Computed steel flow pattern with distributed bubble size
	(Narrow face, 1.854 m slab, 14.8 mm/s, 8.5% gas)123
Figure 6.22	Computed steel flow pattern with distributed bubble size
	(Narrow face, 1.854 m slab, 14.8 mm/s, 16.4% gas)124
Figure 6.23	Flow pattern identification for 0.4-scale water model125
Figure 6.24	Flow pattern identification in steel caster
	(from simulation of real caster)126
Figure 6.25	Flow pattern identification (estimated from simulation of real caster)127
Figure 6.26	Relationship between throughput and submergence depth
	(1.854 m slab)128
Figure 6.27	Velocity vector at centerplane for different nozzle submergence depth
	(1.854 m slab, 23.2 mm/s casting speed and 11% hot argon)129
Figure 6.28	Velocity vector at centerplane for different nozzle submergence depth
	(1.854 m slab, 14.8 mm/s casting speed and 8.5% hot argon)130

## Nomenclature

А	Projected area of bubble
В	body force
B <sub>b</sub>	birth rate due to break-up
B <sub>c</sub>	birth rate due to coalescence
$C_{ii}^{(drag)}$	inter-phase drag terms
$C_{\rm D}$	drag coefficient
C <sub>p</sub>	specific heat
d <sub>i</sub>	Diameter of bubble i
d <sub>s</sub>	Sauter mean diameter of bubble
$\mathrm{D}_{\mathrm{b}}$	Death rate of bubbles in group i
Е	log-layer constant for velocity
$\mathrm{E}_{\mathrm{H}}$	log-layer constant for enthalpy
$f_{\scriptscriptstyle BV}$	breakage volume fraction
$f_{liq}$	volume fraction of liquid phase
$f_{gas}$	volume fraction of gas phase
$g(v_j:v_I)$	Break-up rate of bubbles of volume $v_i$ into bubbles of volume $v_i$
h	static enthalpy
Н	total enthalpy
К	turbulent kinetic energy
$k_{\scriptscriptstyle e\!f\!f}$	effective conductivity of molten steel
$k_o$	molecular conductivity
$k_t$	turbulent conductivity
p	pressure
$Pr_t$	turbulent Prandtl number for velocity
Pr <sub>H</sub>	turbulent Prandtl number for enthalpy
$\mathbf{Q}_{\mathrm{ij}}$	coalescence rate of bubbles between size group i and j
Re <sub>b</sub>	bubble Reynolds number
$\mathbf{S}_{ij}$	collision cross-sectional area of bubbles
$\mathbf{S}_{\text{masss}}, \mathbf{S}_{\text{mom}}$	source term in continuity and momentum equations
Т	temperature

$T_{liq}$	liquidus temperature of steel
V	fluid velocity
$\mathbf{V}_{liq}$	velocity of liquid phase
$\mathbf{V}_{gas}$	velocity of gas phase
x	distance in model domain from centerplane to narrow face
у	distance in model domain from centerplane to wide face
$y^+$	non-dimensional distance from wall
$y_o^+$	crossover point for velocity
$y_H^+$	crossover point for enthalpy
Z.	distance in model domain from top surface to mold exit
Φ	shear production in turbulence models
α	inlet jet angle
8	turbulent energy dissipation rate
κ	Von-Karmen constant
$\mu_{\scriptscriptstyle e\!f\!f}$	effective viscosity
$\mu_o$	molecular viscosity
$\mu_t$	turbulent viscosity
ρ	fluid density
$\theta^{\mathrm{T}}$	collision frequency due to turbulence mechanism
$\theta^{\scriptscriptstyle B}$	collision frequency due to buoyancy mechanism
$\theta^{LS}$	collision frequency due to laminar shear mechanism
$ au_{ij}$	contact time of bubble I and bubble j
η	collision efficiency
Xc	critical dimensionless energy for break-up
ξ	dimensionless size of eddies
λ	turbulence eddy size
$\sigma_{\rm K}$	constant in turbulence models
$\sigma_{\epsilon}$	constant in turbulence models

### Section 1

### Introduction

Most steel in the world is solidified by continuous casting. Although novel casting processes are being invented, continuous casting will remain as the major process for mass production of steel in the foreseeable future. A schematic of the continuous casting process is shown in Figure 1.1 [1]. Liquid steel is poured into a tundish from ladles. From the tundish, liquid steel flows through a ceramic submerged entry nozzle (SEN), where a slide gate controls the flow rate of the steel, then exits through bifurcated ports into the liquid pool in the mold. To prevent clogging in the nozzle, argon gas is often injected into the nozzle through ceramic pores in the nozzle wall. Steel in the mold solidifies against the water-cooled mold wall and forms a shell. The shell holds the liquid steel pool. Out of the mold, rollers roll on the shell to drag it downward at "casting speed". The shell grows as the shell goes down and eventually solidifies through all thickness to form a solid slab.

The flow pattern of liquid steel in the liquid pool has a great influence on quality of the continuous cast steel slab. The flow pattern may affect solidified shell thickness by influencing heat transfer behavior in the mold. High temperature liquid steel flows into the mold from the SEN port with high momentum and impinges the narrow face. The high temperature flow has thinning effect on the shell. If the impingement area is too narrow, heat flow at impingement point will be high and the thinning of the shell will be large. Under extreme conditions, the thinner shell will not be strong enough to hold the liquid steel and a costly "breakout" happens [2]. Flow in the mold should be controlled to deliver heat against the narrow face relatively spreadly to reduce the shell thinning effect. The flow pattern may affect liquid flux filling into the gap between the steel shell and the mold wall and then produce slab surface quality problems such as cracks and high vibration marks [3]. For a single roll flow pattern, the flow hits the top surface with high momentum. Vertical momentum of the flow will lift the level of the top surface, creating a variation in the interface level or "standing wave". A high standing wave at the mold wall will prevent liquid flux filling the gap between the steel shell and the mold wall. Reducing the flux in the gap will increase the friction force between the mold wall and the solidified shell and cause the non-uniform heat flow which increases thermal stress in the shell. The friction force and the thermal stress together may cause cracks in the shell. The flow pattern also affects the trajectory of inclusions and argon bubbles and cause quality problems such as slivers and pencil-pipes. Small gas bubbles and inclusion particles can stay in the flow for a much long time. A bad flow pattern may cause these small inclusions and bubbles to stay in the liquid pool much longer and increase their chance to be trapped by the growing shell [4].

The flow pattern in the mold is associated with many factors, such as flow behavior at the SEN port, SEN port angle, SEN port geometry, submergence depth of the SEN, gas fraction and bubble size in the liquid flow, casting speed and geometry of the mold. The flow pattern can be changed by adjusting above parameters. Previous work studied the effect of gas flow rate, bubble size, SEN submergence depth, casting speed and mold width on the flow pattern [4-8, 17, 22, 29-31]. Although these studies produced positive results, they are also limited in following respects. First, testing data set is often small Most cases have only 3-4 testing cases in [5] and [7]. The small testing set can not give detailed information to support the conclusions. Second, for the bubble size effect, only single sized bubbles are considered [6, 17, 29]. In fact, gas in the flow is multiple sized and bubbles interact with each other. Finally, validation of the flow pattern details with measurements is rare. Even where validation was done, it is not quantitative. The blue ink was used to qualitatively validate the flow pattern in the mold [4].

In this study, both single phase and multiphase flow and single phase heat transfer in the mold are developed, validated quantitatively and applied to investigate the flow pattern in the mold. In section 2, various models are developed to describe the single phase flow, the heat transfer in the single phase flow, the multi-phase flow and the multiple-sized bubbles in the multi-phase flow in the continuous casting mold. In section 3, Particle Image Velocimetry (PIV) is used to quantitatively analyze the flow field in a 0.4 scaled water model. The experimental flow field is compared with the numerically simulated flow field to verify the single phase flow model proposed in section 2. In section 4, temperature in the mold is measured by thermocouples. The single phase heat transfer in the mold is modeled by the heat transfer model developed in section 2. The experimental results are compared with the modeling results. In section 5, a 2-phase flow with uniform bubble size is simulated. A total of 8 different bubble sizes distributed in bubble size range, from 0.5 mm to 5 mm, under normal casting conditions [8] is considered. This section gives a clear understanding on how the flow pattern changes with increasing bubble size. In section 6, a 2-phase flow with multiple bubble sizes is simulated by considering possible coalescence and breakup of bubbles existing in the flow. This section studies how casting conditions influence bubble size distribution in the flow and the flow patterns in the mold.

This study is mainly concerned with flow behavior in the mold and how gas bubbles affect the flow patterns in the mold. It aims to provide perspective to quality problems by analyzing the influence of bubbles on the flow patterns. The computational domain is limited to the mold region of the caster. However, other parts of the continuous casting process, such as the SEN, also influence the flow jet at the port and bubble size. The flow behavior in the SEN has also been studied by Continuous Casting Consortium(CCC) researchers [16]. To address other quality problems, such as cracks, heat conduction and stress in the solidified shell and water-cooled mold wall structure are investigated by CCC researchers [19,20]. All these researches make up a big project to solve the quality problems in the continuous casting. This study is part of this big project.



Figure 1.1 Schematic of tundish, nozzle and mold in continuous casting process [1]

### Section 2

### **Model Formulation**

#### **2.1 Introduction**

In this section, various models will be developed to simulate different fluid flow and heat transfer phenomena in the continuous casting process, including liquid steel (single phase ) problems and argon-liquid steel (multi-phase) problems. Each model selects and solves specific governing equations to give an accurate description of the important physics of the specific problems. All these modeling methods and their governing equations are described. The following models are discussed in this section: single phase turbulent flow model, multiphase turbulent flow model, Multi-Size-Group (MUSIG) model and single phase heat transfer model. Boundary conditions in the modeling of continuous casting are also discussed.

Models of single phase turbulent flow, multi-phase turbulent flow with singlesized bubbles and single phase heat transfer are implemented using the commercial package CFX4.2 [9]. The Multi-Size-Group (MUSIG) model is implemented using CFX4.3 [10].

#### **2.2 Single Phase Model**

Single phase flow in the water mold and continuous casting mold is turbulent flow [5]. The standard K- $\epsilon$  model is used to simulate the turbulent flow with no sources. The continuity equation for non-source single phase standard K- $\epsilon$  model is:

$$\nabla \cdot V = 0 \tag{2.1}$$

The momentum equation for non-source single phase standard K- $\varepsilon$  model is:

$$\nabla \cdot (\rho V \otimes V) = \nabla \left( \mu_{eff} \left( \nabla V + (\nabla V)^T \right) \right) - \nabla p$$
(2.2)

where  $\rho$  is the fluid density, V is the mean velocity, p is the pressure, The effective viscosity,  $\mu_{eff}$  is defined as

$$\mu_{eff} = \mu_o + \mu_t \tag{2.3}$$

where  $\mu_o$  is the molecular viscosity whose value is given in following sections,  $\mu_t$  is the turbulent viscosity. Its value is calculated from the *K* and  $\varepsilon$  turbulence parameters:

$$\mu_{t} = C_{\mu}\rho \frac{K^{2}}{\varepsilon}$$
(2.4)

The turbulent viscosity depends on the turbulent kinetic energy, K, and its rate of dissipation,  $\varepsilon$ . K and  $\varepsilon$  are found by solving two transport equations:

$$\nabla \cdot \left(\rho K \nabla\right) = \nabla \cdot \left( \left( \mu_0 + \frac{\mu_t}{\sigma_k} \right) \nabla K \right) + \Phi - \rho \varepsilon$$
(2.5)

$$\nabla \cdot \left(\rho \varepsilon \mathbf{V}\right) = \nabla \cdot \left( \left( \mu_0 + \frac{\mu_t}{\sigma_{\varepsilon}} \right) \nabla \varepsilon \right) + C_1 \frac{\varepsilon}{K} \Phi + C_2 \rho \frac{\varepsilon^2}{K}$$
(2.6)

where the shear production,  $\Phi$ , is defined by

$$\Phi = \mu_{eff} \nabla V \cdot \left( \nabla V + \left( \nabla V \right)^T \right) - \frac{2}{3} \nabla \cdot V \left( \mu_{eff} \nabla \cdot V + \rho K \right)$$
(2.7)

The empirical constants  $C_1$ ,  $C_2$ ,  $C_\mu$ ,  $\sigma_k$  and  $\sigma_\epsilon$  are given by [9]:

$$C_1=1.44, C_2=1.92, C_{\mu}=0.09, \sigma_k=1.00, \sigma_{\epsilon}=1.30$$

#### **2.3 Multi-Phase Flow Model**

#### 2.3.1 Governing Equations in Multi-Phase Flow

In most continuous casting process, argon is injected into nozzle to prevent clogging of nozzle. Gas and liquid steel form a 2-phase flow. This 2-phase flow is also turbulent flow. If considering gas bubbles with different size as different phases, steel-argon system could be dealt as multi-phase flow. The standard K- $\varepsilon$  model is also used for multi-phase turbulent flow. In the multi-phase model, velocity is solved for each phase separately. K and  $\varepsilon$  are solved only for the liquid phase. The pressure field is shared for the all phases. The continuity and momentum equations for each phase are different from Eqn.2.1 and Eqn.2.2 in three aspects. First, velocity of each phase is not independent from each other. The velocity fields are coupled by an empirical inter-phase drag model to handle the transfer of momentum between the phases (section 2.3.2). Second, the volume fraction of each phase must be introduced into the governing equations. Third, the buoyancy (section 2.3.3) must be considered.

The continuity equation for the gas phase is:

$$\nabla(f_{gas}V_{gas}) = 0 \tag{2.8}$$

The continuity equation for the liquid phase is:

$$\nabla(f_{lia}V_{lia}) = 0 \tag{2.9}$$

The momentum equation for the liquid phase is:

$$\nabla \cdot (f_{liq} \rho_{liq}(V_{liq} \otimes V_{liq})) = \nabla (f_{liq} \mu_{eff} (\nabla V_{liq} + (\nabla V_{liq})^T)) - f_{liq} \nabla p + c^{(drag)}_{liq-gas} (V_{gas} - V_{liq}) + B_{liq}$$

$$(2.10)$$

The momentum equation for the gas phase is:

$$\nabla \cdot (f_{gas} \rho_{gas}(V_{gas} \otimes V_{gas})) = \nabla (f_{gas} \mu_{eff} (\nabla V_{gas} + (\nabla V_{gas})^T)) - f_{gas} \nabla p + c_{gas-liq}^{(drag)} (V_{liq} - V_{gas}) + B_{gas}$$
(2.11)

where *f* is volume fraction, *V* is velocity. *p* is pressure, subscript *liq* and *gas* represent liquid and gas respectively. Inter-phase drag terms  $c^{(drag)}$  are solved in section 2.3.2. B is the buoyancy.

Because relatively smaller density of gas phase, turbulence in gas phase is neglected. Gas phase is set as laminar flow which is solved by Navier-Stokes equations. Turbulent parameters are solved only for liquid phase:

$$\nabla \cdot (f_{liq}(\rho_{liq}K_{liq}V_{liq})) = \nabla \cdot (f_{liq}(\mu_{0liq} + \frac{\mu_{lliq}}{\sigma_{K}})\nabla K_{liq}) + f_{liq}\Phi_{liq} - f_{liq}\rho_{liq}\varepsilon_{liq}$$
(2.12)

$$\nabla \cdot (f_{liq}(\mathbf{p}_{liq} \mathbf{\epsilon}_{liq} V_{liq})) = \nabla \cdot (f_{liq}(\mathbf{\mu}_{0liq} + \frac{\mathbf{\mu}_{tliq}}{\sigma_{\epsilon}}) \nabla \mathbf{\epsilon}_{liq}) + C_1 f_{liq} \frac{\mathbf{\epsilon}_{liq}}{K_{liq}} \Phi_{liq} - C_2 f_{liq} \mathbf{p}_{liq} \frac{\mathbf{\epsilon}^2_{liq}}{K_{liq}}$$
(2.13)

where  $\Phi_{liq}$  is defined as

$$\Phi_{liq} = \mu_{liq_{eff}} \nabla V_{liq} \cdot (\nabla V_{liq} + (\nabla V_{liq})^T) - \frac{2}{3} \nabla \cdot V_{liq} (\mu_{liq_{eff}} \nabla V_{liq} + \rho_{liq} K_{liq})$$
(2.14)

 $\mu_t$  is calculated from Eqn.2.4. The values for the five constants  $C_1$ ,  $C_2$ ,  $C_{\mu}$ ,  $\sigma_k$  and  $\sigma_{\epsilon}$  are the same as those given in section 2.2.

#### 2.3.2 Inter-Phase Drag Model for Multi-Phase Flow

In argon-steel multi-phase system, argon bubbles disperse in continuous liquid steel flow. The drag exerted on an argon bubble by moving liquid steel is assumed to arise from two mechanisms only. The first is due to the viscous surface shear stress, and is called skin friction. The second is due to the pressure distribution around the bubble, and is called form drag. The total drag force can be expressed in terms of the nondimensional drag coefficient  $C_{D_{.}}$  Using drag coefficient  $C_{D_{.}}$  the inter-phase drag terms  $c_{liq\ gas}^{(drag)}$  and  $c_{gas\ liq}^{(drag)}$  in the momentum equations of above multi-phase model are defined as:

$$c_{liq\ gas}^{(drag)} = \frac{3}{4} \frac{C_D}{d} f_{gas} \rho_{liq} |V_{gas} - V_{liq}|$$
(2.15)

$$c_{gas\ liq}^{(drag)} = \frac{3}{4} \frac{C_D}{d} f_{liq} \rho_{gas} |V_{liq} - V_{gas}|$$
(2.16)

The drag coefficient  $C_D$  is defined as:

$$C_D = \frac{D}{\frac{1}{2}\rho V^2 A} \tag{2.17}$$

where D is the magnitude of the drag force, A is the projected area of the bubble in the direction of flow. For a particle of a given shape, undergoing motion in a Newtonian incompressible fluid,  $C_D$  depends only on Reynolds number:

$$\operatorname{Re}_{b} = \frac{\rho_{liq} \left| V_{liq} - V_{gas} \right| d}{\mu_{o_{liq}}}$$
(2.18)

The function  $C_D(\text{Re}_b)$  is determined experimentally, and is known as the drag curve. In argon-steel system in continuous casting, the bubbles move closely with the liquid, and the bubble Reynolds number is about 125 for a 1 mm bubble. The drag curve for the Allen regime was adopted:

$$C_D = \frac{24}{\text{Re}_b} \left( 1 + 0.15 \,\text{Re}_b^{0.687} \right) \tag{2.19}$$

The drag coefficient for spherical particles continues to decrease monotonically with Reynolds number.

#### **2.3.3 Buoyancy Term in Multi-Phase Momentum Equations**

In the gas-liquid multi-phase system, there is a large difference in density between phases. Bubbles will reach a substantial upward velocity due to the great buoyancy force before they leave top surface. The drag force exerted by bubbles to liquid is substantial as well and will affect liquid flow greatly. Meanwhile, for steel continuous casting, it is hoped all gas float up out of bulk liquid to prevent gas related defects. So, it is important to calculate buoyancy force precisely. Buoyancy force introduced into the K- $\varepsilon$  model momentum equations is calculated as:

$$\mathbf{B}_{gas} = \left(\boldsymbol{\rho}_{liq} - \boldsymbol{\rho}_{gas}\right) \mathbf{g} \tag{2.20}$$

where g is the gravity vector  $\{0, 0, 9.8\}$ . The Boussinesq approximation is employed, whereby gas compressibility and thermal effects on  $\rho$  are ignored [9].

#### 2.3.4 Source Terms in Multi-Phase Model

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In the 2-phase flow in continuous casting, liquid steel and argon gas flow into the mold. Liquid steel solidifies and exits the bottom of the mold. Argon gas escapes from the top surface. Different from previous work [5], in this study, the inlet boundary is set for both liquid and gas phases at the SEN port. No additional source term is set for the gas phase. Because the top surface is set as the wall boundary condition, a sink is set for the gas phase at the top boundary of domain to absorb escaping gas. This sink has no influence on the liquid phase. It adds a mass source and a momentum source to the gas equations at the top surface boundary cells. The equations of the gas phase in other cells except top surface boundary cells remain unchanged. After adding the source term, the continuity equation (Eqn.2.8) for the gas phase at the top surface boundary cells becomes:

$$\nabla(f_{gas}\rho_{gas}V_{gas}) = S_{mass} \tag{2.21}$$

The mass source term  $S_{mass}$  is calculated as:

$$S_{mass} = f_{gas} \rho_{gas} W_{gas} A_w \tag{2.22}$$

where  $W_{gas}$  is the upward velocity of gas phase,  $A_w$  is the projection area of the cell at the top surface.

The momentum equation (Eqn.2.11) for the gas phase at the top surface boundary cells is:

$$\nabla \cdot (f_{gas} \rho_{gas} (V_{gas} \otimes V_{gas})) = \nabla (f_{gas} \mu_{eff} (\nabla V_{gas} + (\nabla V_{gas})^T)) - f_{gas} \nabla p + c_{gas-liq}^{(drag)} (V_{liq} - V_{gas}) + B_{gas} + S_{mom}$$
(2.23)

The momentum source term  $S_{mom}$  is calculated as:

$$S_{mom} = S_{mass} V_{gas} \tag{2.24}$$

#### 2.4 Multi-Size-Group Model

In the preceding multiphase model, the dispersed phase is assumed to have the same size and shape through the domain. In reality, the size and shape of the dispersed phase are in a range due to break-up and coalescence. Bubble size is decided by break-up rate and coalescence rate of bubbles. To precisely simulate gas bubble behavior in the liquid and their influence on liquid flow, different bubble sizes and break-up and coalescence between them should be incorporated into the model.

The Multi-Size-Group Model (MUSIG) provides an approximation to solve this problem. In multi-size-group model, bubbles are divided into finite size groups. Break-up

and coalescence occur among these groups. Various MUSIG models are developed. In this study, Luo's MUSIG model [11] as implemented into CFX4.3 [10] was used. In this model, the population balance method together with the break-up and coalescence models was incorporated into three-dimensional CFD calculations. All bubble velocities are related to the average value. Only one momentum equation is solved for all bubbles. For more simplicity, the MUSIG model assumes that all bubbles are moving at the same velocity. This makes the momentum equation for bubbles much easier to solve. However, the results are rougher, because collision between bubbles moving with different velocity is avoided. Only turbulent collision is considered. All gas bubbles are solved as gas phase in a whole using same equations presented in section 2.3. At certain intervals, such as 20 iterations, MUSIG model is invoked to change the volume distribution among size groups. The volume redistribution among size groups will change average size of the bubbles.

To calculate bubble size redistribution, mass transfer between size groups must be calculated first. The mass transfer equation for size group i:

$$\nabla \cdot (\rho V n_i f_{gas}) = S_i \tag{2.25}$$

where  $S_i$  is the rate of mass transfer into the size group i due to break-up or coalescence.  $f_{gas}$  is volume fraction of gas phase.  $n_i$  is the volume fraction of size group i in total gas phase.

The mass transfer rate into size group i is:

$$S_{i} = \rho v_{i} (B_{b} - D_{b} + B_{c} - D_{c})$$
(2.26)

where  $B_b$  is birth rate of bubble i due to breakup,  $D_b$  is death rate of bubble i due to breakup,  $B_c$  is birth rate of bubble i due to coalescence,  $D_c$  is death rate of bubble i due to coalescence.  $v_i$  is the volume of a single bubble in group i.

Combining Eqn.2.25 and 2.26, the following equation is obtained:

$$\nabla \cdot (Vn_i f_{gas}) = v_i (B_b - D_b + B_c - D_c) \tag{2.27}$$

 $B_b, D_b, B_c$  and  $D_c$  are given in following breakup and coalescence models. Then,  $n_I$  is obtained.

For each update of size distribution by MUSIG, the average bubble size is changed. So, the average bubble size is calculated every time MUSIG is invoked. The Sauter mean diameter is used as the bubble average diameter. The Sauter mean diameter is defined as:

$$d_s = \frac{1}{\sum_i \frac{n_i}{d_i}}$$
(2.28)

 $d_s$  is used in Eqn.2.17 and 2.18 to calculate the drag coefficient C<sub>D</sub> between the liquid and gas phases. MUSIG model affects momentum equation by changing average bubble diameter due to mass transfer between size groups. Changed average bubble diameter has changed drag force term in momentum equations. The projected area, A, of the bubbles in the direction of flow in Eqn.2.16 is calculated by:

$$A = \frac{6f_{gas}}{d_s} \tag{2.29}$$

#### 2.4.1 Break-up Model

In this study, the break-up model is based on Luo and Svendsen's break-up model of bubbles [11] as implemented into CFX4.3 [10]. This model is based on the theories of isotropic turbulence and probability and contains no unknown or adjustable parameters. It matches experimental measurements of the water-air system in a high-intensity pipeline flow<sup>[11]</sup>. In this model, only binary breakage is considered. The break-up rate of bubbles of volume  $v_i$  into bubbles of volume  $v_i$  ( $=v_i f_{BV}$ ) is expressed as:

$$g(v_j:v_i) = f_B 0.923(1 - f_{gas}^{total}) \left(\frac{\varepsilon}{d_j^2}\right)^{\frac{1}{3}} \int_{\xi_{\min}}^{1} \frac{(1 + \xi)^2}{\xi^{\frac{11}{3}}} e^{-\chi_c} d\xi$$
(2.30)

where the coefficient  $f_B$  is added for calibration of the model. In this study,  $f_B$  is chosen to be 1.0 according to Luo and Svendsen [11].  $f_{gas}^{total}$  is total gas volume fraction. The critical dimensionless energy for break-up  $\chi_c$  is calculated as:

$$\chi_{c} = \frac{12 \left[ f_{BV}^{\frac{2}{3}} + (1 - f_{BV})^{\frac{2}{3}} - 1 \right] \sigma}{\beta \rho \varepsilon^{\frac{2}{3}} d_{j}^{\frac{5}{3}} \xi^{\frac{11}{3}}}$$
(2.31)

where  $f_{BV}$  is the breakage volume fraction and is defined as:

$$f_{BV} = \frac{v_i}{v_j} \tag{2.32}$$

where  $v_i$  and  $v_j$  are volume of bubble i and j.  $\beta$  is a constant (2.0) [11].  $\sigma$  is surface tension between phases (1.192 for steel-argon system).

The dimensionless size of eddies in the inertial sub-range of isotropic turbulence  $\xi$  is calculated as:

$$\xi = \frac{\lambda}{d_j} \tag{2.33}$$

where  $d_i$  is the size of bubble j,  $\lambda$  is the turbulence eddy size and is calculated as:

$$\lambda = \left(\frac{\mu_0^3}{\varepsilon}\right)^{\frac{1}{4}} \tag{2.34}$$

The birth rate of group i bubbles due to break-up of large bubbles is:

$$B_b = \sum_{j=i+1}^{N} g(v_j; v_i) n_j$$
(2.35)

The death rate of group i bubbles due to break-up to small bubbles is:

$$D_b = \sum_j^i g(v_i : v_j) n_i$$
(2.36)

where  $n_i$  is the volume fraction of bubble i in total gas phase.

#### **2.4.2 Coalescence Model**

Bubbles collide with each other due to their different moving velocities. Some of the collisions result in coalescence between them. Coalescence occurs in three steps. First, two bubbles collide. Second, liquid film becomes thinning. Third, if they stay together for an enough long time, the liquid film becomes thin enough to corrupt. Coalescence happens. Based on collision mechanism, the coalescence process could be considered associating to three different mechanism, turbulence, buoyancy and laminar shear. In this study, the coalescence model is based on Prince and Blanch's coalescence model [12] as implemented into CFX 4.3 [10]. The coalescence rate of bubbles between size group i and j is given by the total collision frequency multiplied by collision efficiency:

$$Q_{ij} = (\theta_{ij}^T + \theta_{ij}^B + \theta_{ij}^{LS})\eta_{ij}$$
(2.37)

where collision frequency  $\theta_{ij}^{r}$ ,  $\theta_{ij}^{b}$  and  $\theta_{ij}^{LS}$  is corresponding to turbulence, buoyancy and laminar shear mechanism. All three collision frequencies form the total collision frequency. The collision efficiency  $\eta_{ij}$  is calculated as:

$$\eta_{ii} = e^{(-t_{ij}/\tau_{ij})} \tag{2.38}$$

 $t_{ii}$  is the required staying time of two bubbles for coalescence:

$$t_{ij} = \left(\frac{r_{ij}^3 \rho}{16\sigma}\right)^{\frac{1}{2}} \ln \frac{h_0}{h_f}$$
(2.39)

where  $h_0$  is the initial film thickness and  $h_f$  is the critical film thickness where coalescence happens. Value of  $h_0$  and  $h_f$  is given by experimental documents. For water-air system [13,14],  $h_0$  is 10<sup>-4</sup>m and  $h_f$  is 10<sup>-8</sup>m.  $\sigma$  is the surface tension and  $\rho$  is density of the gas.

 $\tau_{ii}$  is the contact time of two bubbles. For the large scale eddy, contact time is calculated :

$$\tau_{ij} = \frac{r_{ij}^{\frac{2}{3}}}{\epsilon^{\frac{1}{3}}}$$
(2.40)

where the equivalent radius  $r_{ij}$  is given by:

$$r_{ij} = \frac{1}{2} \left(\frac{1}{r_i} + \frac{1}{r_j}\right)^{-1}$$
(2.41)

where  $r_i$  and  $r_j$  are radius of bubble i and j.

In this study, only one momentum equation is set for gas phase. All bubbles are assumed having same velocity. So collision due to buoyancy and laminar shear are not presented. Only turbulence collision is considered. Meantime, a calibration coefficient  $f_c$ , called coalescence coefficient, is added. Eqn.2.37 becomes:

$$Q_{ij} = f_c \theta_{ij}^{\mathrm{T}} \eta_{ij} \tag{2.42}$$

where  $f_c$  is given by 0.05 in CFX4.3 [10]. The value may be changed to match different systems and set to 0 to ignore coalescence. The turbulence collision rate is given by:

$$\theta_{ij}^{T} = S_{ij} \left[ u_{ii}^{2} + u_{ij}^{2} \right]^{\frac{1}{2}}$$
(2.43)

where collision cross-sectional area of bubbles  $S_{ij}$  is defined by:

$$S_{ij} = \frac{\pi}{4} \left[ d_j + d_i \right]^2$$
(2.44)

Turbulent fluctuating velocity  $u_i$  of bubble i and j is calculated as:

$$u_t = 1.4\varepsilon^{\frac{1}{3}}d^{\frac{1}{3}} \tag{2.45}$$

The birth rate of bubble group i due to coalescence of group-j and group-k bubbles is:

$$B_{c} = \frac{1}{2} \sum_{j}^{i} \sum_{k}^{i} Q_{jk} n_{k} n_{j}$$
(2.46)

The death rate of bubble group i due to coalescence is:

$$D_c = \sum_{j}^{N} Q_{ij} n_i n_j \tag{2.47}$$

#### **2.5 Heat Transfer Model**

1 1

Steel continuous casting is a heat processing. Heat transfer occurs continuously in a nearly steady uncompressible fluid flow. For the convenience of solution, the energy transport equation for that flow is solved for the enthalpy H instead of temperature T. Then temperature is derived from enthalpy solution. The energy equation takes the form:

$$\nabla \cdot \left(\rho HV\right) = \nabla \cdot \left(k_{eff} \nabla T\right) \tag{2.48}$$

where T is the temperature and where H is the total enthalpy. Because liquid flow is incompressible flow, velocity has little influence on H. H is equal to static enthalpy h:

$$H = h \tag{2.49}$$

Because overheat of steel in continuous casting is normally less than 100 °C, constant pressure specific heat  $C_p$  could be approximated as a constant, the  $C_p$  at solidus temperature of steel. Thus, static enthalpy *h* is given by :

$$h = C_p (T - T_{ref}) \tag{2.50}$$

where  $T_{ref}$  is the solidus temperature of steel. T is the temperature of liquid steel. The effective conductivity  $k_{eff}$  in Eqn.2.48 is defined as

$$k_{\text{eff}} = k_o + k_t \tag{2.51}$$

where  $k_o$  is the molecular conductivity and  $k_t$  is the turbulent conductivity, which is defined as

$$k_t = \frac{C_p \mu_t}{\mathrm{Pr}_t} \tag{2.52}$$

where  $Pr_t$  is the turbulent Prandtl number constant. It is value is given as 1.0 [5].  $\mu_t$  is calculated from Eqn.2.4. The values of  $C_p$ ,  $T_{ref}$  and  $k_o$  are given in specific cases in following sections.

#### 2.6 Boundary Conditions

#### 2.6.1 Inlet Boundary Condition

The inlet boundary condition is set as a multi-phase inlet. All phases are assumed to have the same velocity at the inlet. The volume fraction of each phase is set at the inlet. The total mass flowing into the mold should equal the total mass flowing out of mold. The inlet velocity is calculated according to the mass balance of liquid phase. The normal inlet velocity is:

$$V_x = \frac{V_c WN}{A_{jet} f_{liq}}$$
(2.53)

where  $V_x$  is normal inlet velocity,  $V_c$  is the casting speed, W is width of the outlet of the mold, N is thickness of the outlet of the mold,  $A_{jet}$  is the effective area of the inlet,  $f_{liq}$  is volume fraction of the liquid phase.

Volume fraction f at inlet is calculated from liquid and gas flow rate. Gas is subjected to significant volume expansion during injection because it is heated from room temperature to liquid steel temperature. This expansion must be considered. According to B.G.Thomas and X.Huang [17], the volume fraction of gas  $f_{gas}$  is calculated as:

$$f_{gas} = \frac{Q_{gas}\beta}{Q_{liq} + Q_{gas}\beta}$$
(2.54)

where Q is the flow rate of liquid or gas. The factor  $\beta$  is defined as:

$$\beta = \frac{T_0 P_\infty}{T_\infty (P_\infty + \rho g L_n)}$$

where  $T_0$  is liquid temperature at the inlet,  $T_{\infty}$  is room temperature,  $P_{\infty}$  is room pressure,  $\rho$  is liquid density, g is gravity,  $L_n$  is submergence depth of the nozzle. All the above values are given in the specific cases in following sections.

#### 2.6.2 Outlet Boundary Condition

The outlet boundary conditions of both the nozzle and mold are set as pressure boundary condition. A reference pressure, which is set as 0 in this study, is set for pressure. Neumann boundary conditions, i.e. zero normal gradients, are imposed on all other variables.

#### 2.6.3 Symmetry Boundary Condition

For the symmetry boundary, the velocity component normal to the boundary is set to zero. Zero normal gradient boundary conditions are set to other variables.

#### **2.6.4 Wall Boundary Condition**

The wide face and narrow face walls are set to non-slip wall boundary conditions. For continuous caster, the top surface is also set to the non-slip wall boundary condition. In the multi-phase model, a sink is added to the top surface to absorb the escaped gas. The influence of the sink on the governing equations is shown in section 2.3.4. For the water model, the top surface is a free surface. So it is set to the slip wall boundary condition.

Velocity and temperature vary tremendously near the wall. To give accurate enough description of variables near the wall, very fine meshes have to be needed. This requires great memory and time to implement by computer. Wall laws provide an approximate method to avoid this. Wall laws calculate the high gradients of velocity in the boundary layer region using an empirical correlation based on the shear stress at the wall. This allows a coarse mesh to be used and still include the behavior in the boundary layer. The boundary condition for velocity is specified using wall laws in the standard K- $\epsilon$  model. In the following sections, the velocity field is solved using the standard K- $\epsilon$  model. The wall law is used to deal with wall boundary conditions.
In the standard K- $\varepsilon$  model, tangential velocity to the wall is calculated as:

$$V_{t} = \begin{cases} -(C_{\mu}K)^{1/2} y^{+} & y^{+} < y_{0}^{+} \\ \frac{-(C_{\mu}K)^{1/2}}{\kappa} \log(Ey^{+}) & y^{+} \ge y_{0}^{+} \end{cases}$$
(2.55)

where  $V_t$  is the tangential velocity to the wall,  $C_{\mu}$  is a turbulence model constant (0.09) [16]. *K* is turbulence kinetic energy.  $\mu_0$  is molecular viscosity. *n* is normal distance to the wall. *E* is log-layer constant (9.7930) [9].  $\kappa$  is the Von-Karmen constant (0.419) [9]. The non-dimensional distance normal to the wall is:

$$y^{+} = \frac{(\rho^2 C_{\mu}^{-1/2} K)^{1/2}}{\mu} n$$
(2.56)

The cross over point between the viscous sub-layer and the logarithmic region is:

$$y_0^+ = \frac{1}{\kappa} \log(Ey_0^+)$$
(2.57)

which obtains the value  $y_0^+=11.23$ . The turbulence kinetic energy *K* is solved in the control volume adjacent to the wall using transport equations for K- $\varepsilon$  model. The production terms in the K equation are treated differently so that only quantities interior to the flow and the specified boundary conditions on velocity are used. To catch the profile near the wall, the non-dimensional distance of the first cell should be smaller than  $y_0^+$  in Eqn.2.57 [9].

The turbulence dissipation at the wall is calculated from K using the relation:

$$\varepsilon_{(y^{+}=0)} = \frac{C_{\mu}^{3/4} K^{3/2}}{\kappa n}$$
(2.58)

The boundary conditions for heat transfer calculations are similar to the velocity

boundary conditions.

The enthalpy in the wall layer is assumed to be:

$$H = \begin{cases} \frac{\Pr_{H} E_{H} y^{+} \left(\frac{\partial H}{\partial n}\right)_{y^{+}=0}}{\rho C_{\mu}^{1/4} K^{1/2}}, & \text{for } y^{+} \leq y_{H}^{+} \\ \frac{\Pr_{t}}{\kappa} \log \left(E_{H} y^{+}\right) \left(\frac{\partial H}{\partial n}\right)_{y^{+}=0}}{\rho C_{\mu}^{1/4} K^{1/2}}, & \text{for } y^{+} > y_{H}^{+} \end{cases}$$

$$(2.59)$$

where  $y_H^+$  is the cross over point between the viscous sub-layer and the logarithmic region for enthalpy (31.77)[10], and  $E_H$  is the formula of Jayatilleke

$$E_{H} = E \exp\left[9.0\kappa \left(\left(\frac{\mathrm{Pr}}{\mathrm{Pr}_{H}}\right)^{0.75} - 1\right)\left(1 + 0.28 \exp\left(-0.007 \frac{\mathrm{Pr}}{\mathrm{Pr}_{H}}\right)\right)\right]$$
(2.60)

where Pr is the fluid Prandtl number ( $C_p \mu_o / k_o$ ) and  $Pr_H$  is the Prandtl number for enthalpy(0.9) [10]. E has the same value as that in Eqn2.55.

## 2.7 Solution Method

A multi-block numerical grid with body-fitted coordinates is used to create the complex geometry of the nozzle and mold domain. The details of geometry of specific cases are discussed in following sections. The governing equations for above models are discretized using the finite difference method and solved using the commercial finite difference program CFX version 4.2 or 4.3 by AEA Technology [9, 10]. The discretized equations are solved using different iterative solution algorithms [9, 10]]. The velocity and volume fraction equations use full field Stone's algorithm. The K and  $\varepsilon$  equations use line relaxation algorithm. The pressure equation uses preconditioned conjugate gradients

algorithm. The solutions are considered converged when the scaled residuals of the most concerned variables are about 10<sup>-2</sup>. The scaled residual is the ratio of the residual at the current iteration to that of the second iteration. For single phase simulations, the relaxation factors for velocity, mass, K and  $\varepsilon$  are normally set to 0.3 and that for pressure is set to 0.8. For multiphase simulations, relaxation factors for pressure and volume fraction are set to 0.8 and relaxation factors for velocity, mass, K and  $\varepsilon$  are set to 0.005-0.01 for the first 200 iterations and then increased to 0.3 for 1000 iterations or until it is converged. The program is run on the SGI Origin 2000 supercomputer at NCSA at university of Illinois. The CPU time varies with the number of cells and convergence criterion. For a single phase water model (section 3), 90 minutes may be needed to achieve a fully convergence. For a multi-phase steel caster with MUSIG model (section 6), 25 hours may be needed. The CPU time for specific cases are given in following sections when analyzing the cases.

## Section 3

# Validation of Fluid Flow Models

## **3.1 Introduction**

In the preceding section, model formulation was discussed. This section will validate the single phase turbulent K- $\epsilon$  model by comparing modeling results with experimental measurements. Because velocity measurements are difficult in a real caster, to validate the model, experiments and simulation were conducted on scaled water models. Validation with a water model is reasonable because of the hydrodynamic similarity existing between water and liquid steel which is characterized by Reynolds and Froude numbers [4]. In this section, scaled simulation was also compared with full size simulation to evaluate proportional scaling. An alternative scaling method, Froude scaling, is also evaluated.

## **3.2** Comparison of Simulations and PIV Measurements

#### **3.2.1 Scaled Water Model Experimental Description**

Water model experiments are widely applied in steelmaking industry for two major reasons. First, there is dynamic similarity between the water and steel flow as characterized by the Reynolds and Froude numbers. These two dimensionless numbers account for the relative influences of inertial, gravitational and viscous forces so that the internal fluid flow between the water and steel can be comparable. Second, a water model is easy to operate and allows direct observation of the liquid flow so that simulated steel flow can be evaluated [18]. A full size water model with the same size mold as a real continuous caster is very large and inconvenient to be built in the lab. So scaled water models are often set up for experiments. This study uses a 0.4 scaled water model constructed by LTV steel. The mold dimensions and other details are given in Figure 3.1. There are three round 39-mm diameter outlet holes at the bottom of the mold. Tiny particles (50-100  $\mu$ m) are added into the water to trace the flow. Particle Image Velocimetry (PIV) technology is used to trace these particles to give velocity field images of the flow. PIV provides practical quantitative turbulence information on a plane section. This makes it easy to compare simulation and water model quantitatively.

In some measurements, air was injected at a proportion same as that in practice to simulate argon injection through the submerged entry nozzle at the real continuous caster. PIV data were processed by computer and compared with simulation results.

### **3.2.2 Numerical Simulation of 0.4 Scaled Water Model**

Numerical simulations of 0.4 proportionally scaled water model used the single phase standard K- $\epsilon$  model. The geometric, dynamic and other parameters in numerical simulations are given in Table 3.2 and Figure 3.1 and are similar to the PIV measurements except for approximating the 39-mm round outlet as 35×35 mm square holes which have the same area. This treatment is reasonable because it has little influence on flow behavior at the upper part of the mold where defects are believed to form. Because of the symmetry of the mold, only one quarter of the mold is analyzed. Figure 3.1 shows the quarter of 0.4 scaled water model. Water goes in from the SEN (Submergence Entry Nozzle) and goes out from the three holes at the bottom. In the numerical simulation, total number of cells is 41,960. The boundary conditions are shown in Table 3.1. Table 3.2 shows the data used in simulation. The simulation runs 1200 iterations spending 1.56 CPU hours to reach the acceptable convergence. The convergence plot (Figure 3.2) shows that residuals go down to 1% of the initial level. This suggests a good convergence.

#### 3.2.3 Proportional Scaling and Froude Scaling

The water model is set up according to similarity between water and steel flow which is characterized by the Reynolds number and Froude number. To decide the correct flow rates in the scale water model, a straightforward scaling method is proportional scaling which follows  $V_{model} = V_{real}(\frac{L_{model}}{L_{real}})$ . This is based on matching Reynolds number  $\text{Re} = \frac{\rho VL}{\mu}$ . Another alternative scaling mechanism is Froude scaling which is based on Froude similarity. Froude number is defined as  $Fr = \frac{V^2}{gL}$ . According to Froude similarity,  $\frac{V_1^2}{gL_1} = \frac{V_2^2}{gL_2}$ . So Froude scaling produces  $V_{model} = V_{real}\sqrt{\frac{L_{model}}{L_{real}}}$ . In this section, Both of these two scaling methods are compared.

Simulations of a 0.4 proportionally scaled water model, a 0.4 Froude scaled water model and a full size water model are implemented in this section. The 0.4 proportionally scaled model and the 0.4 Froude scaled model have the same domain but inlet velocities. Figure 3.1 shows the geometry of the domain for these cases. The data for these three simulations are listed in Table 3.2. Boundary conditions are all same for three cases as listed in Table 3.1. The convergence history for three cases is shown in Figure 3.2. The residuals for velocities in three cases all go down to less than 2% of original value. This indicates all velocities converged well.

Flow patterns of three cases are compared in Figure 3.3. Flow in Froude scaled water model is more similar with that in full size water model than that in proportionally scaled water model. The magnitude of velocities in Froude scaled water model is larger than that in the proportionally water model because the velocities are scaled by the square root of the ratio between geometries according to Froude method so that its inlet velocity is larger. The positions of the circulation eyes and impingement points are listed in Table 3.3. From the table, the geometric features of flow pattern in the Froude scaled water model are slightly closer to that in the full size water model than that in the proportionally scaled water model.

Figure 3.4 and Figure 3.5 compare the speed along several typical directions between the above three simulations. The speeds are scaled by their respective inlet speeds to provide comparability. The dimensionless speed is defined as  $V^*_{dimensionless} = \frac{V}{V_{miner}}$ . The dimensionless distance below the meniscus is defined as  $Z^*_{dimensionless} = \frac{Z}{Z_{domain}}$ , where  $Z_{domain}$  is the height of the domain and Z is the distance below the meniscus. The dimensionless distance from SEN is defined as  $X^* = \frac{X}{X_{domain}}$ , where  $X_{domain}$  is half of the mold width and X is the horizontal distance from SEN. The quantitative comparisons did not show too much different between flows in Froude scaled water model and proportionally scaled water model. Although there is a difference between the flow patterns in the Froude scaled water model and the proportionally scaled water model, the dimensionless velocities remain almost identical at most areas in the

domain. These two different scaling methods do not bring in too much difference to the flow.

# 3.2.4 Comparison of Numerical Simulation Results and PIV Measurements on 0.4 Scaled Water Model

In the previous section, it shows that proportional scaling (based on Reynolds number) and Froude scaling (based on Froude number) have similar computed results after dimensionless treatment by inlet speed. This simulation uses the same inlet velocities and flow rate as those of LTV water model experiments which adopted Froude scaling. Model predictions are validated in several ways: comparison with flow pattern on the centerplane parallel to wide face, speed and velocities along jet and velocities along centerline between the SEN and narrow face. Figure 3.6 compares the computed flow pattern on the centerplane parallel to the wide face with corresponding PIV measurements. The flow pattern of numerical simulation is qualitatively very similar to that in the water model. The positions of the upper and lower eyes of flow and impingement points for both cases are very close (Table 3.3).

To quantitatively compare the computed and measured flow in the water model, velocities along two lines are compared. Figure 3.7 compares velocities along the center line of the jet. The speed, horizontal velocity and vertical velocity of numerical simulation match those of the water model within +/- 10%. The difference in vertical velocity at the end of the jet between numerical simulation and PIV measurements is because the simulated jet bends up a little bit and the straight line with initial jet angle can not catch the jet at the end. The data at the end does not reflect the feature of the jet. Figure 3.8 shows comparison of velocities along center line between SEN and narrow

face. In this case, only the region above the jet is processed because this region is most concerned. The horizontal velocity changes form negative to positive with increasing depth from the meniscus because of its double-roll flow characteristic. The vertical velocity remains low because the center line goes through the upper eye. Vertical velocity above and below the upper eye is very small (Figure 3.6), Both horizontal velocity and vertical velocity of numerical simulation matches PIV measurements well. The difference in vertical velocity between the numerical simulation and the PIV measurements is due to the offset of the upper eye in the simulation. In the numerical simulation, the center line does not go through exactly the upper eye. So the vertical velocity along the center line is a little higher than PIV measurements. The discrepancies are not larger than 5%.

## 3.3 Numerical Simulation of Flow in Real Continuous Caster

The water model is often scaled from the real continuous caster to avoid the inconvenience in constructing a full size mold. However, similarity between the flow in the water model and the real caster has not been verified. The fluid is different. The geometry is also different. Thus, the flow in the water model and real caster might be different in some respects. Because of the difficulty in measurements in continuous caster, it is hard to collect adequate accurate data of the flow in the caster. The similarity is hard to be verified directly. However, if the simulations of the water model and the real caster match, that would increase confidence in the scale water model. The relationship is shown in Figure 3.10. If relation 1 and 2 can be validated, then relation 3 and 4 could be inferred based on similarity.

Figure 3.1 shows the geometry of the full size mold of continuous caster. The simulation domain for the real full size mold measures  $1825 \times 200 \times 3000$  mm, which is 2.5 times larger than water model except for the longer length of 3000 mm. The boundary conditions are the same as those for 0.4 scaled water model except the top surface is set non-slip wall which considers the less mobile flux layer. Data used in simulation are listed in Table 3.2. The simulation runs in 1800 iterations spending 6.24 CPU hours to reach acceptable convergence. The scaled residual plot is shown in Figure 3.2. The convergence plots show that residuals of all variables go down to 3% of the initial level.

Figure 3.9 shows comparison of the flow patterns on the centerplane parallel to wide face of the 0.4 scaled water model simulation with that of the full size continuous casting simulation. The two cases have similar flow patterns. The positions of the upper eye and the impingement point match fairly well (Table 3.3). The position of the lower eye is different because the outlet condition is different. For the water model, the holes are used to drain the water. For the real caster, the whole bottom of the domain is outlet. The total area of the outlet in the water model is smaller than that in the real caster. Therefore, the outlet velocity of the water is relatively larger. This lowers the position of the lower eye in the real caster.

The quantitative comparison of full size continuous casting simulation with 0.4 scaled water model simulation is shown in Figure 3.7 and Figure 3.8. The velocities and speed are scaled by the corresponding inlet velocity to provide comparability. Along the jet direction, the two cases show almost identical results for velocity vectors. Along the center line between the narrow face and the SEN, velocities also match each other well. The largest difference is less than 5%.

From comparison of the full size continuous casting simulation with the 0.4 scaled water model simulation, the two simulations have very similar results. Because the simulation of the 0.4 scaled model matches the PIV measurements as well, the simulation of the full size continuous caster hopefully matches the flow in the real caster.

## **3.4 Conclusions**

- Numerical K-ε model simulations of flow in the 0.4 proportionally scaled water model match PIV measurements both qualitatively and quantitatively.
- 2. Dimensionless comparisons of the flow simulations in the 0.4 proportionally scaled water model, full size water model and liquid steel flow in the full size continuous caster show that these flows are very similar. The standard K-ε model which is validated in 0.4 scaled water model could be applied to model a full size continuous caster with reasonable accuracy.
- 3. Analysis of flow in Froude scaling and proportionally scaling show that these two scaling methods produce negligible differences at least for single phase flow.

Boundary	Boundary Condition	Boundary	Boundary Condition
Inlet Port	Inlet	Top Surface	Slip Wall
Narrow Face	Non-slipWall	Centerplane	Symmetry
Wide Face	Non-slip Wall	Outlet Holes	Pressure

Table 3.1 Boundary conditions in numerical simulation

Table 3.2 Data used in simulations

	0.4 Proportionally	0.4 Froude	1:1 Water	1:1 Liquid
	Scaled Water	Scaled Water Scaled Water		Steel Caster
	Model	Model		
Domain Width W (mm)	730	730	1825	1825
Domain Thickness D (mm)	80	80	200	200
Domain Height H (mm)	950	950	2375	3000
Nozzle Submergence Depth	80	80	200	200
(top surface to top of port)				
Jet Angle	30° down	30° down	30° down	30° down
Inlet Jet Spread Angle	$0^{\circ}$	0°	0°	0°
Inlet Speed, $V_c$ ( <i>m</i> / <i>s</i> )	0.4239	0.6703	1.0597	1.0597
Inlet Velocity, $V_x(m/s)$	0.3671	0.5805	0.9178	0.9178
Inlet Velocity, $V_z(m/s)$	0.2119	0.3351	0.5298	0.5298
Inlet Turbulent Kinetic	0.0440	0.0440	0.0440	0.0502
Energy, $K_o(m^2/s^2)$				
Inlet Turbulence Dissipation	0.999	0.999	0.999	0.457
Rate, $\varepsilon_{o}(m^{2}/s^{3})$				
Liquid Density, $\rho$ (kg/m <sup>3</sup> )	1000	1000	1000	7.02×10 <sup>3</sup>
Liquid Viscosity $(m^2/s)$	1×10 <sup>-3</sup>	1×10 <sup>-3</sup>	1×10 <sup>-3</sup>	5.55×10 <sup>-3</sup>

Dimensionless	PIV measurements	0.4 Froude	0.4 proportionally	Full size	Real caster
	on 0.4 scale	scaled	scaled water	water model	simulation
position	water model	water model	model		
(x*, z*)					
Upper eye					
x*	0.638	0.639	0.694	0.643	0.610
Z*	0.118	0.113	0.114	0.116	0.095
Lower eye					
2					
x*	0.552	0.634	0.634	0.571	0.559
Z*	0.433	0.532	0.413	0.526	0.459
Impingement					
point					
x*	1.000	1.000	1.000	1.000	1.000
Z*	0.231	0.231	0.234	0.229	0.210

 Table 3.3 Dimensionless position of circulation eyes and impingement point



Figure 3.1 Simulation domain for three molds





Residuals for 0.4 proportionally scaled mold simulation





Figure 3.2 Convergence history of simulations of 0.4 proportionally scaled water model, 0.4 Froude scaled water model, full size water model and full size real continuous casting

Distance unit : m



Figure 3.3 Comparison of computed flow patterns between 0.4 Froude scaled, 0.4 proportionally scaled and full size single phase water model



Speed along Center Line between SEN and NF

Figure 3.4 Comparison of speed along jet and centerline between 0.4 Froude scaled, 0.4 proportionally scaled and full size water model





Speed along Horizontal Line 0.05 m below Meniscus

Speed along Horizontal Line 0.11 m below Meniscus



Figure 3.5 Comparison of speed along horizontal lines between 0.4 Froude scaled, 0.4 proportionally scaled and full size water model



Figure 3.6 Comparison of flow patterns between PIV measurements and simulation of 0.4 scaled water model



Figure 3.7 Comparison of velocities along jet between simulation and PIV measurements of 0.4 scaled water model and real caster



Horizontal Velocity along Center Line between SEN and NF

 $(V_{x0}=0.3671 \text{m/s for } 0.4 \text{ scaled model}, V_{x0}=0.9178 \text{m/s for real caster})$ 



Vertical Velocity along Center Line between SEN and NF

 $(V_{z0}=0.2119 \text{ m/s for } 0.4 \text{ scaled model}, V_{z0}=0.5298 \text{ m/s for real caster})$ 

Figure 3.8 Comparison of velocities along centerline between simulation and PIV measurements of 0.4 scaled water model and real caster



Figure 3.9 Comparison of flow patterns between simulations of continuous casting and 0.4 proportionally scaled water model



Figure 3.10 A schematic of relationship between numerical simulation, scaled water model and real caster

## Section 4

# Validation of Heat Transfer Models

## 4.1 Introduction

Liquid steel is poured into the mold at a temperature above the liquidus. The sensible heat contained in the liquid steel represented by this temperature difference is known as the superheat. The superheat must be transferred out of the steel mass before it can solidify. The heat transfer behavior in the mold affects solidified microstructure and formation of defects, such as meniscus hook formation, breakouts, centerline segregation etc. A good understanding of heat transfer in the mold may help to improve product quality.

Heat transfer in continuous casting is associated with flow in the mold. The complex nature of heat transfer in the continuous casting mold and the high temperature phenomena make it very difficult to analyze it by traditional methods such as water model and plant measurements. Numerical simulation provides a feasible method. In this work, assuming flow and heat transfer is steady in the liquid pool, temperature is computed by solving the single phase Navier-Stokes equations and corresponding Eulerian heat conduction equation (see section 2.4). The heat flux on the wide face and narrow face is also computed based on the temperature gradient. By adjusting factors such as nozzle geometry, thickness of the powder/flux layer or other parameters, the temperature distribution in the liquid pool and the corresponding heat flux to the shell may be changed as desired. These changes are easy to make and evaluated in the mathematical simulation before expensive experiments are conducted.

In this study, three-dimensional turbulent flow and the associated heat transfer is modeled with the finite difference code CFX4.2, using the standard K- $\epsilon$  model with a user defined wall law. To validate the results of the simulation of heat transfer, temperature at different positions in the upper liquid pool is measured in an operating caster and compared with the simulation.

## 4.2 Experiments

Temperature distribution in the liquid steel pool is difficult to measure. In the experiment in this work, an apparatus was constructed by AK Steel's Research Center in Middletown, OH to measure vertical temperature profiles in the liquid pool [21]. The apparatus allows a thermocouple to go down and up vertically at a slow speed (0.6mm/sec) so that thermal equilibrium could be approached at the thermocouple tip. The maximum insertion depth of the thermocouple is 180mm. Temperature signals are transmitted to a computer for data processing. The computer stored the temperature associated with insertion depth. Each cycle of insertion and withdraw makes a single data set. The instrument schematic is shown in Figure 4.1.

Figure 4.2 shows a typical set of the measured data. From the changes of the temperature gradient (slope of this plot), it is possible to identify the powder/flux and flux/steel boundary. This is because these different materials have different heat conduction coefficient. The temperature slope in different materials should have steepness inversely proportional to their different thermal conductivity. In Figure 4.2, slope 1, 2 and 3 should represent powder, flux and steel respectively. From Figure 4.2,

the position of top surface of steel could be determined. This method is used in data analysis in this section.

Table 4.1 lists the common parameters used in the experiments. All other experimental conditions are listed in Table 4.2. The first five measurements have the same submergence depth (130 mm) and represent different positions in the mold. The other two measurements repeat at two positions with different submergence depth (159 mm). To compare measurements under the same experimental conditions, only the first five measurements are considered. All five points are located along the center plane between the wide faces. The 5 locations are 50 mm from SEN, 125 mm from SEN, the middle point between SEN and the narrow face, 125 mm from the narrow face and 50 mm from the narrow face (Figure 4.3). The pouring temperature is about 1559°C (2838 F). The casting speed is 25.4mm/sec (60 inch/min).

## **4.3 Numerical Simulation of Heat Transfer**

Modeling of heat transfer in the liquid pool is based on the results of modeling of flow. The commercial package CFX4.2 is used to solve the Navier-Stokes equations (Eq.2.10 in section 2) and corresponding Eulerian heat conduction equation (Eq 2.25). The model standard K- $\varepsilon$  model with user defined wall law (described in section 2.5.4) is used. Considering the symmetry of the mold, one quarter of the mold is modeled and the solidified shell is not considered. The geometry was set up by D. Creech in previous work [5]. The boundary conditions along the narrow faces and wide faces are set as constant temperature, non- slip walls. The temperature of the wall is set at the liquidus temperature of the steel corresponding to the solidification interface between the solidified shell and the liquid. The top surface is set as the constant heat flux wall

corresponding to the assumption of constant thickness of powder layers. The inlet temperature is set to the pouring temperature. The outlet temperature is set at the liquidus temperature and Table.4.3 lists parameters used in heat transfer simulation. Figure 4.4 shows the convergence history of the heat transfer simulation. The residual of enthalpy reflects convergence level of temperature. In Figure 4.4, the residual of enthalpy is reduced by 4 orders of magnitude. This means temperature converged to an excellent level.

### **4.4 Modeling and Experimental Results**

The velocity on the center plane between wide faces is shown in Figure 4.5. The predicted temperature distribution on the center plane based on the flow modeling is shown in Figure 4.6. Figure 4.7 to Figure 4.11 compares the experimental and predicted temperatures at those positions. Each figure shows the entire temperature signal recorded by the thermocouple from the beginning to the end of each experiment and the corresponding superheat. The measurement at 125 mm from the narrow face has a lower temperature than the other measurements (Figure 4.10) because the ladle was being changed during the measurement and temperature in the tundish is only 1551°C, 8°C lower than normal. The thermocouple hit the shell and broke while measuring at 50 mm from the narrow face so temperature was recorded in only one direction. Measurements were lost during other experiments also (Figure 4.7). From Figure 4.7 to Figure 4.11, the predicted temperature matches the measurements well. The largest difference is less than 4°C, which is within a 7% error with respect to superheat.

In Figure 4.6, two jets from the SEN create a high temperature zone respectively. The right jet impinges the narrow face and forms a peak of temperature and heat flux on the narrow face. Then the jet splits into two flows. One goes down to form a relatively high temperature area along the narrow face. Another flow goes up to meet meniscus and heat up upper part of the liquid pool. The lowest temperatures appear at the upper and lower recirculation eyes. The corner of the meniscus and the narrow face is another low temperature zone. This low temperature zone is dangerous because hooks might form in this area if the liquid temperature drops below the solidus temperature of steel. Flow parameters should be adjusted to prevent this happening.

Heat flux along center line of the narrow face could be used to calculate shell growth by shell solidification model CON1D [15]. Comparison of measured and predicted shell growth can validate prediction of heat transfer [21]. Figure 4.12 shows the predicted heat flux along three vertical lines. From the plot, most of the superheat is conducted through narrow faces. The peak superheat flux is about 1500 kw/m<sup>2</sup>.

## **4.5 Conclusions**

- In this study, an efficient model of heat transfer for 3-D turbulent flow in a thin slab caster has been developed. This model uses the standard K-ε model, the Eulerian heat conduction equation and a wall law to simulate fluid flow and heat transfer.
- Temperature along 5 vertical lines on the center plane between wide faces of the liquid steel pool is measured with thermocouple probes.
- 3. Comparisons of measured and predicted temperature show that the modeling results match experimental results within 7%. The heat transfer model developed in this section can be applied to simulate temperature in the liquid pool and heat flux through solidified shell.

Parameters	Regular Value				
		Total	Left	Right	
Mold Width (mm)	Тор	970.25	485.11	485.14	
	Bottom	958.85	479.44	479.41	
Mold Thickness (mm)	132.08				
Casting Speed (mm/sec)	25.4				
Gas Inflow Rate	No Gas				
Nozzle Port Height (mm)	77				
Nozzle Port Width (mm)	32				
Nozzle Wall Thickness (mm)(near port)		17.5			
Nozzle Bore Inner Diameter (mm) (near port)	70				
Nozzle Submergence Depth (mm)	129.54				
Nominal Vertical Angle of Port Edges	15°C				
Inlet Jet Spread Angle	0°C				

Table.4.1 Parameters of mold and nozzle

No.	Position of Measurement	Pouring Temperature	Submergence Depth	Thickness of Powder (mm)	Thickness of Flux (mm)	Casting Speed
1	Midway of NF vs SEN (295 mm from CL)	1832 K (2838 F)	129 mm (5.1 inch)	60	6	25.4 mm/s (60 IPM)
2	50mm from SEN (150mm from CL)	1831 K (2836 F)	129 mm (5.1 inch)	62	10	25.4 mm/s (60 IPM)
3	50mm from NF (440 mm from CL)	1831 K (2836 F)	129 mm (5.1 inch)	83	5	25.4 mm/s (60 IPM)
4	125mm from SEN (225mm from CL)	1831 K (2836 F)	129 mm (5.1 inch)	68	7	25.4 mm/s (60 IPM)
5	125mm from NF (365mm from CL)	1824 K (2824 F)	129 mm (5.1 inch)	68	5	25.4 mm/s (60 IPM)
6	50mm from NF (440mm from CL)	1826 K (2827 F)	159 mm (6.25 inch)	28	5	25.4 mm/s (60 IPM)
7	Midway of NF vs SEN (295 mm from CL)	1828 K (2830 F)	159 mm (6.25 inch)	34	7	25.4 mm/s (60 IPM)

# Table.4.2 Experimental parameters

Mold Width	984 mm		
Mold Thickness	132 mm		
Nozzle Submergence Depth (top surface to top of port)	127 mm		
Nozzle port Angle	15° down		
Inlet Jet Spread Angle	0°		
Casting Speed, V <sub>c</sub>	25.4 <i>mm/s</i>		
Inlet Velocity, V <sub>x</sub> ,V <sub>y</sub>	0		
Inlet Velocity, V <sub>z</sub>	0.857 <i>m/s</i>		
Inlet Turbulent Kinetic Energy, K <sub>o</sub>	$0.00425 \ m^2/s^2$		
Inlet Turbulence Dissipation Rate, $\varepsilon_{o}$	$0.020 \ m^2/s^3$		
Liquid Steel Density, p	7020 $kg/m^3$		
Laminar Conductivity, k <sub>0</sub>	26 kg-m/ s <sup>3</sup> -K		
Specific Heat, C <sub>p</sub>	$680 \ m^2 / \ s^2 - K$		
Laminar Viscosity, $\mu_0$	0.0056 kg/m s		
Pouring Temperature	1832 <i>K</i>		
Liquidus Temperature	1775 K		
Heat Flux from Top Surface	12,000 W/m <sup>2</sup>		

## Table.4.3 Parameters used in heat transfer simulation



Figure 4.1 Equipment used in measuring temperature in liquid pool



Figure 4.2 Temperature slope of three different material layers



Figure 4.3 Positions of measuring points



Figure 4.4 Residuals of heat transfer modeling



Figure 4.5 Predicted velocity on center plane between wide faces



Figure 4.6 Predicted temperature on center plane between wide faces


Figure 4.7 Measurement and modeling at 50 mm from SEN



Figure 4.8 Measurement and modeling at 125 mm from SEN



Figure 4.9 Measurement and modeling at midway between SEN and narrow face



Figure 4.10 Measurement and modeling at 125 mm from narrow face



Figure 4.11 Measurement and modeling at 50 mm from narrow face



Figure 4.12 Heat flux along three typical lines

# Section 5

# Effect of Gas Bubble Size on Fluid Flow in Continuous Casting Mold

## 5.1 Introduction

Argon gas injection is a common practice in continuous casting. Argon gas is applied at several stages in the continuous casting process to encourage mixing, to help prevent nozzle clogging, and to promote the flotation of solid inclusion particles from liquid steel. However, gas bubbles change the flow pattern greatly [17, 22] and might cause some new problems.

Bubble size is one of the major factors which influence flow pattern. B.G.Thomas [17] et al investigated the effects of bubble size on flow pattern and heat transfer and gained important results. But only two different bubble sizes were present in this work. This limits the detailed understanding of effects of bubble size on flow pattern and other parameters. Another possible drawback of this work is that, to simplify the calculation, no momentum equation was solved for the gas phase. Instead, the bubbles were assumed to have a constant velocity in the vertical direction. This assumption might bring an inaccurate description of movement of bubbles and its influence on liquid flow because the distance might be too short to reach a terminal velocity.

In this study, a three-dimensional gas-liquid multiphase turbulent flow model is developed. Coupled liquid phase and gas phase equations are solved simultaneously so that movement of bubbles can be predicted more accurately.

With this model, the effects of bubble size on the flow pattern and other parameters are investigated. In this section, only one bubble size is set for gas phase. So

all gas bubbles are identical. To investigate the importance of difference in bubble size expected to exist in the liquid steel, seven cases are run, each with a constant bubble size ranging from 0.5mm to 4.0mm. This range covers most of the bubbles forming in the liquid steel according H. Bai's simulation [16]. For each case, the gas volume fraction is kept at 8.5%. To compare with single-phase flow, flow without gas is also modeled.

The flow pattern has a great influence on the formation of inclusion defects such as slivers and pencilpipes. Using these simulation results, the effects of bubble size on the formation of gas entrapment defects and inclusion particles are also investigated.

## **5.2 Model Description**

In this study, a 3-D gas-liquid multiphase turbulent flow model is developed. Using CFX 4.2, the standard K- $\varepsilon$  model is used to solve the turbulent liquid steel flow in the caster mold. Only large scale turbulent eddies could be simulated in this model. The gas phase is assumed to be laminar. An Eulerian multiphase multi-fluid model is used to simulate the motion of gas bubbles in the liquid steel. Each phase has its own velocity field. But only liquid phase has K and  $\varepsilon$  fields. The pressure field is shared by all phases. The velocity fields of the bubbles and liquid are coupled by an empirical inter-phase drag model. Details of this model are addressed in section 2.

The gas bubbles and particles concentrate mainly in the upper liquid pool in the caster. To save computing time, the length of the domain is set 3m from meniscus instead of the entire liquid pool. Again, considering the symmetry of the geometry, only a quarter of the strand is modeled.

At the nozzle port inlet boundary, the liquid and gas have the same velocity. The fraction of each phase is also specified. Inlet pressure values are extrapolated from downstream.

The meniscus is set as the non-slip wall boundary. Because almost all of the argon gas escapes from the meniscus (very little gas is entrapped in the slab), a gas sink is set at the top surface of the domain. A user subroutine is written to set up to absorb the escaping gas [subroutine degas in Appendix D].

The wide face and narrow face is set as non-slip wall boundary. According to the work of D. Creech [5], the solidified shell of wide face and narrow face can be neglected in the modeling without excessive error. The movement of the shell is also neglected. The velocity of the shell is set to zero.

The bottom of domain is set to a pressure boundary. Previous work showed that this condition can handle outlets where the flow is not fully developed better than the mass flow boundary [23]. Normal gradients for all variables are set to zero at the pressure boundary. In this study, the pressure is also specified to equal zero.

The two symmetry planes are set to symmetry boundary. The velocity normal to the plane is set to zero. For all other variables, including pressure, zero normal gradients are specified at symmetry boundary.

A schematic of the simulation domain with boundary conditions and a typical mesh are shown in Figure 5.1. The parameters used in the simulations are listed in Table 5.1. A single phase case with the same casting speed and the same geometry in Table 5.1 is computed first. All 2-phase flow simulations start from the results of the single-phase simulation as an initial guess. This method makes convergence of the 2-phase simulations

much easier and faster. The single-phase simulation converged in 1500-200 iterations, needing 5-8 hours on NCSA Origin 2000 supercomputer. The 2-phase simulations converged in 2000-3000 iterations, needing 20-30 hours on the same supercomputer.

## **5.3 Results and Analysis**

### 5.3.1 Effect of Bubble Size on Flow Pattern

Gas injection affects the casting process through its influence on the liquid flow pattern. The extent of this effect depends on factors such as the gas injection rate, the bubble size, jet velocity etc. In this study, all variables are fixed except bubble size. Mean bubble diameter changes from 0.5mm to 4.0mm, which covers most of the bubble size range met in practice. The effect of bubble size on the flow pattern is shown in Figure 5.2. Small bubbles (d<1.5mm) have greater influence on the liquid flow than big bubbles (d>1.5mm). They tend to stay in the main jet of the flow longer and bend all the flow upward to impinge the top surface, thus tending to form a single roll flow. Big bubbles tend to float up faster and split the jet. In this case, the jet divides into two branches of flow. One branch floats up quickly when the jet enters the domain and impinges the top surface. Another flow, which has little gas left with it, moves on and impinges the narrow face. A generally double roll flow pattern is formed. For the non-gas case, jet goes straight to the narrow face and forms a classic double roll flow pattern after impinging the narrow face.

#### **5.3.2 Effect of Bubble Size on Gas Entrapment**

Most gas bubbles float upward to escape from the top surface due to their strong buoyancy force. Bubbles entering lower circulation zone have a larger chance to be trapped. Although no more than 5% of the gas enter the lower circulation zone of the liquid flow [17], it is still dangerous because the volume fraction for serious gas defects is as low as  $10^{-4}$ % [24]. Thus, it is important to know how much gas is left in lower region of the mold and how large the region where the gas fraction is over  $10^{-4}$ % is.

A parameter "penetration depth" is used to tell how deep the dangerous gas fraction level can extend downward from meniscus. Penetration depth is defined as the deepest distance from the meniscus where the 10<sup>-5</sup>% gas volume fraction contour reaches. It is assumed that the deeper the penetration depth, the more difficult it is for gas to float up to meniscus and the more possible it is for gas to be entrapped. This parameter gives a clue of the chance that gas might be trapped.

Figure 5.3 shows the penetration depth of different bubbles sizes. For small bubbles (0.5mm ~ 1.0mm in diameter), the penetration depth increases with bubble size. In this size range, bubbles stay with liquid steel flow. The smaller the bubble, the stronger tendency it has to stay with the flow. Most of the flow will meet the top surface and let gas escape. With increasing bubble size, the drag forces by the bubbles decrease. Some part of the flow will not bend upward to meet the top surface. This part will go down before it meets the top surface. Gas in this part of flow will also go down the domain and has a deeper penetration. The 1.0mm bubble has the deepest penetration. Large bubbles (1.5mm~4.0mm in diameter) have a very small penetration depth. Big bubbles have a smaller total surface area and thus smaller drag force to liquid steel flow. Most bubbles leave the jet quickly after entering the domain. Very few large bubbles go down the less the domain with flow. The larger the bubble, the less gas leaves in the flow, and the less the penetration depth.

For large bubbles, the gas percentage near the narrow face is much smaller than  $10^{-4}$ % (see Figure 5.4 ). This indicates that there is very small chance for large bubbles (1.5mm~4.0mm in diameter) to be trapped and form pencil-pipe defects on the narrow face.

Another parameter that should be concerned is the area which is enclosed by the  $10^{-5}$ % gas fraction contour on the solidifying shell. It is reasonable to assume that the larger this area is, the higher the possibility of gas entrapment. This is because entrapment can only occur when the gas bubble moves very slowly relative to the solidifying interface. Figure 5.5 shows the area enclosed by the  $10^{-5}$ % gas fraction contours for different bubbles. Small bubbles (0.5mm ~ 1.0mm) have much more area enclosed by the  $10^{-5}$ % gas fraction contour than large bubbles (1.5mm~4.0mm). The 1.0mm bubbles have largest area. This indicates that small bubbles (especially around 1.0mm) are more prone to gas entrapment. Large bubbles (2.0mm~4.0mm) have very small potential entrapment areas so might not have gas entrapment.

It is assumed that gas entrapment happens when the speed of gas flowing down the shell approaches the casting speed. It is much more possible for gas to be entrapped when the gas volume fraction is over 10<sup>-5</sup>% and the gas downward velocity is close to casting speed simultaneously. Figure 5.6 shows the overlap of 10<sup>-5</sup>% gas fraction contour and gas vertical velocity contour for four different bubbles. The red-shaded area is the overlap region. These regions might be the position where gas entrapment happens.

### **5.3.3 Effect of Bubble Size on Inclusion Formation**

There are two major sources of inclusion particles: inclusion particles entering with the inlet flow and inclusion particles sheared off from flux layer and entrapped in the flow. The first kind of inclusion particles comes from steel making and erosion of refractory materials. Most of these particles circulate in the liquid pool and become entrained harmlessly into the flux layer. Some of them are entrapped in the growing shell and initiate quality problems such as surface slivers. According to experiments by R.C.Sussman etc.[4], 22.3% of particles will not float up and finally remain in the solidified shell when gas is not injected into inlet flow. With such a high percentage of entrapment of particles, inclusion defects should be very common in continuous casting. Fortunately, this is not the fact. Serious inclusion defects happen only occasionally. This indicates that steel inlet flow is rather clean and particles in inlet flow might not be the major cause of inclusion problem. Flux particles sheared away and entrapped by liquid steel flow might be the major source of inclusion particles. There are two possible mechanisms by which flux is sheared away and entrained in the steel. The first mechanism is that flux could be sheared away by liquid level fluctuations of the interface between the liquid steel and the flux layer (Figure 5.7). The second is that flux also could be sheared by downward steel flow as shown in Figure 5.8. This suggests that both the flow pattern (single vs double roll) and the velocity magnitude are important.

## 5.3.3.1. Effect of Bubble Size on Level Fluctuations

In this study, turbulence is assumed to be isotropic. The average velocity fluctuation can be obtained from the Kinetic Energy K [9]:

$$\overline{u'} = \sqrt{\frac{2}{3}K}$$

The level fluctuation height is reported to related to K also [17]:

$$h = \frac{\rho_{steel} * K}{0.5(\rho_{steel} - \rho_{flux})g}$$

Figure 5.9 shows the maximum kinetic energy and maximum level fluctuations on the top surface for different bubble sizes. Generally, small bubbles have higher level fluctuation than larger bubbles. Flow without gas has the smallest level fluctuation on the top surface. This simply shows that the stronger surface flow caused by the small bubbles leads to more level fluctuations and consequently more surface defects.

### 5.3.3.2 Effect of Bubble Size on Downward Velocity near Solidifying Shell

During operation of the continuous caster, the top liquid steel level is always moving up and down with minor fluctuations of the inlet flow rate and casting speed. This might sometimes lower the liquid steel level below the top of solidified shell sometimes. The top edge of the solidified shell then inserts into flux layer like a blade. When the flux moves to the gap, the lower part of the flux layer is peeled off and stays at the solidified shell. If liquid steel flow is upward, the peeled flux would be put back and has less chance to form inclusion particles. Otherwise, the peeled flux would be pushed down into the liquid pool and form large inclusion particles. Because these particles form just near the growing shell, they would be easily entrapped by the shell.

The maximum downward velocity can give clues on the tendency to entrap flux particles by this mechanism. Figure 5.10 shows the effect of the bubble size on vertical velocity around the top surface perimeter of the mold. The maximum downward vertical velocity appears on the wide face in all eight cases. This indicates that the wide face would likely suffer more inclusion defects than the narrow face. The inclusion particles concentrate more near the SEN on the wide face with a shallow depth. The maximum downward vertical velocity increases with increasing bubble size and reaches a maximum at 1.5 mm bubble. The maximum downward vertical velocity decreases with further increasing the bubble size. Bubbles between 1.0mm and 2.0mm have the largest maximum downward vertical velocity so might have the largest tendency to suffer inclusion defects caused by the second mechanism. Combined with the first entrapment mechanism, it is believed that smaller bubbles (d<2.0mm) have more tendency to entrap flux into liquid flow.

What should be noticed is that flow without gas has the smallest vertical velocity around top surface perimeter of the mold. Considering the flow without gas also has the smallest level fluctuations, it might have the smallest tendency to suffer inclusion entrapment. Flow without gas should have less inclusion defects than flow with gas if the liquid steel is clean. If this is the case, we will find more flux inclusions in gas injection practice than in no gas practice. This suggests that gas injection is a dangerous practice that must be carefully optimized.

## **5.4 Conclusions**

1. The jet without gas goes across the mold and impinges on the narrow face forming a classic double roll flow pattern. Jets with small gas bubbles (<1.5mm) bend upward and impinge first on the top surface to form a classic single roll flow pattern. Jets with large gas bubbles (>1.5mm) split into two branches. One, which is mostly gas, goes up and impinges on the top surface. The other, which is mostly liquid steel, impinges first on the narrow face. The flow with large bubbles (>1.5mm) shows a generally double roll flow pattern.

- 2. Medium size bubbles (around 1.0mm) appear to have the largest tendency to be entrapped deep below the meniscus as gas defects such as pencil pipes because they encourage a single roll flow pattern which imparts the largest downward velocity to the liquid jet and they are easily carried by the jet deep into the caster. With decreasing bubble size below 1 mm, the tendency of forming gas defects also decreases. Large bubbles (>1.5mm) have the least tendency to form gas defects because they float out the liquid quickly.
- 3. Smaller bubbles induce larger level fluctuations on the top surface and increase the possibility of entrapping flux particles into the liquid flow by liquid level fluctuation.
- 4. Downward velocity along the solidified shell at the meniscus reaches its maximum value at 1.5mm diameter. Bubbles with size between 1.0mm and 2.0mm have the largest possibility to cause flux to be trapped in the growing shell at meniscus due to downward movement of liquid steel and flux. The most possible position for inclusions happening by this mechanism is near the SEN at wide faces.
- 5. Compared with the flow with gas, the flow without gas has much less tendency to entrap flux particles and might have fewer inclusion problems in cases when the liquid steel is clean.
- 6. The study in this section is based on the assumption of uniform bubble size, which is not the case. The next section will explore the importance of bubble size distribution.

	Case B(6.3L/min,35"/min)
Mold Width	1854 mm
Mold Thickness	228 mm
Nozzle Submergence Depth	120 mm
(top surface to top of port)	
Nozzle Bore Inner Diameter	78 mm
Port Wall Thickness	27.5 mm
Nozzle Port Height	78 mm
Nozzle Port Width	78 <i>mm</i>
Inlet Jet Height, L <sub>h</sub>	50 mm
Inlet Jet Width, L <sub>w</sub>	78 mm
Nominal Vertical Angle of Port Edges	15° down
Inlet Jet Spread Angle	0°
Casting Speed, V <sub>c</sub>	14.8 <i>mm/s</i>
Inlet Velocity, V <sub>x</sub>	0.8766 <i>m/s</i>
Inlet Velocity, V <sub>z</sub>	0.3018 <i>m/s</i>
Inlet Turbulent Kinetic Energy, K <sub>o</sub>	$0.0502 \ m^2/s^2$
Inlet Turbulence Dissipation Rate, $\epsilon_{\!\scriptscriptstyle o}$	$0.457 \ m^2/s^3$
Liquid Steel Density, p	7020 $kg/m^3$
Steel Laminar (Molecular) Viscosity, $\mu_{o}$	0.00560 kg/m s
Inlet Gas Flow Rate(whole slab)	6.3 <i>l/min</i>
Inlet Gas Volume Fraction, $f_{gas}$	8.5% or 0%(No gas)
Average Gas Bubble Diameter, $D_0$	0.5, 0.7, 1.0, 1.5, 2.0, 3.0, 4.0 mm
Gravitational Acceleration, g	9.8 $m/s^2$

Table 5.1. Parameters used in the simulations



Figure 5.1 Schematic of simulation domain with boundary conditions and typical meshes





0.7 mm bubble

2.5m/s

# 1.0 mm bubble

Figure 5.2 Effects of bubble size on fluid flow pattern in continuous casting mold



3.0 mm bubble

4.0 mm bubble

Figure 5.2 Effects of bubble size on fluid flow pattern in continuous casting mold(continued)



Figure 5.3 Effect of bubble size on gas penetration depth



Figure 5.4 Gas fraction contour for different sized bubbles

150mm (Scale for top three plots)



150mm (Scale for following four plots)



Figure 5.5 Gas percentage contours for different bubbles



1.0 mm bobble

1.5 mm bubble





Figure 5.7 Schematic of the first mechanism of flux entrapment



Figure 5.8 Schematic of the second mechanism of flux entrapment



Figure 5.9 Effect of bubble size on the maximum kinetic energy and level fluctuation on the top surface



Figure 5.10 Effect of bubble size on vertical velocity around top surface perimeter

# Section 6

# Effect of Argon Flow Rate on Flow Pattern in Continuous Slab Casting

## 6.1 Introduction

The influence of argon on flow in continuous slab casting varies with factors such as the gas volume fraction in liquid steel flow, gas distribution in the jet, bubble size and jet direction etc. Increasing gas volume fraction gives more lift force to the liquid jet and tends to cause single roll. A jet with most gas concentrating in upper part of the jet tends to cause double roll due to less interaction with jet. Larger bubbles leave jet more quickly and tends to cause double roll. Those factors are influenced by liquid flow rate, argon flow rate, argon injection method and nozzle geometry. The mold flow is also influenced by submergence depth, mold geometry and nozzle orientation.

Flow in the nozzle and mold is difficult to measure due to the high temperature. Computational simulation is therefore an efficient way to know the flow behavior in the nozzle and mold. In this section, typical argon flow rates, casting speeds and slab widths used in practice are investigated. Details of the nozzle are given in Table 6.1. The computational simulation takes two steps. In the first step, flow in the nozzle is simulated. In the second step, flow in the mold is simulated using inlet conditions from the results recorded at the nozzle port in the first step. The inlet mesh is chosen exactly the same as the nozzle port in the first step. This integrates the nozzle and mold computational simulation into a complete computational simulation set. The validity of this procedure was demonstrated by Hershey [28]. In Section 5, it was shown that bubble size has a great influence on the flow pattern. Bubble size and its distribution in the nozzle varies greatly due to different bubble formation mechanisms which depend on liquid flow velocity and gas injection methods [16]. Gas bubbles will also breakup and coalesce after they enter the nozzle and mold. Treating gas bubbles as uniform-sized is too rough an approximation and is not able to capture the influence of bubbles on the liquid flow in detail.

In this section, the MUSIG model (described in Section 2) is used to model breakup and coalescence of bubbles. The initial bubble size and its distribution in the nozzle are measured from a two-needle water model experiment. Because the bubble size under the same casting conditions differs in water-air system and steel-argon system, Bai's model [16] is used to extrapolate the results from water-air system to the steelargon system. For the 0.4-scale water model, the bubble size is measured from still photographs. The objective of this study is to validate the multiphase model with MUSIG model by the comparison of computed results with the water model measurements. The model is then used to identify the flow pattern in the real caster for different casting conditions.

## 6.2 Definition of the Mean Bubble Size in This Study

In this section, the average bubble diameter is calculated by the Average Volume Method. First, the total volume of all bubbles in the sample pool is calculated. Then, the average volume of a single bubble is calculated by dividing the total volume by the total number of bubbles. Finally, the average diameter is calculated by assuming spherical bubbles having the same volume as the average volume. The equation to calculate the average bubble diameter is:

$$D_{avg} = 2 * \sqrt[3]{3 * (\sum_{i=1}^{n} (Ni * V_i) / \sum_{i=1}^{n} N_i) / 4 / \pi}$$

where  $V_I$  is the volume of a single bubble at size i.  $N_I$  is the number of bubbles at size I in a sample pool.

The volume fraction of bubbles at size i is straightforward:

$$VF_i = N_i * V_i / \sum_{i=1}^n N_i * V_i$$

where  $VF_I$  is the volume fraction of bubbles at size i.

## 6.3 Bubble Size and Its Distribution in the 0.4-scale Water Model

In this study, instantaneous photographs are used to measure the bubble size and its distribution in the water model for the conditions in Table 6.4 and 6.5. Two such photographs are shown in Figs. 6.1 and 6.2. The air injection method of the 0.4-scale water model at LTV steel is different from the steel casting nozzle. Instead of a porous ceramic filter, the 0.4-scale water model uses a slot ring to inject air. This method produces very large bubbles at low casting speed (> 10 mm). The results are shown in Table 6.2 and 6.3. Figure 6.3 shows the bubble size and its distribution in the water model. With slot ring injection, bubble size is relatively uniform at high casting speed (23.2 mm/s). Most bubbles have the size between 2.5 mm and 3.5 mm. Average bubble size is 2.59 mm. The bubble size is widely distributed at low casting speed (14.8 mm/s). Although few in numbers, huge bubbles (9.5 mm and 10.5mm) account for about 47% of total gas volume. Most bubbles are smaller than 5.0 mm. The number of bubbles smaller

than 1.0 mm is larger than that at high casting speed (case A). This causes the average bubble size (2.43 mm) to be smaller than that at high casting speed (2.59 mm).

### 6.4 Validation of MUSIG Model by 0.4-scale Water Model

The MUSIG model was validated by comparison with experiments in the 0.4scale water model. Two water model cases (case A and B) were performed. These two cases are scaled from two typical plant cases at LTV Steel. Case A corresponds to 1.854 m slab with casting speed of 23.2 mm/s and 13SLPM argon flow rate. Case B corresponds to 1.854 m slab with casting speed of 14.8 mm/s and 6.3 SLPM argon flow rate. The conditions for these two water model cases are listed in Table 6.4 and 6.5. In the practice, case A tends to have more pencil pipe defects and case B tends to have more sliver defects. The computational simulation domain for the 0.4-scale water model is the same as that in section 3 (Figure 3.1). The bubble size and its distribution were obtained by still photographs as discussed in section 6.3. Figure 6.4 compares the still photograph, the result of computational simulation and that of PIV measurement on the case A. On the left is a still photograph of the flow for case A. In the middle is the computational simulation result of velocity on the center plane of the 0.4-scale water model. On the right is the PIV measurement of velocity on the center plane of the 0.4-scale water model. Figure 6.6 shows a similar comparison but for case B. From these comparisons, it is obvious that the computational simulation matches the observations and measurements qualitatively for both cases. Both case A and case B have a single roll flow under normal casting conditions.

An interesting phenomenon was observed in the experiments. The flow rate control was not so smooth so that flow rate changes sometimes occurred. When the water level is low, the flow rate was increased by up to 15% to restore the standard water level. In the steel continuous casting, the same phenomenon might exist when ladle is changed (As a practice in LTV steel, casting speed is decreased from 23.2 mm/s to 14.8 mm/s and then increased back to 23.2 mm/s with no gas flow rate changes when ladle is changed). The 15% increase of flow rate in the 0.4-scale water model changes the flow pattern. Figure 6.5 and Figure 6.7 show still photographs of the flow during the 15% increase in flow rate for the conditions the same as case A and B respectively. For case A, the flow changes to a slight double roll from the single roll. However, for case B, the flow is always single roll, even for the 15% increase in the flow rate. The computational simulation matches the measurements in both cases. Although the observations, measurements and computational simulations in the 0.4-scale water model match each other, they don't match the plant observations and measurements in steel continuous casting [26]. Further discussion is given in section 6.8.5.

## 6.5 Bubble Size and Its Distribution in the Steel Casting Nozzle

In the steel casting nozzle, the initial bubble size depends mainly on the vertical liquid velocity, the gas flow rates per orifice and the contact angle [16]. The steel throughput, nozzle bore diameter, argon flow rate and ceramic filter structure are the main factors that determine the bubble size and its distribution. Under different casting conditions, the bubble size and its distribution in the nozzle are different. The nozzle geometry parameters and casting conditions for both case A and B are listed in Table 6.1.

The average liquid velocity is calculated from liquid flow rate:

$$V = \frac{Q_{liq}}{\pi R^2}$$

where R is the radius of the nozzle. The liquid flow rate  $Q_{liq}$  is calculated according to:

$$Q_{liq} = V_c * W * T$$

where  $V_c$  is casting speed, W is mold width, T is mold thickness.

The vertical liquid velocity shearing gas at pores as it across the ceramic pores in the nozzle is 2.05 m/s and 1.31 m/s for case A and B respectively. The number of active pores in the ceramic filter is estimated based on Zhang's experiments [25] which suggests 60 mm<sup>2</sup> of surface area per active site. For the nozzle in case A and B, the area of the porous ceramic filter is 50 mm high which gives total inner surface area of ceramic filter of 12246 mm<sup>2</sup>. The total active pores in the ceramic filter is then 204. The gas flow rate per pore is about 2.9 ml/s and 1.8 ml/s .

The bubble size for the above parameters is based on measurements in the twoneedle water model experiments by Hua Bai [16]. The water model is composed of a 35×35 cm square section nozzle and two needles (Figure 6.8). The nozzle wall is made of transparent plastic for easy observation. Needles with 0.2, 0.3 and 0.4 mm in diameter were used. Water flow through the nozzle and air is injected into the nozzle through two needles. Adjusting the water flow velocity and air injection rate obtains different bubble sizes. The bubble size is measured from high speed video camera frames. Bai's experiments don't have the exact case A and B. This study chooses the closest cases to case A and B. Table 6.6 and 6.7 show the results. For the high liquid flow rate case (case A), bubbles distribute over a relatively wide size range. Small bubbles are much more numerous than larger bubbles. But large bubbles (>3.0 mm) take more than 76% volume. The mean bubble size for case A is 1.94 mm. For the low liquid flow rate case (case B), bubbles concentrate around the mean bubble size. Bubbles show a relatively uniform size. The mean bubble size for case B is 2.12 mm. A statistical analysis of the bubble size distribution for the two cases is shown in Figure 6.9. The above bubble sizes and their distribution are used in this study.

The bubble in the real caster might be slightly larger for two reasons. Firstly, the mean bubble size in the steel argon system is larger than the corresponding bubble size in the water-air system according to Bai's bubble formation model (Figure 6.10) [16]. For case A (vertical liquid velocity = 1.9 m/s and gas flow rate per pore = 2.9 ml/s), the steel-argon system has a mean bubble size of 2.7 mm while water-air system has a mean bubble size of 2.6 mm. For case B (vertical liquid velocity = 1.4 m/s and gas flow rate per pore = 1.8 ml/s), the steel-argon system has a mean bubble size of 2.75 mm while water-air system has a mean bubble size of 2.38 mm. Secondly, Bai measured a larger mean bubble diameter than in the nozzle computational simulation in this study.

### 6.6 Computational simulation of Flow in the Nozzle

Before every mold computational simulation, a nozzle computational simulation is run for the particular throughput and gas injection rate. Figure 6.11 shows the geometry of the nozzle. Liquid flows in from the top surface where the nozzle connects to the tundish. Argon is injected into nozzle through porous ceramic in the gas injection zone. The top surface and gas injection zone are set as inlet boundaries. The nozzle wall is set as non-slip wall boundary. The nozzle port is set as pressure boundary. Linear pressure is applied at the nozzle port based on submergence depth. The geometry and other parameters for nozzle computational simulation are listed in Table 6.7. About 2 CPU hours on NCSA SGI supercomputer are needed for 500 iterations which achieve a fully convergence.

Some typical results for case A and B are shown in Figure 6.12 and 6.13. There is a swirl on the center plane parallel to SEN port in case A while flow on that plane in case B is symmetric. The inflow area is larger in case B than that in case A. The spread angle of the jet in case A is also larger than that in case B.

### 6.7 Computational simulation of Flow in the Continuous Caster

The inlet condition of the mold computational simulation is from the nozzle computational simulation output. Variables are passed by a data file for initial conditions at inlet (see Appendix E and F ). Figure 6.14 and 6.15 shows that these variables are passed consistently. Because the slide gate orientation is 90°, the flow at the two ports is assumed to be identical. Thus, flow in the caster is assumed to be symmetric on left and right sides. Only half of the mold is simulated. Figure 6.16 shows the geometry of mold domain. The top surface is set as a non-slip wall boundary considering the slow movement of the high viscosity slag. The wide face and narrow face are set as non-slip wall as well. The outlet is set as a pressure boundary. Other parameters are shown in Table 6.8. The standard K- $\varepsilon$  model is used to simulate multiphase flow in the mold and the MUSIG model is used to model bubble breakup and coalescence. About 13 CPU hours are needed for 1000 iterations on NCSA SGI supercomputer which gives a fully convergence (decreasing final residuals below 0.1% of original residual).

## 6.8 Results and Analysis

### 6.8.1 Influence of Casting speed on Flow in the Continuous Caster

Casting speed has a great influence on flow in the nozzle. High casting speed needs a high liquid velocity and gate opening in nozzle. High liquid velocity tends to produce unbalanced flow and a large jet horizontal angle at nozzle port. For case A (high casting speed), the horizontal angle of the jet is 7°. For case B (low casting speed), the horizontal angle of the jet is 3°. The high horizontal angle directs more liquid toward the wide face and makes the flow more complex at high casting speed. Low casting speed makes liquid velocity in the nozzle small. Small liquid velocity produces larger bubbles according to Bai's model [16]. Large bubbles leave the jet more quickly after it enters the mold and imparts less buoyancy to the jet. The jet is bent up less and has more tendency to form a double roll flow pattern as explained in section 5. Figure 6.17 and 6.18 show that high casting speed has a complex flow pattern at the center plane while lower casting speed has a barely double roll flow pattern. The horizontal flow might help to prevent changing from single roll to double roll or vice versa by causing a relatively wide complex flow range. With gas percentage varying in a small range (11% for Figure 6.17 and 8.9% for Figure 6.18), increasing casting speed has no great influence to the flow. With gas flow rate constant (13 SLPM for both Figure 6.17 and Figure 6.19), increasing casting speed from 14.8 mm/s (Figure 6.19) to 23.2 mm/s (Figure 6.17) changes the flow from single roll to complex flow. Therefore, gas fraction is a much better indicator to identify the flow pattern in the mold than casting speed and gas flow rate.

### 6.8.2 Influence of Gas Fraction on the Flow Pattern.

Gas fraction is a good indicator to identify the flow pattern in the mold. High gas fraction increases buoyancy applied to the jet and bends up the jet more to produce a single roll flow pattern. Low gas fraction has more tendency to form a double roll. With increasing gas fraction, the flow pattern will evolve through three different flow patterns: double roll, complex and single roll. Changing the flow patterns brings liquid level fluctuation at meniscus and is very detrimental. Gas fraction should be controlled to avoid flow pattern switching. Plant measurements show that there is a critical gas fraction for flow switching. Below the gas fraction, the flow is double roll. Above the gas fraction, the flow is single roll. The critical gas fraction does not change with the casting speed or throughput. Figure 6.24 shows the computational simulation results of the relationship between the gas fraction and the flow patterns for four different slab widths. It gives the critical gas fractions where double roll changes to single roll for different slab widths. The critical gas fraction at which the flow switch from double roll to single roll or vice versa is not simply a point. There is a range of gas fraction which is called "complex zone" here for the flow pattern transition. Above this zone, the flow pattern is single roll. Below it, the flow pattern is double roll. In the complex zone, the flow pattern is complex.

Even the complex flow is relatively stable and is expected to have no harmful effect to slab quality such as slivers (case A). However, a flow pattern change which increases liquid level fluctuations is suspected to be more harmful in this aspect (case B) [26]. Plant measurements from LTV steel [26] show that, during ladle changing, the casting speed dropped from 23.2 mm/s (case A) to 14.8 mm/s (case B) without

decreasing the gas flow rate. The computational simulations show that this process experiences a flow pattern change from complex flow to single roll due to the increase in gas fraction from 11% to 16.4% [Figure 6.24]. Production records show that slabs produced in this period had more sliver problems [26]. The flow pattern switching happening in this period could explain the sliver defects. Although a stable complex flow pattern should have less tendency to cause high liquid level fluctuation and entrap flux, its complex feature is still harmful. Gas is less likely to escape the upper recirculation zone and is eventually entrapped in the slab to form pencil pipes. Case A (11% gas) experienced more pencil pipe defects than case B even though case B has a much higher gas fraction (16.4% gas). Gas fraction should be kept stable and away from complex zone to give a stable flow pattern. Figure 6.24 and Figure 6.25 show the relationship between gas fraction and flow pattern.

### 6.8.3 Influence of Slab Width on the Flow Pattern

Decreasing slab width decreases the distance between the nozzle port and the narrow face. The jet has less time to be lifted up to hit the top surface. Thus, the narrower slab is more likely to form a double roll flow pattern than a wider slab. Figure 6.25 shows this tendency. With a gas flow rate of 15 SLPM and throughput of 2 ton/min, 1.854 m and 1.600 m slab show single roll, 1.321 m slab shows complex flow and 1.016 m slab shows double roll. With a casting speed of 16.9 mm/s and gas percentage of 15%, 1.854 m slab shows a single roll, 1.600 m slab shows a complex flow, and 1.321 m and 1.016 m slab shows a single roll, 1.600 m slab shows a complex flow, and 1.321 m and 1.016 m slab shows a single roll, 1.600 m slab shows a complex flow, and 1.321 m and 1.016 m slab show double roll (Figure 6.24). It is concluded that, keeping casting speed and gas fraction constant, decreasing slab width is likely to have double roll flow in the mold.
## 6.8.4 Influence of Submergence Depth on the Flow Pattern

Increasing submergence depth increases the vertical distance from the jet to the top surface. It is more difficult for the bent jet to hit the top surface for deep submergence depth. Thus, deeper submergence depth has more tendency to form a double roll. Figure 6.26 shows how the flow pattern changes with submergence depth for a given slab width of 1.854 m. Figure 6.27 and Figure 6.28 show velocity vector plots at center plane which show the relationship between SEN submergence depth and flow patterns. For a given slab width, submergence depth and gas fraction, throughput has little effect on the flow pattern.

## 6.8.5 Difference between flow in water model and steel caster

In Figure 6.4 through Figure 6.7, PIV measurements in water model verified the computational simulations. The model then is applied to the corresponding steel cases. The results of steel caster computational simulation are different from that of the computational simulation of water models. For case A, the steel caster computational simulation shows a complex flow which is more closer to a double roll than a single roll. The water model computational simulation and PIV measurements show that its water model counterpart is normally single roll. For case B, the steel caster computational simulation and PIV measurements show that its water model counterpart is normally single roll. For case B, the steel caster computational simulation and PIV measurements show that its water model counterpart is always single roll. Plant measurements [26] showed that case A is normally double roll and case B is mostly double roll but experiences some flow pattern switching. The plant measurements also show that, keeping gas flow rate constant and decreasing casting speed from 23.2 mm/s

to 14.8 mm/s, the flow pattern changes from double roll to single roll. This phenomenon is matched by the computational simulations. Keeping gas flow rate at 13 SLPM and changing casting speed from 23.2 mm/s (Figure 6.17) to 14.8 mm/s (Figure 6.19), the flow pattern changes from double roll to single roll. The plant measurements are consistent with computational simulations of flow in the steel caster. Although both plant measurements in steel caster and 0.4-scale water model experiments match their computational simulations, they are not consistent with each other. Computational simulation and measurements show that the flow pattern in the steel caster is sometimes very different from that in a scale water model and the steady, multiphase K- $\varepsilon$ computation can match both. The most important cause for the difference is the reducedscale of the water model combined with the Froude-based velocity scaling criterion used to choose the water model flow rates.

## 6.9 Conclusions

- Multiphase computational simulations of a 0.4-scale water model using the MUSIG model match experimental measurements. Plant measurements of the flow pattern by MFC sensors also verify multiphase computational simulation of the real caster. The 0.4-scale water model does not always match flow in the real caster and in some cases predicts the wrong flow pattern (e.g. 1.854 m slab, casting speed of 14.8 mm/s and gas fraction of 8.5%).
- 2. Transient increases in the liquid flow rate, such as required to accommodate level changes, may produce significant changes in the flow pattern, if the conditions are near critical.

- 3. Gas changes the flow pattern by applying buoyancy force to the liquid jet. Increasing gas fraction changes the flow from double roll to an intermediate state--complex flow, then to single roll. For a given slab width, there is a critical gas fraction range where the flow pattern is in transition from double roll to single roll. This range does not change too much with throughput for a given slab width and submergence depth. The critical gas fraction increases when decreasing slab width.
- 4. Slab width has influence on flow patterns. For a given casting speed and gas fraction, narrow slab tends to form double roll while wide slab tends to form single roll.
- 5. Nozzle submergence depth also affects the flow pattern in the mold. Deeper submergence depth has more tendency to form double roll.

Case A Nozzie	Case B Nozzle		
100			
310			
45			
70			
66			
776			
78			
165			
78X78			
28.5			
15°			
12			
90°			
58.6%	52%		
1.40	0.89		
590.5	375.8		
0.927			
13	6.3		
11%	8.9%		
	100   310   45   70   66   776   78   165   78X78   28.5   15°   12   90°   58.6%   1.40   590.5   0.927   13   11%		

Table 6.1 Nozzle parameters (steel caster computational simulation)

\*Blank in second column is the same as the first column.

Region	≤1mm	1-2mm	2-3mm	3-4mm	4-5mm	≥5mm		
А	>10	16	20	17	2	0		
В	>5	8	16	10				
С	>10	9	13	5				
Total	25	33	49	32	2			
Volume (ml)	13.08	58.29	400.68	718.01	95.38			
Volume percentage	1.02%	4.53%	31.15%	55.83%	7.42%			
Average	2.59mm							

Table 6.2 Measured bubble size distribution in LTV 0.4-scale water model for case A (23.2 mm/s (55ipm) + 13 SLPM/11% hot gas)

Table 6.3 Measured bubble size distribution in LTV 0.4-scale water model for case B (14.8 mm/s (35ipm) + 6.3 SLPM/8.5% hot gas)

Region	≤1mm	1-2mm	2-3mm	3-4mm	4-5mm	5-6mm	6-7mm	7-8mm	8-9mm	9-10mm	≥10mm
11	5	2	1								
12	>11	3	1	1							
13	>14	6	4	2	1						
14	>16	4	2	1	1						1
21	9	3									
22	>13	4	1		1						
23	>20	3	2	1							
24	>20	4	2	2		2				1	
31	>10	4									
32	>10	3	2								
33	>17	4	3	1		1					
34	5	6	2		1						
35	6	1	3		1						
44	6	4	1								
45	>12	5	2								
Total	>174	57	26	8	5	3				1	1
Volume (ml)	91.05	100.68	212.60	179.50	238.44	261.20				448.69	523.69
Volume percent	4.43%	4.90%	10.34 %	8.73%	11.60 %	12.71 %				21.83%	25.48%
Averag	ge size	2.43mm									

	Normal Conditions
Mold Width W (mm) x Thickness H (mm)	730 x 80
Mold Height (mm)	950
Nozzle Submergence Depth	80
(top surface to top of port)	
Nozzle Inner Diameter (mm)	31
Nozzle Port Width (mm) x Height (mm)	31 x 31
Jet Angle	30° down
Inlet Jet Spread Angle	0°
Water Flow Rate Q <sub>w</sub> (SLPM)	58.59 (15.5 GPM)
Equivalent Steel Casting Speed (mm/s)	22.9
$V_c = \frac{Q_w}{0.4 \times W \times 0.4 \times H \times \sqrt{0.4}}$	
Gas Flow Rate (SLPM, hot volume)	7.43 (15.8 SCFH)
Gas Volume Fraction (%)	11.3
Inlet Velocity, $V_x$ (m/s)	0.571
Inlet Velocity, $V_z(m/s)$	0.330
Inlet Turbulent Kinetic Energy, $K_o(m^2/s^2)$	0.044
Inlet Turbulence Dissipation Rate, $\varepsilon_0$ (m <sup>2</sup> /s <sup>3</sup> )	0.999
Water Density (kg/m <sup>3</sup> )	1000
Water Viscosity (m <sup>2</sup> /s)	1×10 <sup>-3</sup>
Gas Density (kg/m <sup>3</sup> )	1.20
Gas Viscosity (m <sup>2</sup> /s)	1.7×10 <sup>-5</sup>
Average Bubble Diameter (mm)	2.590
Breakup Coefficient	0.5
Coalescence Coefficient	0

Table 6.4 Parameters for scale water model for case A

	Normal Conditions
Mold Width W (mm) x Thickness H (mm)	730 x 80
Mold Height (mm)	950
Nozzle Submergence Depth	80
(top surface to top of the port)	
Nozzle Inner Diameter (mm)	31
Nozzle Port Width (mm) x Height (mm)	31 x 31
Jet Angle	30° down
Inlet Jet Spread Angle	0°
Water Flow Rate (SLPM)	37.80 (10.0 GPM)
Equivalent Steel Casting Speed (mm/s)	14.8
$V_c = \frac{Q_w}{0.4 \times W \times 0.4 \times H \times \sqrt{0.4}}$	
Gas Flow Rate (SLPM, hot volume)	3.71 (7.9 SCFH)
Gas Volume Fraction (%)	8.9
Inlet Velocity, $V_x$ (m/s)	0.358
Inlet Velocity, $V_z(m/s)$	0.207
Inlet Turbulent Kinetic Energy (m <sup>2</sup> /s <sup>2</sup> )	0.044
Inlet Turbulence Dissipation Rate (m <sup>2</sup> /s <sup>3</sup> )	0.999
Water Density (kg/m <sup>3</sup> )	1000
Water Viscosity (m <sup>2</sup> /s)	1×10 <sup>-3</sup>
Gas Density (kg/m <sup>3</sup> )	1.20
Gas Viscosity (m <sup>2</sup> /s)	1.7×10 <sup>-5</sup>
Average Bubble Diameter (mm)	2.43
Breakup Coefficient	0.1
Coalescence Coefficient	0

Table 6.5 Parameters for scale water model for case B

Diameter	<1.0	1.0	1.5	2.0	2.5	3.0	3.0-	>5.0	Photo
of bubbles(mm)							5.0		Frame No.
Data Set #1	4	3	1	1	0	1	1	0	159
Data Set #2	2	0	0	0	0	0	1	0	219
Data Set #3	0	0	0	1	1	0	1	0	265
Data Set #4	2	1	1	0	1	0	1	0	285
Data Set #5	2	2	1	0	1	0	0	0	315
Data Set #6	3	2	0	0	0	1	2	0	368
Data Set #7	0	1	1	0	0	0	2	0	397
Data Set #8	3	2	0	1	0	0	1	0	532
Data Set #9	2	1	1	0	1	0	2	0	631
Data Set #10	1	2	0	1	0	0	2	0	668
Total Number	19	14	5	4	4	2	13	0	
Total Volume	3.73	11.00	8.84	12.57	19.63	14.14	163.3	0	
$(mm^3)$									
Volume %	1.6	4.72	3.79	5.39	8.42	6.06	70	0	
Average	1.94 mm								

Table 6.6 Measured bubble size distribution for case A in double-needle experiments

Table 6.7 Measured bubble size distribution for case B in double-needle experiments

Diameter	<1.0	1.0	1.5	2.0	2.5	3.0	>3.0	Photo
of bubbles(mm)								Frame No.
Data Set #1	0	0	2	5	4	2	0	141
Data Set #2	0	0	3	4	2	3	0	177
Data Set #3	0	0	1	7	3	0	0	211
Data Set #4	0	0	2	7	3	1	0	246
Data Set #5	0	1	6	4	4	0	0	286
Data Set #6	0	0	2	3	6	2	0	341
Data Set #7	0	1	3	7	3	0	0	380
Data Set #8	0	0	2	8	4	2	0	444
Data Set #9	0	0	1	6	5	0	0	481
Data Set #10	0	0	2	5	4	2	0	522
Total Number	0	2	24	56	38	12	0	
Total Volume	0	2.09	56.55	234.57	248.71	113.1	0	
$(mm^3)$								
Volume %	0	0.32	8.63	35.81	37.97	17.27	0	
Average	2.12 mm							

	Case A	Case B
Mold Width (mm)	1854	
Mold Thickness (mm)	228	
Nozzle Submergence Depth	165	
(top surface to top of port) (mm)		
Nozzle Bore Inner Diameter (mm)	78	
Nozzle Port Height (mm)	78	
Nozzle Port Width (mm)	78	
Vertical Velocity in Nozzle (m/s)	2.05	1.31
Nominal Vertical Angle of Port Edges	15° down	
Inlet Jet Spread Angle	0°	
Casting Speed, V <sub>c</sub> (mm/s)	23.2	14.8
Liquid Steel Density, $\rho_1$ (kg/m <sup>3</sup> )	7020	
Gas Density, $\rho_{gas}$ (kg/m <sup>3</sup> )	0.27	
Steel Laminar (Molecular) Viscosity, $\mu_o$ (kg/m/s)	0.00560	
Gas Viscosity, $\mu_{gas}$	7.42E-5	
Surface Tension Coeff. (Steel-Argon) (N/m)	1.192	
Inlet steel flow rate (m <sup>3</sup> /min)	0.584	0.376
Throughput (ton/min)	4.10	2.64
Inlet Gas Flow Rate (SLPM)	13	6.3 / 13
Inlet Gas Volume Fraction, $f_{gas}$	11%	8.5%/16.4%
Gravitational Acceleration, g (m/s <sup>2</sup> )	9.8	
Breakup Coefficient	0.5	0.1
Coalescence Coefficient	0	0

Table 6.8 Parameters in the steel caster computational simulation

\*Blank in second column is the same as the first column.



Figure 6.1 Bubble size in the water model (case A, 23.2 mm/s + 13 SLPM/ 11% hot gas)



Figure 6.2 Bubble size in the water model (case B, 14.8 mm/s + 6.3 SLPM/ 8.5% hot gas)



Figure 6.3 Bubble size distribution in the mold (measurements in 0.4 scaled water model)



0.4m/s



Figure 6.4 Comparison of velocity at centerplane between PIV measurements, computational simulation and eyeview for case A



Flow Picture of Water Model (55ipm +11% gas)



Liquid Velocity Vectors of Modeling (55ipm+11%gas)

Figure 6.5 Comparison of computational simulation and eyeviews while adjusting liquid level with 15% increase in liquid flow rate for case A



Figure 6.6 Comparison of velocity at centerplan between PIV measurements, computational and eyeview for case B





Liquid Velocity Vectors of Modeling (35ipm+8.5%gas)

Figure 6.7 Comparison of computational simulation and eyeviews while adjusting liquid level with 15% increase in liquid flow rate for case B



Figure 6.8 Schematic of water model of bubble formation



Figure 6.9 Assumed bubble size distribution in the steel caster nozzle (based on double-needle experimental measurements)



Figure 6.10 Average diameter of bubbles vs gas flow rate (Rearranged from Figure 18 in [27]),



Figure 6.11 Geometry of the nozzle



Figure 6.12 Liquid steel velocity in the nozzle for case A



Centerplane parallel to SEN port

Centerplane perpendicular to SEN port

Figure 6.13 Liquid steel velocity in the nozzle for case B



Nozzle port output by nozzle modeling (Velocity vector U)



Real caster nozzle port input (Velocity vector U)



Nozzle port output by nozzle modeling (Velocity vector V)

>2.0000 E-01 1.3846E-01 6.9231E-02 0.0000E+00-6.9231E-02 -1.3846E-01 <-2.0000E-01

>1.2430E+00 8.7806E-01 5.1306E-01 1.4805E-01 -2.1696E-01 -5.8197E-01 -9.4697E-01

(same key for both plots)

>5.0000E-01 3.1538E-01 1.0769E-01 -1.0000E-01 -3.0769E-01 -5.1538E-01 <-7.0000E-01



Real caster nozzle port input (Velocity vector V)



Nozzle port output by nozzle modeling (Velocity vector W)



Real caster nozzle port input (Velocity vector W)

## Figure 6.14 Correspondence of nozzle computational simulation output to mold computational simulation input

117





Nozzle port output by nozzle modeling (Volume fraction of liquid)

Real caster nozzle port input (Volume fraction of liquid)



(same key for all four plots)



Nozzle port output by nozzle modeling (Volume fraction of gas)



Real caster nozzle port input (Volume fraction of gas)





Figure 6.16 Computational simulation domain for mold with typical meshes



Figure 6.17 Computed steel flow pattern with distributed bubble size for case A (wide face)



Figure 6.18 Computed steel flow pattern with distributed bubble size for case B (wide face)



Figure 6.19 Computed steel flow pattern with distributed bubble size for case B with high gas fraction (wide face, 16.4% gas)



Figure 6.20 Computed steel flow pattern with distributed bubble size for case A (narrow face)



Figure 6.21 Computed steel flow pattern with distributed bubble size for case B (narrow face)



Figure 6.22 Computed steel flow pattern with distributed bubble size for case B with high gas fraction (narrow face, 16.4% gas)



Figure 6.23 Flow pattern identification for 0.4-scale water model (modified from M. Assar, P. Dauby and G. Lawson [26])



Figure 6.24 Flow pattern identification for steel caster (from computational simulation of steel caster)



Figure 6.25 Flow pattern identification (estimated from computational simulations of steel caster)



Figure 6.26 Relationship between throughput and submergence depth (1.854 m slab width)



Figure 6.27 Computed velocity at centerplane for different SEN submergence depth (1.854 m slab, 23.2 mm/s castingspeed and 11% hot gas)


submergence depth 190mm (7.5 inch)

submergence depth 165 mm (6.5 inch)

Figure 6.28 Computed velocity at centerplane for different SEN submergence depth (1.854 m slab, 14.8 mm/s castingspeed and 8.5% hot gas)

# Section 7

# Conclusions

In this study, computational models of turbulent flow in continuous casting of steel have been developed, including single phase flow, single phase heat transfer, multiphase flow with uniform size and multi-size bubbles. Experiments are performed on a 0.4-scale water model to verify the single phase and multiphase turbulent flow models by comparing their prediction with the velocity measurements using PIV (Particle Image Velocity) technology. The 0.4-scale water model measurements and computational results for multiphase with multi-size bubbles are then compared with plant measurements using MFC (Mold Flow Control) sensors to verify the 0.4-scale water model and the multiphase model. The validated models are then applied in a parametric study to investigate the effects of casting conditions on the defects such as pencil pipes and slivers. Plant measurements using thermocouples in the steel caster are performed to verify the single phase heat transfer model. The validated model was applied to predict the heat flux in the steel caster. The heat flux is used in the calculation of the solidification of the shell using CON1D program.

#### 7.1 Validation of Fluid Flow Models in Continuous Casting

The velocities predicted by both the single phase and multiphase K- $\varepsilon$  model developed in this study match PIV measurements in the 0.4-scale water model both qualitatively and quantitatively. The comparison between the simulations of 0.4-scale water model, full size water model and steel caster showed that the flows in these cases are similar with each other and the single phase K- $\varepsilon$  model developed in this study can

be used to model a full size continuous caster with reasonable accuracy. Comparison between the simulations of the proportionally scaled and Froude scaled water model showed that the difference between these two scaling methods is negligible for single phase flow.

The multiphase flow model with the MUSIG model developed in this study matches the PIV measurements in the 0.4-scale water model. Different simulations using the same model also match plant measurements of the flow pattern using MFC sensors in the steel caster. However, the 0.4-scale water model does not always match flow in the real caster, due to the effects of gas injection, and in some cases produced very different flow patterns. This indicates that the 0.4-scale water model can not reliably reproduce multiphase flow in the steel caster at least when using Froude similarity.

#### 7.2 Validation of Heat Transfer Model

The single phase heat transfer model of temperature in the molten pool developed in this study was verified by plant measurements using thermocouples in the steel caster. The thermocouple measurements show that the current single phase heat transfer model can predict temperature in the upper part of the mold well (with an error less than 7%). The heat transfer model developed in this study can be used to simulate the temperature in the liquid pool of the steel caster and the heat flux through the solidified shell.

#### 7.3 Effects of Gas Injection on Fluid Flow in the Steel Caster

The effects of gas bubble size on the flow pattern in the steel caster are first investigated based on the oversimplified assumption of uniform size bubbles. Then, The effects of multi-size bubbles on the liquid flow were investigated with the MUSIG model (multiphase flow with multi-size bubbles) developed in this study. Specific findings include:

- Gas changes the flow pattern by applying buoyancy force to the liquid jet. Increasing gas fraction changes the flow from double roll to an intermediate state--complex flow, then to single roll. For a given slab width, there is a critical gas fraction range where the flow pattern is in transition from double roll to single roll. This range does not change too much with steel throughput for a given slab width and submergence depth.
- The critical gas fraction range decreases with smaller bubbles. For a typical case, the jet without gas goes across the mold and impinges on the narrow face forming a classic double roll flow pattern. With 8.5% small uniform gas bubbles (<1.5mm), the jets bend upward and impinge first on the top surface to form a classic single roll flow pattern. With 8.5% large gas bubbles (>1.5mm), the jets split into two branches. One, which is mostly gas, goes up and impinges on the top surface. The other, which is mostly liquid steel, impinges first on the narrow face. The flow with 8.5% large bubbles (>1.5mm) shows a generally double roll flow pattern.
- For the same gas fraction, small bubbles appear to have more tendency to cause a single roll flow pattern. Increasing a uniform bubble size above1.5 mm, flow pattern changes from single roll to double roll. The tendency to form gas defects is also decreased because large bubbles float out of the liquid quickly.

- The Single roll flow pattern such as induced by smaller bubbles induces larger level fluctuations on the top surface and increases the possibility of entrapping flux particles into the liquid flow.
- The downward velocity along the solidified shell at the meniscus reaches its maximum value for the strong single roll flow pattern which accompanied 1.5 mmdiameter uniform bubbles. Bubbles with size between 1.0 mm and 2.0 mm have the largest possibility to cause flux to be trapped in the growing shell at meniscus due to downward movement of liquid steel and flux. The most likely position for inclusion entrapment by this mechanism is near the SEN at wide faces.
- The single roll flow pattern imparts larger downward velocity to the liquid jet than double roll and is likely to carry bubbles deep into the caster to form gas defects such as pencil pipes.
- Compared with the flow with gas, the flow without gas has much less tendency to entrap flux particles and might have fewer inclusion problems in cases when the liquid steel is clean.

# 7.4 Effects of Slab Width and Nozzle Submergence Depth on the Fluid Flow in the Steel Caster

The effects of slab width and nozzle submergence depth were investigated with the MUSIG model developed in this study. The main findings include:

• For a given casting speed, gas fraction and submergence depth, narrow slabs tend to form double roll while wide slab tends to form single roll. Thus, the critical gas fraction range for this flow transition increases when decreasing slab width.

• Nozzle submergence depth greatly affects the flow pattern in the mold. Deeper nozzle submergence depth has more tendency to form double roll, so thereby increases the critical gas fraction range.

### Appendix A. A sample command file for heat transfer model

>>CFX4 >>OPTIONS THREE DIMENSIONS BODY FITTED GRID CARTESIAN COORDINATES TURBULENT FLOW HEAT TRANSFER **INCOMPRESSIBLE FLOW** STEADY STATE **USER SCALAR EQUATIONS 2** >>USER FORTRAN USRWTM >>VARIABLE NAMES U VELOCITY 'U VELOCITY' V VELOCITY 'V VELOCITY' W VELOCITY 'W VELOCITY' PRESSURE 'PRESSURE' DENSITY 'DENSITY' VISCOSITY 'VISCOSITY' K 'K' EPSILON 'EPSILON' **TEMPERATURE** 'TEMPERATURE' ENTHALPY 'ENTHALPY' USER SCALAR1 'YPLUS' USER SCALAR2 'WALL HEAT FLUX' >>MODEL DATA >>AMBIENT VARIABLES K 3.5000E-03 TEMPERATURE 1.7750E+03 >>TITLE PROBLEM TITLE 'ARMCO SINGLE-PHASE HEAT TRANSFER MODEL' >>PHYSICAL PROPERTIES >>FLUID PARAMETERS VISCOSITY 5.5500E-03 DENSITY 7.0200E+03 >>HEAT TRANSFER PARAMETERS THERMAL CONDUCTIVITY 2.6000E+01 FLUID SPECIFIC HEAT 1\* 6.800000E+02 ENTHALPY REFERENCE TEMPERATURE 1.7980E+03 >>TURBULENCE PARAMETERS >>TURBULENCE MODEL TURBULENCE MODEL 'K-EPSILON' >>SOLVER DATA >>PROGRAM CONTROL MAXIMUM NUMBER OF ITERATIONS 1000 MASS SOURCE TOLERANCE 1.0000E-06 **ITERATIONS OF TURBULENCE EQUATIONS 1** ITERATIONS OF VELOCITY AND PRESSURE EQUATIONS 1 **ITERATIONS OF TEMPERATURE AND SCALAR EQUATIONS 1 ITERATIONS OF HYDRODYNAMIC EQUATIONS 1** SOLVER DEBUG PRINT STREAM 20 >>DEFERRED CORRECTION **K START 1900** K END 2000 **EPSILON START 1900 EPSILON END 2000** >>REDUCTION FACTORS

U VELOCITY 3.5000E-01 V VELOCITY 3.5000E-01 W VELOCITY 3.5000E-01 PRESSURE 1.0000E-01 K 2.5000E-01 TEMPERATURE 2.5000E-01 >>UNDER RELAXATION FACTORS U VELOCITY 4.5000E-01 V VELOCITY 4.5000E-01 W VELOCITY 4.5000E-01 PRESSURE 1.0000E+00 VISCOSITY 1.0000E+00 K 7.000E-01 EPSILON 7.000E-01 TEMPERATURE 8.5000E-01 >>MODEL BOUNDARY CONDITIONS >>INLET BOUNDARIES PATCH NAME 'NOZZLE INLET' NORMAL VELOCITY 8.5700E-01 K 4.2500E-03 TEMPERATURE 1.8320E+03 >>PRESSURE BOUNDARIES PATCH NAME 'MOLD OUTLET' PRESSURE 0.0000E+00 STATIC PRESSURE SPECIFIED >>WALL BOUNDARIES PATCH NAME 'MOLD NARROW WALL' TEMPERATURE 1.7750E+03 >>WALL BOUNDARIES PATCH NAME 'MOLD WIDE WALL' TEMPERATURE 1.7750E+03 >>WALL BOUNDARIES PATCH NAME 'MOLD TOP SURFACE' TEMPERATURE ABC 0.0 0.025 300.0 >>OUTPUT OPTIONS >>FRONTEND PRINTING NO TOPOLOGY STRUCTURE >>PRINT OPTIONS >>WHAT NO VARIABLES >>LINE GRAPH DATA EACH ITERATION FILE NAME 'RESIDUALS' ALL VARIABLES RESIDUAL >>STOP

# Appendix B. A sample command file for multiphase MUSIG model using output of nozzle simulation as input conditions

>>CFX4 >>SET LIMITS TOTAL INTEGER WORK SPACE 1500000 TOTAL CHARACTER WORK SPACE 9000 TOTAL REAL WORK SPACE 9500000 MAXIMUM NUMBER OF BLOCKS 20 MAXIMUM NUMBER OF PATCHES 100 MAXIMUM NUMBER OF INTER BLOCK BOUNDARIES 40 >>OPTIONS THREE DIMENSIONS BODY FITTED GRID CARTESIAN COORDINATES TURBULENT FLOW **ISOTHERMAL FLOW** INCOMPRESSIBLE FLOW **BUOYANT FLOW** STEADY STATE **USER SCALAR EQUATIONS 1** NUMBER OF PHASES 2 NUMBER OF MUSIG SIZE GROUPS 11 >>USER FORTRAN USRSRC USRTRN USRBCS >>VARIABLE NAMES U VELOCITY 'U VELOCITY' V VELOCITY 'V VELOCITY' W VELOCITY 'W VELOCITY' PRESSURE 'PRESSURE' VOLUME FRACTION 'VOLUME FRACTION' DENSITY 'DENSITY' VISCOSITY 'VISCOSITY' K 'K' EPSILON 'EPSILON' USER SCALAR1 'YPLUS' >>MODEL DATA >>AMBIENT VARIABLES K 5.0200E-02 EPSILON 4.5700E-01 >>TITLE PROBLEM TITLE 'LTV MULTIPHASE MUSIG' >>PHYSICAL PROPERTIES >>BUOYANCY PARAMETERS ALL PHASES GRAVITY VECTOR 0.000000E+00 0.000000E+00 -9.800000E+00 >>FLUID PARAMETERS PHASE NAME 'PHASE1' VISCOSITY 5.5500E-03 DENSITY 7.0200E+03 >>FLUID PARAMETERS PHASE NAME 'PHASE2' VISCOSITY 7.4200E-05 **DENSITY 5.5900E-01** >>MULTIPHASE PARAMETERS >>PHASE DESCRIPTION PHASE NAME 'PHASE1'

LIQUID CONTINUOUS >>PHASE DESCRIPTION PHASE NAME 'PHASE2' GAS DISPERSE MEAN DIAMETER 2.5900E-03 >>MULTIPHASE MODELS >>MOMENTUM INTER PHASE TRANSFER SINCE >>INTER PHASE TRANSFER MODELS >>MOMENTUM FIRST PHASE NAME 'PHASE1' SECOND PHASE NAME 'PHASE2' SURFACE TENSION COEFFICIENT 1.1920E+00 >>PARTICLE DRAG MODEL FLOW REGIME 'VISCOUS' >>MUSIG MODEL PHASE NAME 'PHASE2' MAXIMUM DIAMETER 1.0500E-02 MINIMUM DIAMETER 5.0000E-04 EQUAL DIAMETER DIVISION >>BREAKUP MODEL LUO SVENDSEN BREAKUP COEFFICIENT 4.0000E+00 >>COALESCENCE MODEL PRINCE BLANCH COALESCENCE COEFFICIENT 5.0000E-02 INITIAL FILM THICKNESS 1.0000E-04 CRITICAL FILM THICKNESS 1.0000E-08 >>TURBULENCE PARAMETERS >>TURBULENCE MODEL PHASE NAME 'PHASE1' TURBULENCE MODEL 'K-EPSILON' PARTICLE INDUCED TURBULENCE 'NONE' >>TURBULENCE MODEL PHASE NAME 'PHASE2' TURBULENCE MODEL 'LAMINAR' PARTICLE INDUCED TURBULENCE 'NONE' >>SOLVER DATA >>PROGRAM CONTROL MAXIMUM NUMBER OF ITERATIONS 500 MASS SOURCE TOLERANCE 1.0000E-08 **ITERATIONS OF TURBULENCE EQUATIONS 1** ITERATIONS OF VELOCITY AND PRESSURE EQUATIONS 1 **ITERATIONS OF TEMPERATURE AND SCALAR EQUATIONS 1 ITERATIONS OF HYDRODYNAMIC EQUATIONS 1** SOLVER DEBUG PRINT STREAM 200 >>DEFERRED CORRECTION **K START 2000** K END 2201 **EPSILON START 2000 EPSILON END 2201** >>UNDER RELAXATION FACTORS PHASE NAME 'PHASE1' U VELOCITY 5.5000E-01 V VELOCITY 5.5000E-01

W VELOCITY 5.5000E-01 PRESSURE 9.0000E-01

VOLUME FRACTION 9.000E-01 VISCOSITY 1.0000E+00 K 7.0000E-01 EPSILON 6.0000E-01 >>UNDER RELAXATION FACTORS PHASE NAME 'PHASE2' U VELOCITY 5.0000E-01 V VELOCITY 5.0000E-01 W VELOCITY 5.0000E-01 PRESSURE 9.0000E-01 VOLUME FRACTION 9.000E-01 VISCOSITY 1.0000E+00 K 5.0000E-01 EPSILON 5.0000E-01 ALL MUSIG VOL FRACTIONS 8.0000E-01 >>MODEL BOUNDARY CONDITIONS >>INLET BOUNDARIES PHASE NAME 'PHASE1' PATCH NAME 'NOZZLE INLET' U VELOCITY 1.41270E-01 V VELOCITY 0.0000E+00 W VELOCITY -4.05100E-01 VOLUME FRACTION 8.900E-01 K 5.0200E-02 EPSILON 4.5700E-01 >>INLET BOUNDARIES PHASE NAME 'PHASE2' PATCH NAME 'NOZZLE INLET' U VELOCITY 1.41270E-01 V VELOCITY 0.0000E+00 W VELOCITY -4.05100E-01 VOLUME FRACTION 1.1000E-01 K 1.0000E-04 EPSILON 1.0000E-04 MUSIG VOL FRACTION1 0.0107 MUSIG VOL FRACTION2 0.0453 MUSIG VOL FRACTION3 0.3115 MUSIG VOL FRACTION4 0.5583 MUSIG VOL FRACTION5 0.0742 MUSIG VOL FRACTION6 0.0 MUSIG VOL FRACTION7 0.0000 MUSIG VOL FRACTION8 0.0000 MUSIG VOL FRACTION9 0.0000 MUSIG VOL FRACTION10 0.0 MUSIG VOL FRACTION11 0.0 >>PRESSURE BOUNDARIES PHASE NAME 'PHASE1' PATCH NAME 'OUTLET' PRESSURE 0.0000E+00 VOLUME FRACTION 1.0000E+00 >>PRESSURE BOUNDARIES PHASE NAME 'PHASE2' PATCH NAME 'OUTLET' PRESSURE 0.0000E+00 >>WALL BOUNDARIES PHASE NAME 'PHASE1' PATCH NAME 'TOP SURFACE' U VELOCITY 0.0000E+00 V VELOCITY 0.0000E+00 W VELOCITY 0.0000E+00

>>WALL BOUNDARIES ALL PHASES PATCH NAME 'WIDE FACE' U VELOCITY 0.0000E+00 V VELOCITY 0.0000E+00 W VELOCITY 0.0000E+00 >>WALL BOUNDARIES ALL PHASES PATCH NAME 'NARROW FACE' U VELOCITY 0.0000E+00 V VELOCITY 0.0000E+00 W VELOCITY 0.0000E+00 >>OUTPUT OPTIONS >>FRONTEND PRINTING NO TOPOLOGY STRUCTURE >>PRINT OPTIONS >>WHAT NO VARIABLES >>LINE GRAPH DATA EACH ITERATION FILE NAME 'RESIDUALS' ALL VARIABLES RESIDUAL >>STOP

#### Appendix C. Subroutine for heat transfer model

SUBROUTINE USRWTM(CNAME, IEQN, IPHASE, CTROP, CWALL + ,GAMMA,YWALL,TMULT,ISTART,NCV ,DIFLAM,XYPLUS,ELOG,PRT + ,U,V,W,P,VFRAC,DEN,VIS,TE,ED,RS,T,H,RF,SCAL + ,XP,YP,ZP,VOL,AREA,VPOR,ARPOR,WFACT,IPT + ,IBLK,IPVERT,IPNODN,IPFACN,IPNODF,IPNODB,IPFACB + ,WORK,IWORK,CWORK) С С C USER SUBROUTINE TO OVERWRITE DEFAULT WALL MULTIPLIERS (TMULT) C FLUX = TMULT \* (VARIABLE AT WALL - VARIABLE AT NEAREST NODE) С С C\* С \*\*\*\*\*\*\* C\* С LOGICAL LDEN, LVIS, LTURB, LTEMP, LBUOY, LSCAL, LCOMP ,LRECT,LCYN,LAXIS,LPOROS,LTRANS + С CHARACTER\*(\*) CWALL, CWORK CHARACTER CNAME\*6,CTROP\*6 С C---- AREA FOR USERS EXPLICITLY DECLARED VARIABLES С С COMMON + /ALL/ NBLOCK,NCELL,NBDRY,NNODE,NFACE,NVERT,NDIM + /ADDIMS/ NPHASE.NSCAL.NVAR.NPROP ,NDVAR,NDPROP,NDXNN,NDGEOM,NDCOEF,NILIST,NRLIST,NTOPOL + + /CHKUSR/ IVERS, IUCALL, IUSED + /DEVICE/ NREAD,NWRITE,NRDISK,NWDISK + /IDUM/ ILEN,JLEN + /LOGIC/ LDEN,LVIS,LTURB,LTEMP,LBUOY,LSCAL,LCOMP ,LRECT,LCYN,LAXIS,LPOROS,LTRANS + + /MLTGRD/ MLEVEL,NLEVEL,ILEVEL + /SGLDBL/ IFLGPR,ICHKPR + /SPARM/ SMALL,SORMAX,NITER,INDPRI,MAXIT,NODREF,NODMON + /TRANSI/ NSTEP,KSTEP,MF,INCORE + /TRANSR/ TIME, DT, DTINVF, TPARM + /TURBKE/ CMU,C1,C2,C3,CAPPA,CMU14,CMU34 С C---- AREA FOR USERS TO DECLARE THEIR OWN COMMON BLOCKS C THESE SHOULD START WITH THE CHARACTERS 'UC' TO ENSURE С NO CONFLICT WITH NON-USER COMMON BLOCKS С С DIMENSION CWALL(NPHASE), GAMMA(NCELL, NPHASE,\*) +.YWALL(NCV).TMULT(NCV.NPHASE) +,DIFLAM(NVAR,NPHASE),XYPLUS(NVAR,NPHASE),ELOG(NVAR,NPHASE) +,PRT(NVAR,NPHASE)

```
DIMENSION
 + U(NNODE,NPHASE),V(NNODE,NPHASE),W(NNODE,NPHASE),P(NNODE,NPHASE)
 +,VFRAC(NNODE,NPHASE),DEN(NNODE,NPHASE),VIS(NNODE,NPHASE)
 +,TE(NNODE,NPHASE),ED(NNODE,NPHASE),RS(NNODE,NPHASE,6)
 +,T(NNODE,NPHASE),H(NNODE,NPHASE),RF(NNODE,NPHASE,4)
  +,SCAL(NNODE,NPHASE,NSCAL)
  DIMENSION
 + XP(NNODE), YP(NNODE), ZP(NNODE)
  +,VOL(NCELL),AREA(NFACE,3),VPOR(NCELL),ARPOR(NFACE,3),WFACT(NFACE)
  +,IPT(*),IBLK(5,NBLOCK)
 +,IPVERT(NCELL,8),IPNODN(NCELL,6),IPFACN(NCELL,6),IPNODF(NFACE,4)
 +, IPNODB(NBDRY, 4), IPFACB(NBDRY)
  +,IWORK(*),WORK(*),CWORK(*)
С
C---- AREA FOR USERS TO DIMENSION THEIR ARRAYS
С
C---- AREA FOR USERS TO DEFINE DATA STATEMENTS
С
С
C---- STATEMENT FUNCTION FOR ADDRESSING
  IP(I,J,K)=IPT((K-1)*ILEN*JLEN+(J-1)*ILEN+I)
С
C----VERSION NUMBER OF USER ROUTINE AND PRECISION FLAG
С
  IVERS = 1
  ICHKPR = 1
С
C---- TO USE THIS USER ROUTINE FIRST SET IUSED=1
С
  IUSED = 1
С
С
  IF (IUSED.EQ.0) RETURN
С
C---- FRONTEND CHECKING OF USER ROUTINE
  IF (IUCALL.EQ.0) RETURN
С
С
C-----EXAMPLE - TO USE A MOLECULAR VISCOSITY DEPENDENT ON TEMPERATURE
C IN A SMOOTH WALL CALCULATION.
С
  THE EXAMPLE HAS MOLECULAR VISCOSITY = ATEMP*TEMPERATURE + BTEMP.
C THE USER WOULD IN GENERAL SUPPLY A FUNCTIONAL RELATIONSHIP FOR
C MOLECULAR VISCOSITY IN TERMS OF OTHER VARIABLES.
  CALL IPALL('MOLD NARROW WALL','*','PATCH','CENTRES',IPT,NPT,
        CWORK, IWORK)
  +
С
C FIND VARIABLE NUMBER FOR VELOCITY
   MAX_NF=0.0
   WALLHF_NF=0.0
  DO 101 I=1,NPT
  INODE=IPT(I)
  IBDRY=INODE-NCELL
  IFACE=IPFACB(IBDRY)
  INODE2=IPNODF(IFACE,1)
  AREAM=SQRT(AREA(IFACE,1)*AREA(IFACE,1)
```

```
+AREA(IFACE,2)*AREA(IFACE,2)
  +
          +AREA(IFACE,3)*AREA(IFACE,3))
  +
   IF (ABS(SCAL(INODE,1,2)).GT.MAX_NF) THEN
     MAX_NF=ABS(SCAL(INODE,1,2))
   END IF
   WALLHF_NF=WALLHF_NF+SCAL(INODE,1,2)*AREAM
101 CONTINUE
   CALL IPALL('MOLD WIDE WALL','*','PATCH','CENTRES',IPT,NPT,
  +
         CWORK, IWORK)
C FIND VARIABLE NUMBER FOR VELOCITY
    MAX WF=0.0
    WALLHF WF=0.0
   DO 102 J=1,NPT
   INODE=IPT(J)
   IBDRY=INODE-NCELL
   IFACE=IPFACB(IBDRY)
   AREAM=SQRT(AREA(IFACE,1)*AREA(IFACE,1)
          +AREA(IFACE,2)*AREA(IFACE,2)
  +
          +AREA(IFACE,3)*AREA(IFACE,3))
   +
   IF (ABS(SCAL(INODE,1,2)).GT.MAX_WF) THEN
     MAX_WF=ABS(SCAL(INODE,1,2))
   END IF
   WALLHF_WF=WALLHF_WF+SCAL(INODE,1,2)*AREAM
102 CONTINUE
С
   CALL IPALL('MOLD TOP SURFACE', '*', 'PATCH', 'CENTRES', IPT, NPT,
         CWORK, IWORK)
  +
С
C FIND VARIABLE NUMBER FOR VELOCITY
   AREA_TOP=0.0
   WALLHF_TOP=0.0
   DO 103 I=1,NPT
   INODE=IPT(I)
   IBDRY=INODE-NCELL
   IFACE=IPFACB(IBDRY)
   INODE2=IPNODF(IFACE,1)
   AREAM=SQRT(AREA(IFACE,1)*AREA(IFACE,1)
          +AREA(IFACE,2)*AREA(IFACE,2)
  +
          +AREA(IFACE,3)*AREA(IFACE,3))
  +
   AREA_TOP=AREA_TOP+AREAM
   WALLHF_TOP=WALLHF_TOP-SCAL(INODE,1,2)*AREAM
103 CONTINUE
   TOP_HF=-12000*AREA_TOP
С
   CALL IPALL('NOZZLE INLET', '*', 'PATCH', 'CENTRES', IPT, NPT,
         CWORK, IWORK)
  +
С
C FIND VARIABLE NUMBER FOR VELOCITY
   SUM=0.0
    HF_IN=0.0
   TMASS_IN=0.0
   DO 104 I=1,NPT
   INODE=IPT(I)
   IBDRY=INODE-NCELL
   IFACE=IPFACB(IBDRY)
   AREAM=SQRT(AREA(IFACE,1)*AREA(IFACE,1)
   +
          +AREA(IFACE,2)*AREA(IFACE,2)
```

```
+AREA(IFACE,3)*AREA(IFACE,3))
  +
    TMASS_IN=TMASS_IN+7020*W(INODE,1)*AREAM
    HF_IN=HF_IN+680*7020*W(INODE,1)*AREAM*(T(INODE,1)-1775)
104 CONTINUE
С
   CALL IPALL('MOLD OUTLET', '*', 'PATCH', 'CENTRES', IPT, NPT,
         CWORK, IWORK)
  +
С
C FIND VARIABLE NUMBER FOR VELOCITY
   HF OUT=0.0
   AREA_OUT=0.0
   TEMP3=0.0
   TMASS OUT=0.0
   DO 105 I=1,NPT
   INODE=IPT(I)
   IBDRY=INODE-NCELL
   IFACE=IPFACB(IBDRY)
   ISURF1=IPNODB(IBDRY,1)
   ISURF2=IPNODB(IBDRY,2)
   INODE2=IPNODF(IFACE,1)
   AREAM=SQRT(AREA(IFACE,1)*AREA(IFACE,1)
          +AREA(IFACE,2)*AREA(IFACE,2)
  +
          +AREA(IFACE,3)*AREA(IFACE,3))
   +
   TMASS_OUT=TMASS_OUT+7020*W(ISURF1,1)*AREAM
    HF_OUT=HF_OUT-680*7020*W(ISURF1,1)*AREAM*(T(ISURF1,1)-1775)
    TEMP3=TEMP3+T(INODE,1)*AREAM
С
   AREA_OUT=AREA_OUT+AREAM
105 CONTINUE
С
С
    AVRG_TEMP=TEMP3/AREA_OUT
С
     SUM2=VOLUME*(AVRG_TEMP-298)*680*7020
С
С
   OUT_HF=HF_OUT+WALLHF_NF+WALLHF_WF+TOP_HF
   HEAT_BALANCE=HF_IN+OUT_HF
   ERROR=100*HEAT_BALANCE/HF_IN
   TMASS_DIFF=ABS(TMASS_OUT-TMASS_IN)
   TMASS_ERR=TMASS_DIFF/ABS(TMASS_IN)
С
   WRITE(NWRITE,*) ' '
   WRITE(NWRITE,*) '**** INPUT RATE OF HEAT AND MASS****'
   WRITE(NWRITE,*) 'INPUT HEAT RATE FROM INLET=',HF_IN
   WRITE(NWRITE,*) 'INPUT MASS RATE FROM INLET=',ABS(TMASS_IN)
   WRITE(NWRITE,*) ' '
   WRITE(NWRITE,*) '**** OUTGOING RATE OF HEAT AND MASS****'
   WRITE(NWRITE,*) 'HEAT RATE ON NARROW FACE=',WALLHF_NF
   WRITE(NWRITE,*) 'HEAT RATE ON WIDE FACE=',WALLHF_WF
   WRITE(NWRITE,*) 'HEAT RATE ON TOP SURFACE=',WALLHF_TOP
   WRITE(NWRITE,*) 'OUTGOING HEAT RATE FROM OUTLET=',HF_OUT
   WRITE(NWRITE,*) 'TOTAL OUTGOING HEAT RATE=',OUT_HF
   WRITE(NWRITE,*) 'TOTAL OUTGOING MASS RATE=',ABS(TMASS_OUT)
   WRITE(NWRITE,*) ' '
   WRITE(NWRITE,*) '****HEAT BALANCE****'
   WRITE(NWRITE,*) 'HEAT BALANCE=',HEAT_BALANCE,"W"
   WRITE(NWRITE,*) 'ERROR of HEAT BALANCE=',ERROR,"%"
   WRITE(NWRITE,*) 'MASS BALANCE=',TMASS_DIFF,"KG"
   WRITE(NWRITE,*) 'ERROR of MASS BALANCE=',TMASS_ERR,"%"
   WRITE(NWRITE,*) ' '
```

```
WRITE(NWRITE,*) '****MONITORING TERMS****'
   WRITE(NWRITE,*) 'HEAT RATE ON TOP SURFACE(1)=',WALLHF_TOP
   WRITE(NWRITE,*) 'HEAT RATE ON TOP SURFACE(2)=',TOP_HF
   WRITE(NWRITE,*) 'TOTAL AREA of OUTLET=',AREA_OUT
   WRITE(NWRITE,*) 'TOTAL NUMBER OF CELLS=',NCELL
   WRITE(NWRITE,*) 'MAX. WALL HEAT FLUX on NF=',MAX_NF
   WRITE(NWRITE,*) 'MAX. WALL HEAT FLUX on WF=',MAX_WF
С
   CALL IPALL('*','WALL','PATCH','CENTRES',IPT,NPT,CWORK,IWORK)
С
C FIND VARIABLE NUMBER FOR ENTHALPY
   CALL GETVAR('USRWTM','H ',IVAR)
C IF ENTHALPY EQUATION SET MULTIPLIER
  IF (IVAR.EQ.IEQN) THEN
    PRANDT = PRT(IVAR,1)
   DO 120 I = 1, NPT
    INODE = IPT(I)
    IBDRY = INODE - NCELL
    LCV = IBDRY - ISTART + 1
    INODE1 = IPNODB(IBDRY,1)
    CMU = 0.09
    DENS = DEN(INODE1,1)
    DENSQK = DEN(INODE1,1) * SQRT( TE(INODE1,1) )
    EPSILON = ED(INODE, 1)
    AKE = TE(INODE1,1)
    AKK = AKE^*AKE
C NORMAL DISTANCE FROM NODE TO WALL
С
    DN = YWALL(LCV)
С
    WRITE(NWRITE,*) YWALL(LCV)
С
C COMPUTE MULTIPLIER
    TMULT(LCV,1) = (CMU*DENS*AKK)/(PRANDT*DN*EPSILON)
120 CONTINUE
С
  END IF
С
С
  RETURN
  END
```

# Appendix D. Subroutine for multiphase flow with MUSIG model

```
SUBROUTINE USRSRC(IEQN,ICALL,CNAME,CALIAS,AM,SP,SU,CONV
            ,U,V,W,P,VFRAC,DEN,VIS,TE,ED,RS,T,H,RF,SCAL
  +
            ,XP,YP,ZP,VOL,AREA,VPOR,ARPOR,WFACT,IPT
  +
            ,IBLK,IPVERT,IPNODN,IPFACN,IPNODF,IPNODB,IPFACB
  +
            ,WORK,IWORK,CWORK)
С
  LOGICAL LDEN, LVIS, LTURB, LTEMP, LBUOY, LSCAL, LCOMP
       ,LRECT,LCYN,LAXIS,LPOROS,LTRANS
  +
С
   CHARACTER*(*) CWORK
  CHARACTER CNAME*6. CALIAS*24
С
С
  COMMON
   + /ALL/ NBLOCK,NCELL,NBDRY,NNODE,NFACE,NVERT,NDIM
   + /ALLWRK/ NRWS,NIWS,NCWS,IWRFRE,IWIFRE,IWCFRE
   + /ADDIMS/ NPHASE,NSCAL,NVAR,NPROP
        ,NDVAR,NDPROP,NDXNN,NDGEOM,NDCOEF,NILIST,NRLIST,NTOPOL
  + /CHKUSR/ IVERS,IUCALL,IUSED
  + /DEVICE/ NREAD,NWRITE,NRDISK,NWDISK
  + /IDUM/ ILEN, JLEN
  + /LOGIC/ LDEN,LVIS,LTURB,LTEMP,LBUOY,LSCAL,LCOMP
        ,LRECT,LCYN,LAXIS,LPOROS,LTRANS
  +
  + /MLTGRD/ MLEVEL,NLEVEL,ILEVEL
  + /SGLDBL/ IFLGPR,ICHKPR
  + /SPARM/ SMALL,SORMAX,NITER,INDPRI,MAXIT,NODREF,NODMON
  + /TRANSI/ NSTEP, KSTEP, MF, INCORE
   + /TRANSR/ TIME, DT, DTINVF, TPARM
С
С
  COMMON /UCSURF/ JISURF,NISURF,JDSURF,NDSURF
  COMMON /UCINJR/ FINJ1,FINJ2,UGASIN,VGASIN,WGASIN
  COMMON /UCFLUX/ GASIN, GASOUT
С
С
  DIMENSION AM(NCELL,6,NPHASE),SP(NCELL,NPHASE),SU(NCELL,NPHASE)
  +,CONV(NFACE,NPHASE)
С
  DIMENSION
  + U(NNODE,NPHASE),V(NNODE,NPHASE),W(NNODE,NPHASE),P(NNODE,NPHASE)
  +,VFRAC(NNODE,NPHASE),DEN(NNODE,NPHASE),VIS(NNODE,NPHASE)
  +,TE(NNODE,NPHASE),ED(NNODE,NPHASE),RS(NNODE,NPHASE,6)
  +,T(NNODE,NPHASE),H(NNODE,NPHASE),RF(NNODE,NPHASE,4)
  +,SCAL(NNODE,NPHASE,NSCAL)
С
  DIMENSION
  + XP(NNODE), YP(NNODE), ZP(NNODE)
  +,VOL(NCELL),AREA(NFACE,3),VPOR(NCELL),ARPOR(NFACE,3)
  +,WFACT(NFACE)
  +,IPT(*),IBLK(5,NBLOCK)
  +,IPVERT(NCELL,8),IPNODN(NCELL,6),IPFACN(NCELL,6),IPNODF(NFACE,4)
  +,IPNODB(NBDRY,4),IPFACB(NBDRY)
  +,IWORK(*),WORK(*),CWORK(*)
С
C---- STATEMENT FUNCTION FOR ADRESSING
```

IP(I,J,K)=IPT((K-1)\*ILEN\*JLEN+(J-1)\*ILEN+I) С C----VERSION NUMBER OF USER ROUTINE AND PRECISION FLAG С IVERS=4 ICHKPR = 1 С С IUSED=1 IF (IUSED.EQ.0) RETURN С C---- FRONTEND CHECKING OF USER ROUTINE IF (IUCALL.EQ.0) RETURN С C---- ADD TO SOURCE TERMS IF (ICALL.EQ.1) THEN С С С C-----C VOLUME FRACTION [M]/[T] С IF(CALIAS.EQ.'VOLUME FRACTION') THEN С C.....SINK AT TOP CALL USRDEGAS(CALIAS,SP,W,VFRAC,DEN,AREA,IWORK(JDSURF),NDSURF + ,IPNODB,IPFACB,TOTAL) WRITE(NWRITE,\*) ' VOLUME GASOUT=:',TOTAL WRITE(NWRITE,\*) ' VF :',TOTAL С С ENDIF С C-----C PRESSURE [M]/[T] С IF(CALIAS.EQ.'PRESSURE') THEN С GASIN = 0.0С C.....SINK AT TOP CALL USRDEGAS(CALIAS,SU,W,VFRAC,DEN,AREA,IWORK(JDSURF),NDSURF ,IPNODB,IPFACB,TOTAL) + GASOUT = TOTAL WRITE(NWRITE,\*) 'PRESSURE GASOUT=:',TOTAL С ENDIF С C-----C U VELOCITY [M][L]/[T]^2 С IF(CALIAS.EQ.'U VELOCITY') THEN С С C.....SINK AT TOP CALL USRDEGAS(CALIAS,SP,W,VFRAC,DEN,AREA,IWORK(JDSURF),NDSURF ,IPNODB,IPFACB,TOTAL) С WRITE(NWRITE,\*) ' U :',TOTAL С WRITE(NWRITE,\*) ' U GASOUT=:',TOTAL

С

```
C.....SET MINIMUM VOID FRACTION
    VFMIN=1.0E-8
   DO 101 INODE=1,NNODE
    VFRAC(INODE,2)=MAX(VFRAC(INODE,2),VFMIN)
101 CONTINUE
С
  ENDIF
С
C-----
      _____
C V VELOCITY [M][L]/[T]^2
С
  IF(CALIAS.EQ.'V VELOCITY') THEN
С
С
C.....SINK AT TOP
   CALL USRDEGAS(CALIAS,SP,W,VFRAC,DEN,AREA,IWORK(JDSURF),NDSURF
         ,IPNODB,IPFACB,TOTAL)
  +
C WRITE(NWRITE,*) ' V :',TOTAL
C.....SET MINIMUM VOID FRACTION
   VFMIN=1.0E-8
   DO 201 INODE=1,NNODE
    VFRAC(INODE,2)=MAX(VFRAC(INODE,2),VFMIN)
201 CONTINUE
С
  ENDIF
С
C-----
C W VELOCITY [M][L]/[T]^2
С
  IF(CALIAS.EQ.'W VELOCITY') THEN
С
C.....SINK AT TOP
   CALL USRDEGAS(CALIAS,SP,W,VFRAC,DEN,AREA,IWORK(JDSURF),NDSURF
         ,IPNODB,IPFACB,TOTAL)
  +
С
    WRITE(NWRITE,*) ' W :',TOTAL
С
C.....SET MINIMUM VOID FRACTION
   VFMIN=1.0E-8
   DO 301 INODE=1,NNODE
    VFRAC(INODE,2)=MAX(VFRAC(INODE,2),VFMIN)
301 CONTINUE
С
  ENDIF
С
ENDIF
С
C---- OVERWRITE SOURCE TERMS
  IF (ICALL.EQ.2) THEN
С
  ENDIF
  RETURN
  END
С
С
  SUBROUTINE USRDEGAS(CALIAS,SUP,W,VFRAC,DEN,AREA,IPTSRF,NDSURF
  +
            ,IPNODB,IPFACB,TOTAL)
С
```

С C UTILITY SUBROUTINE FOR USER-SUPPLIED FREE SURFACE DEGASSING С C, С С THIS SUBROUTINE IS CALLED BY THE FOLLOWING SUBROUTINES С USRSRC С C\*\*\*\*\* \*\*\*\*\* CHARACTER CALIAS\*24 С C---- AREA FOR USERS EXPLICITLY DECLARED VARIABLES С С COMMON + /ALL/ NBLOCK,NCELL,NBDRY,NNODE,NFACE,NVERT,NDIM + /ALLWRK/ NRWS,NIWS,NCWS,IWRFRE,IWIFRE,IWCFRE + /ADDIMS/ NPHASE,NSCAL,NVAR,NPROP ,NDVAR,NDPROP,NDXNN,NDGEOM,NDCOEF,NILIST,NRLIST,NTOPOL + + /DEVICE/ NREAD,NWRITE,NRDISK,NWDISK С С DIMENSION SUP(NCELL,NPHASE) +,W(NNODE,NPHASE),VFRAC(NNODE,NPHASE),DEN(NNODE,NPHASE) +,AREA(NFACE,3),IPTSRF(NDSURF),IPNODB(NBDRY,4),IPFACB(NBDRY) С С C---- STATEMENT FUNCTION FOR ADRESSING С IP(I,J,K)=IPT((K-1)\*ILEN\*JLEN+(J-1)\*ILEN+I)С С ----DEGASSING PHASE INDEX C-TOTAL=0.0 С C-------C VOLUME FRACTION С IF(CALIAS.EQ.'VOLUME FRACTION') THEN С DO 101 I=1,NDSURF INODE = IPTSRF(I)IBDRY = INODE-NCELL INODE1= IPNODB(IBDRY,1) IFACE = IPFACB(IBDRY)AREAM = SQRT ( AREA(IFACE,1)\*AREA(IFACE,1) +AREA(IFACE,2)\*AREA(IFACE,2) + +AREA(IFACE,3)\*AREA(IFACE,3) ) FLUX = DEN(INODE1,2)\* AREAM\*MAX(W(INODE1,2),0.0) SUP(INODE1,2)=SUP(INODE1,2)-FLUX IF(I.EQ.5) THEN WRITE(NWRITE,\*) ' AREA= ', AREAM WRITE(NWRITE,\*) 'W= ', W(INODE1,2) WRITE(NWRITE,\*) 'FLUX= ', FLUX WRITE(NWRITE,\*) 'TOTAL== ', TOTAL

WRITE(NWRITE,\*) 'VOL.FRAC= ', VFRAC(INODE1,2)

```
ENDIF
   TOTAL = TOTAL+FLUX
С
   WRITE(NWRITE,*) ' FLUX=',FLUX
101 CONTINUE
С
C-----
C OTHER VARIABLES
С
  ELSE
С
  DO 201 I=1,NDSURF
   INODE = IPTSRF(I)
   IBDRY = INODE-NCELL
   INODE1= IPNODB(IBDRY,1)
   IFACE = IPFACB(IBDRY)
   AREAM = SQRT ( AREA(IFACE,1)*AREA(IFACE,1)
          +AREA(IFACE,2)*AREA(IFACE,2)
  +
          +AREA(IFACE,3)*AREA(IFACE,3) )
  +
   FLUX = VFRAC(INODE1,2)*DEN(INODE1,2)
        *AREAM*MAX(W(INODE1,2),0.0)
  +
   SUP(INODE1,2)=SUP(INODE1,2)-FLUX
      TOTAL = TOTAL+FLUX
   WRITE(NWRITE,*) ' FLUX=',FLUX
С
201 CONTINUE
С
  ENDIF
401 CONTINUE
С
  RETURN
  END
С
С
  SUBROUTINE USRTRN(U,V,W,P,VFRAC,DEN,VIS,TE,ED,RS,T,H,RF,SCAL,
  +
           XP,YP,ZP,VOL,AREA,VPOR,ARPOR,WFACT,CONV,IPT,
           IBLK, IPVERT, IPNODN, IPFACN, IPNODF, IPNODB, IPFACB,
  +
           WORK, IWORK, CWORK)
  +
С
С
C USER SUBROUTINE TO ALLOW USERS TO MODIFY OR MONITOR THE SOLUTION AT
C THE END OF EACH TIME STEP
C THIS SUBROUTINE IS CALLED BEFORE THE START OF THE RUN AS WELL AS AT
C THE END OF EACH TIME STEP
С
С
C THIS SUBROUTINE IS CALLED BY THE FOLLOWING SUBROUTINES
С
  CUSR TRNMOD
С
C****
         С
        *****
C,
С
  LOGICAL LDEN, LVIS, LTURB, LTEMP, LBUOY, LSCAL, LCOMP
      ,LRECT,LCYN,LAXIS,LPOROS,LTRANS
  +
С
  CHARACTER*(*) CWORK
С
```

C---- AREA FOR USERS EXPLICITLY DECLARED VARIABLES

С

```
C
  COMMON
  + /ALL/ NBLOCK,NCELL,NBDRY,NNODE,NFACE,NVERT,NDIM
  + /ALLWRK/ NRWS,NIWS,NCWS,IWRFRE,IWIFRE,IWCFRE
  + /ADDIMS/ NPHASE,NSCAL,NVAR,NPROP
       ,NDVAR,NDPROP,NDXNN,NDGEOM,NDCOEF,NILIST,NRLIST,NTOPOL
  + /CHKUSR/ IVERS,IUCALL,IUSED
  + /CONC/ NCONC
  + /DEVICE/ NREAD,NWRITE,NRDISK,NWDISK
  + /IDUM/ ILEN, JLEN
  + /LOGIC/ LDEN,LVIS,LTURB,LTEMP,LBUOY,LSCAL,LCOMP
       ,LRECT,LCYN,LAXIS,LPOROS,LTRANS
  +
  + /MLTGRD/ MLEVEL,NLEVEL,ILEVEL
  + /SGLDBL/ IFLGPR,ICHKPR
  + /SPARM/ SMALL,SORMAX,NITER,INDPRI,MAXIT,NODREF,NODMON
  + /TIMUSR/ DTUSR
  + /TRANSI/ NSTEP,KSTEP,MF,INCORE
  + /TRANSR/ TIME,DT,DTINVF,TPARM
С
C---- AREA FOR USERS TO DECLARE THEIR OWN COMMON BLOCKS
  THESE SHOULD START WITH THE CHARACTERS 'UC' TO ENSURE
C
С
  NO CONFLICT WITH NON-USER COMMON BLOCKS
С
  COMMON / UCSURF/ JISURF, NISURF, JDSURF, NDSURF
  COMMON /UCINJR/ FINJ1, FINJ2, UGASIN, VGASIN, WGASIN
  COMMON /UCFLUX/ GASIN, GASOUT
С
С
  DIMENSION
 + U(NNODE,NPHASE),V(NNODE,NPHASE),W(NNODE,NPHASE),P(NNODE,NPHASE)
 +,VFRAC(NNODE,NPHASE),DEN(NNODE,NPHASE),VIS(NNODE,NPHASE)
 +,TE(NNODE,NPHASE),ED(NNODE,NPHASE),RS(NNODE,NPHASE,6)
 +,T(NNODE,NPHASE),H(NNODE,NPHASE),RF(NNODE,NPHASE,4)
 +,SCAL(NNODE,NPHASE,NSCAL)
  DIMENSION
 + XP(NNODE), YP(NNODE), ZP(NNODE)
 +,VOL(NCELL),AREA(NFACE,3),VPOR(NCELL),ARPOR(NFACE,3)
 +,WFACT(NFACE),CONV(NFACE,NPHASE)
  +,IPT(*),IBLK(5,NBLOCK)
 +,IPVERT(NCELL,8),IPNODN(NCELL,6),IPFACN(NCELL,6),IPNODF(NFACE,4)
 +, IPNODB(NBDRY, 4), IPFACB(NBDRY)
  +,IWORK(*),WORK(*),CWORK(*)
С
С
C---- STATEMENT FUNCTION FOR ADDRESSING
   IP(I,J,K)=IPT((K-1)*ILEN*JLEN+(J-1)*ILEN+I)
С
C----VERSION NUMBER OF USER ROUTINE AND PRECISION FLAG
С
  IVERS=3
  ICHKPR = 1
С
C---- TO USE THIS USER ROUTINE FIRST SET IUSED=1
С
   IUSED=1
```

```
IF (IUSED.EQ.0) RETURN
С
C---- FRONTEND CHECKING OF USER ROUTINE
   IF (IUCALL.EQ.0) RETURN
С
С
  IF (KSTEP.EQ.0) THEN
С
C-----INJECTION BOUNDARY LOCATIONS
С
   CALL IPALL('Entrance', '*', 'PATCH', 'CENTRES', IPT, NPT,
         CWORK, IWORK)
  +
С
  NISURF=NPT
  SET INJECTION BOUNDARY LIST INTO INTERGER WORK SPACE
С
   CALL SETPER('USRTRN','IWORK ','ISURF ',NISURF,JISURF)
С
   DO 101 I=1,NISURF
   IWORK(JISURF+I-1)=IPT(I)
101 CONTINUE
  WRITE(NWRITE,*) ' '
  WRITE(NWRITE,*) '******* USRTRN - INLET **********
  WRITE(NWRITE,*) ' NISURF=',NISURF,' JISURF=',JISURF
  WRITE(NWRITE,*) ' IWORK(JISURF)=',IWORK(JISURF)
  WRITE(NWRITE,*) ' IWORK(JISURF+NISURF-1)=',IWORK(JISURF+NISURF-1)
С
C-----INJECTION RATES: MASS FLUX IN kg/s, VELOCITY IN m/s
  FGASIN = 0
  UGASIN = 0
  VGASIN = 0
  WGASIN = 0
С
C-----DETERMINE TOTAL AREA OF INLET BOUNDARY AND MASS FLUX PER AREA
   TOTALA = 0.0
   DO 201 I=1,NISURF
   INODE = IWORK(JISURF+I-1)
   IBDRY = INODE-NCELL
   IFACE = IPFACB(IBDRY)
   AREAM = SQRT ( AREA(IFACE,1)*AREA(IFACE,1)
            +AREA(IFACE,2)*AREA(IFACE,2)
  +
            +AREA(IFACE,3)*AREA(IFACE,3))
    TOTALA = TOTALA+AREAM
201 CONTINUE
С
  WRITE(NWRITE,*) ' TOTAL AREA OF INLET =',TOTALA
   FINJ2 = FGASIN/TOTALA
С
С
С
C-----DEGASSING BOUNDARY LOCATIONS
С
   CALL IPALL('Top', '*', 'PATCH', 'CENTRES', IPT, NPT,
         CWORK, IWORK)
   +
С
  NDSURF=NPT
С
  SET DEGASSING BOUNDARY LIST INTO INTERGER WORK SPACE
   CALL SETPER('USRTRN','IWORK ','DSURF ',NDSURF,JDSURF)
```

```
С
```

```
DO 401 I=1,NDSURF
   IWORK(JDSURF+I-1)=IPT(I)
401 CONTINUE
  WRITE(NWRITE,*) ' '
  WRITE(NWRITE,*) '******** USRTRN - OUTLET *********
  WRITE(NWRITE,*) 'NDSURF=',NDSURF,' JDSURF=',JDSURF
  WRITE(NWRITE,*) ' IWORK(JDSURF)=',IWORK(JDSURF)
  WRITE(NWRITE,*) ' IWORK(JDSURF+NDSURF-1)=',IWORK(JDSURF+NDSURF-1)
С
  ENDIF
С
C-----
       .....
С
   END OPERATION
С
  IF(KSTEP.EQ.NSTEP) THEN
С
C====CHECK MASS BALANCE
С
C.....GAS FLOW RATE FROM SOURCE IS SET IN USRSRC
  INFLOE_SOURCE=GASIN
C....INLET
   CALL IPALL('Entrance', '*', 'PATCH', 'CENTRES', IPT, NPT,
         CWORK, IWORK)
  +
С
   GASIN=0.0
   GASIN_2=0.0
   WATIN = 0.0
C.....GASIN INITIALISED IN USRSRC
   DO 501 I=1,NPT
   INODE = IPT(I)
   IBDRY = INODE-NCELL
    IFACE = IPFACB(IBDRY)
    WATIN = WATIN+CONV(IFACE,1)
    GASIN_2 = GASIN_2+CONV(IFACE,2)
501 CONTINUE
  GASIN=GASIN 2
   SUM=0.0
  SUM=INFLOE_SOURCE+GASIN
  WRITE(NWRITE,*) ' '
  WRITE(NWRITE,*) '&&&&&& SHOWING OF OUTPUT OF INFLOW DATA &&&&&
  WRITE(NWRITE,*) ' '
  WRITE(NWRITE,*) 'GAS INFLOW RATE FROM SOURCE=',INFLOE_SOURCE
  WRITE(NWRITE,*) 'GAS INFLOW RATE FROM INLET=',GASIN
  WRITE(NWRITE,*) 'THE SUM OF GAS INFLOW RATE FROM SOURCE AND
  +INLET=',SUM
  WRITE(NWRITE,*) ' '
  WRITE(NWRITE,*) ' '
С
C.....OUTLET
  GASOUT_SINK=GASOUT
  WRITE(NWRITE,*) '&&&&&& SHOWING OF OUTPUT OF OUTFLOW DATA &&&&&
  WRITE(NWRITE,*) ' '
```

WRITE(NWRITE,\*) 'GAS OUTFLOW RATE FROM TOP SURFACE=', GASOUT\_SINK

WRITE(NWRITE,\*) ' '

```
CALL IPALL('OUTLET','*','PATCH','CENTRES',IPT,NPT,
   +
         CWORK, IWORK)
С
  GASOUT = 0.0
   GASOUT_2 = 0.0
   WATOUT = 0.0
C.....GASOUT INITIALISED IN USRSRC
   DO 601 I=1,NPT
   INODE = IPT(I)
   IBDRY = INODE-NCELL
   IFACE = IPFACB(IBDRY)
    WATOUT = WATOUT+CONV(IFACE,1)
    GASOUT_2 = GASOUT_2+CONV(IFACE,2)
601 CONTINUE
   GASOUT = GASOUT_2
   SUM1=0.0
   SUM1= GASOUT_SINK+GASOUT
  WRITE(NWRITE,*) 'GAS OUTFLOW RATE FROM BOTTUM=', GASOUT
  WRITE(NWRITE,*) 'THE SUM OF GAS OUTFLOW RATE FROM TOP SURFACE AND
  + BOTTOM=',SUM
  WRITE(NWRITE,*) ' '
  WRITE(NWRITE,*) ' '
С
  WRITE(NWRITE,*) ' *** MASS BALANCE ***'
  WRITE(NWRITE,*) ' '
  WRITE(NWRITE,*) ' AIR'
  WRITE(NWRITE,*) ' TOTAL INFLOW (kg/s) =',SUM
  WRITE(NWRITE,*) ' TOTAL OUTFLOW (kg/s) =',SUM1
   GASERR=(SUM1-SUM)*100.0/SUM
  WRITE(NWRITE,*) ' % ERROR
                               =',GASERR
  WRITE(NWRITE,*) ' '
  WRITE(NWRITE,*) 'WATER'
  WRITE(NWRITE,*) ' INFLOW (kg/s) =',WATIN
  WRITE(NWRITE,*) ' OUTFLOW (kg/s) =',WATOUT
   WATERR=(WATOUT+WATIN)*100.0/WATIN
  WRITE(NWRITE,*) ' % ERROR
                                =',WATERR
  WRITE(NWRITE,*) ' '
С
C.....SET IPLT=1 TO CLEAN UP VELOCITY FOR PLOTTING
C*****NOTE THAT THIS OPERATION WILL SLIGHTLY ALTER THE VELOCITY FIELD
C HENCE THE DUMP FILE GENERATED SHOULD BE USED FOR PLOTTING ONLY
   NOT FOR RESTART.
С
С
    IPLT=0
С
C.....FILTER OUT VELOCITY WITH SMALL VOLUME FRACTION
   IF(IPLT.EQ.1) THEN
     DO 701 INODE=1,NNODE
      FILTER=1.0-EXP(-3.0*VFRAC(INODE,2)/1.E-4)
      U(INODE,2)=U(INODE,2)*FILTER
      V(INODE,2)=V(INODE,2)*FILTER
      W(INODE,2)=W(INODE,2)*FILTER
      CONTINUE
701
С
   ENDIF
С
  ENDIF
С
```

C RETURN END

# **Appendix E. Subroutine for nozzle simulation**

SUBROUTINE USRTRN(U,V,W,P,VFRAC,DEN,VIS,TE,ED,RS,T,H,RF,SCAL, XP,YP,ZP,VOL,AREA,VPOR,ARPOR,WFACT,CONV,IPT, + IBLK, IPVERT, IPNODN, IPFACN, IPNODF, IPNODB, IPFACB, + WORK, IWORK, CWORK) + С \*\*\*\*\* C, С C USER SUBROUTINE TO ALLOW USERS TO MODIFY OR MONITOR THE SOLUTION AT C THE END OF EACH TIME STEP THIS SUBROUTINE IS CALLED BEFORE THE START OF THE RUN AS WELL AS AT С С THE END OF EACH TIME STEP С \*\*\*\*\* C\* С С THIS SUBROUTINE IS CALLED BY THE FOLLOWING SUBROUTINES С CUSR TRNMOD С С С \*\*\*\*\* С LOGICAL LDEN, LVIS, LTURB, LTEMP, LBUOY, LSCAL, LCOMP ,LRECT,LCYN,LAXIS,LPOROS,LTRANS + С CHARACTER\*(\*) CWORK С C---- AREA FOR USERS EXPLICITLY DECLARED VARIABLES С С COMMON + /ALL/ NBLOCK,NCELL,NBDRY,NNODE,NFACE,NVERT,NDIM + /ALLWRK/ NRWS,NIWS,NCWS,IWRFRE,IWIFRE,IWCFRE + /ADDIMS/ NPHASE,NSCAL,NVAR,NPROP ,NDVAR,NDPROP,NDXNN,NDGEOM,NDCOEF,NILIST,NRLIST,NTOPOL + /CHKUSR/ IVERS,IUCALL,IUSED + /CONC/ NCONC + /DEVICE/ NREAD.NWRITE.NRDISK.NWDISK + /IDUM/ ILEN,JLEN + /LOGIC/ LDEN,LVIS,LTURB,LTEMP,LBUOY,LSCAL,LCOMP ,LRECT,LCYN,LAXIS,LPOROS,LTRANS + /MLTGRD/ MLEVEL,NLEVEL,ILEVEL + /SGLDBL/ IFLGPR,ICHKPR + /SPARM/ SMALL,SORMAX,NITER,INDPRI,MAXIT,NODREF,NODMON + /TIMUSR/ DTUSR + /TRANSI/ NSTEP,KSTEP,MF,INCORE + /TRANSR/ TIME,DT,DTINVF,TPARM С C---- AREA FOR USERS TO DECLARE THEIR OWN COMMON BLOCKS C THESE SHOULD START WITH THE CHARACTERS 'UC' TO ENSURE NO CONFLICT WITH NON-USER COMMON BLOCKS С С COMMON /UCSURF/ JISURF,NISURF,JDSURF,NDSURF COMMON /UCINJR/ FINJ1, FINJ2, UGASIN, VGASIN, WGASIN COMMON / UCFLUX/ GASIN, GASOUT

 $TOT_AV_E = 0$  $GAS_AV_U = 0$  $GAS_AV_V = 0$ 

```
С
С
  DIMENSION
 + U(NNODE,NPHASE),V(NNODE,NPHASE),W(NNODE,NPHASE),P(NNODE,NPHASE)
 +,VFRAC(NNODE,NPHASE),DEN(NNODE,NPHASE),VIS(NNODE,NPHASE)
 +,TE(NNODE,NPHASE),ED(NNODE,NPHASE),RS(NNODE,NPHASE,6)
 +,T(NNODE,NPHASE),H(NNODE,NPHASE),RF(NNODE,NPHASE,4)
 +,SCAL(NNODE,NPHASE,NSCAL)
  DIMENSION
 + XP(NNODE), YP(NNODE), ZP(NNODE)
 +,VOL(NCELL),AREA(NFACE,3),VPOR(NCELL),ARPOR(NFACE,3)
 +,WFACT(NFACE),CONV(NFACE,NPHASE)
  +,IPT(*),IBLK(5,NBLOCK)
 +,IPVERT(NCELL,8),IPNODN(NCELL,6),IPFACN(NCELL,6),IPNODF(NFACE,4)
 +,IPNODB(NBDRY,4),IPFACB(NBDRY)
  +,IWORK(*),WORK(*),CWORK(*), TMP(20,20,11)
С
С
C---- STATEMENT FUNCTION FOR ADDRESSING
  IP(I,J,K)=IPT((K-1)*ILEN*JLEN+(J-1)*ILEN+I)
С
C----VERSION NUMBER OF USER ROUTINE AND PRECISION FLAG
С
  IVERS=3
  ICHKPR = 1
С
C---- TO USE THIS USER ROUTINE FIRST SET IUSED=1
С
  IUSED=1
  IF (IUSED.EQ.0) RETURN
С
C---- FRONTEND CHECKING OF USER ROUTINE
  IF (IUCALL.EQ.0) RETURN
С
С
  IF (KSTEP.EQ.NSTEP) THEN
С
C-----INJECTION BOUNDARY LOCATIONS
С
  OPEN(5,FILE='DATA1.DAT',STATUS='NEW',FORM='FORMATTED')
С
  VEL = 0.0
  TOTAL_BOOT = 0
  TOT_UP_U = 0
  TOT_UP_V = 0
  TOT_UP_W = 0
  TOT_UP_K = 0
  TOT_UP_E = 0
  TOT_AV_U = 0
  TOT_AV_V = 0
  TOT_AV_W = 0
  TOT_AV_K = 0
```

```
GAS_AV_W = 0
   GAS_FLOWRATE = 0
  LIQ_FLOWRATE = 0
   AREA_FRAC1 = 0
   AREA_FRAC2 = 0
   X_LIQ_OUT = 0
   Y_LIQ_OUT = 0
   Z_LIQ_OUT = 0
   TOTALA = 0
   IHEAD = 0
   CALL IPREC('BLOCK-NUMBER-19', 'BLOCK', 'CENTRES', IPT, ILEN,
   +
          JLEN, KLEN, CWORK, IWORK)
С
С
   DO 203 K=1,KLEN
   DO 202 I=1,ILEN
         INODE = IP(I,8,K)
     IFACE = INODE+NCELL
С
C Record the variables
     TMP(KLEN+1-K,IHEAD+I,1) = U(INODE,1)
         TMP(KLEN+1-K,IHEAD+I,2) = V(INODE,1)
         TMP(KLEN+1-K,IHEAD+I,3) = W(INODE,1)
     TMP(KLEN+1-K,IHEAD+I,4) = TE(INODE,1)
         TMP(KLEN+1-K,IHEAD+I,5) = ED(INODE,1)
         TMP(KLEN+1-K,IHEAD+I,6) = U(INODE,2)
         TMP(KLEN+1-K,IHEAD+I,7) = V(INODE,2)
         TMP(KLEN+1-K,IHEAD+I,8) = W(INODE,2)
         TMP(KLEN+1-K,IHEAD+I,9) = VFRAC(INODE,1)
         TMP(KLEN+1-K,IHEAD+I,10) = VFRAC(INODE,2)
C End of recording variables
C VEL is speed at current node
```

```
\label{eq:VEL} \begin{array}{l} \mathsf{VEL} = \mathsf{SQRT}(\mathsf{U}(\mathsf{INODE},1)^*\mathsf{U}(\mathsf{INODE},1)+\\ + & \mathsf{V}(\mathsf{INODE},1)^*\mathsf{V}(\mathsf{INODE},1)+\mathsf{W}(\mathsf{INODE},1)^*\mathsf{W}(\mathsf{INODE},1))\\ \mathsf{AREAX} = \mathsf{AREA}(\mathsf{IFACE},1)\\ \mathsf{AREAY} = \mathsf{AREA}(\mathsf{IFACE},2)\\ \mathsf{AREAZ} = \mathsf{AREA}(\mathsf{IFACE},3) \end{array}
```

C Following average variable is calculated according to Hua's Formula

TOTAL\_BOOT=TOTAL\_BOOT+VEL\*AREAX\*VFRAC(INODE,1)

- TOT\_UP\_U=TOT\_UP\_U+U(INODE,1)\*VEL\*AREAX\*VFRAC(INODE,1) TOT\_UP\_V=TOT\_UP\_V+V(INODE,1)\*VEL\*AREAX\*VFRAC(INODE,1) TOT\_UP\_W=TOT\_UP\_W+W(INODE,1)\*VEL\*AREAX\*VFRAC(INODE,1) TOT\_UP\_K=TOT\_UP\_K+TE(INODE,1)\*VEL\*AREAX\*VFRAC(INODE,1) TOT\_UP\_E=TOT\_UP\_E+ED(INODE,1)\*VEL\*AREAX\*VFRAC(INODE,1)
- C End of Hua's Average variables
- C Followinf variables are based on area weighted average
  - TOT\_AV\_U = TOT\_AV\_U+AREAX\*U(INODE,1)\*VFRAC(INODE,1) TOT\_AV\_V = TOT\_AV\_V+AREAX\*V(INODE,1)\*VFRAC(INODE,1) TOT\_AV\_W = TOT\_AV\_W+AREAX\*W(INODE,1)\*VFRAC(INODE,1)

TOT\_AV\_K = TOT\_AV\_K+AREAX\*TE(INODE,1)\*VFRAC(INODE,1) TOT\_AV\_E = TOT\_AV\_E+AREAX\*ED(INODE,1)\*VFRAC(INODE,1)

C End of area weighted average

GAS\_AV\_U = GAS\_AV\_U+AREAX\*U(INODE,2)\*VFRAC(INODE,2) GAS\_AV\_V = GAS\_AV\_V+AREAX\*V(INODE,2)\*VFRAC(INODE,2) GAS\_AV\_W = GAS\_AV\_W+AREAX\*W(INODE,2)\*VFRAC(INODE,2)

AREA\_FRAC1 = AREA\_FRAC1 + AREAX\*VFRAC(INODE,1) AREA\_FRAC2 = AREA\_FRAC2 + AREAX\*VFRAC(INODE,2)

C Liquid flowrate is calculate in three faces respectively. They are C added up in final stage

X\_LIQ\_OUT=X\_LIQ\_OUT+U(INODE,1)\*AREAX\*VFRAC(INODE,1) Y\_LIQ\_OUT=Y\_LIQ\_OUT+V(INODE,1)\*AREAY\*VFRAC(INODE,1) Z\_LIQ\_OUT=Z\_LIQ\_OUT+W(INODE,1)\*AREAZ\*VFRAC(INODE,1)

GAS\_FLOWRATE=GAS\_FLOWRATE+AREAX\*U(INODE,2)\*VFRAC(INODE,2)

- +AREAY\*V(INODE,2)\*VFRAC(INODE,2) + +
  - +AREAZ\*W(INODE,2)\*VFRAC(INODE,2)

TOTALA = TOTALA+AREAX 202 CONTINUE 203 CONTINUE

IHEAD = IHEAD+ILEN

CALL IPREC('BLOCK-NUMBER-20', 'BLOCK', 'CENTRES', IPT, ILEN, JLEN,KLEN,CWORK,IWORK) +

DO 303 K=1,KLEN DO 302 I=1,ILEN INODE = IP(I,8,K)IFACE = INODE+NCELL

С

C Record the variables

```
TMP(KLEN+1-K,IHEAD+I,1) = U(INODE,1)
TMP(KLEN+1-K,IHEAD+I,2) = V(INODE,1)
TMP(KLEN+1-K,IHEAD+I,3) = W(INODE,1)
TMP(KLEN+1-K,IHEAD+I,4) = TE(INODE,1)
TMP(KLEN+1-K,IHEAD+I,5) = ED(INODE,1)
TMP(KLEN+1-K,IHEAD+I,6) = U(INODE,2)
TMP(KLEN+1-K,IHEAD+I,7) = V(INODE,2)
TMP(KLEN+1-K,IHEAD+I,8) = W(INODE,2)
TMP(KLEN+1-K,IHEAD+I,9) = VFRAC(INODE,1)
TMP(KLEN+1-K,IHEAD+I,10) = VFRAC(INODE,2)
```

```
C End of recording variables
```

C VEL is speed at current node

```
VEL = SQRT(U(INODE,1)*U(INODE,1)+
       V(INODE,1)*V(INODE,1)+W(INODE,1)*W(INODE,1))
+
  AREAX = AREA(IFACE, 1)
  AREAY = AREA(IFACE, 2)
  AREAZ = AREA(IFACE,3)
```

C Following average variable is calculated according to Hua's Formula

TOTAL\_BOOT=TOTAL\_BOOT+VEL\*AREAX\*VFRAC(INODE,1)

TOT\_UP\_U=TOT\_UP\_U+U(INODE,1)\*VEL\*AREAX\*VFRAC(INODE,1) TOT\_UP\_V=TOT\_UP\_V+V(INODE,1)\*VEL\*AREAX\*VFRAC(INODE,1) TOT\_UP\_W=TOT\_UP\_W+W(INODE,1)\*VEL\*AREAX\*VFRAC(INODE,1) TOT\_UP\_K=TOT\_UP\_K+TE(INODE,1)\*VEL\*AREAX\*VFRAC(INODE,1) TOT\_UP\_E=TOT\_UP\_E+ED(INODE,1)\*VEL\*AREAX\*VFRAC(INODE,1)

- C End of Hua's Average variables
- C Followinf variables are based on area weighted average

 $TOT_AV_U = TOT_AV_U + AREAX^*U(INODE,1)^*VFRAC(INODE,1)$  $TOT_AV_V = TOT_AV_V + AREAX^*V(INODE,1)^*VFRAC(INODE,1)$  $TOT_AV_W = TOT_AV_W + AREAX^*W(INODE,1)^*VFRAC(INODE,1)$  $TOT_AV_K = TOT_AV_K + AREAX^*TE(INODE,1)^*VFRAC(INODE,1)$  $TOT_AV_E = TOT_AV_E + AREAX^*ED(INODE,1)^*VFRAC(INODE,1)$ 

C End of area weighted average

GAS\_AV\_U = GAS\_AV\_U+AREAX\*U(INODE,2)\*VFRAC(INODE,2) GAS\_AV\_V = GAS\_AV\_V+AREAX\*V(INODE,2)\*VFRAC(INODE,2) GAS\_AV\_W = GAS\_AV\_W+AREAX\*W(INODE,2)\*VFRAC(INODE,2)

AREA\_FRAC1 = AREA\_FRAC1 + AREAX\*VFRAC(INODE,1) AREA\_FRAC2 = AREA\_FRAC2 + AREAX\*VFRAC(INODE,2)

C Liquid flowrate is calculate in three faces respectively. They are C added up in final stage

X\_LIQ\_OUT=X\_LIQ\_OUT+U(INODE,1)\*AREAX\*VFRAC(INODE,1) Y\_LIQ\_OUT=Y\_LIQ\_OUT+V(INODE,1)\*AREAY\*VFRAC(INODE,1) Z\_LIQ\_OUT=Z\_LIQ\_OUT+W(INODE,1)\*AREAZ\*VFRAC(INODE,1)

 $\mathsf{GAS\_FLOWRATE=GAS\_FLOWRATE+AREAX^*U}(\mathsf{INODE},2)^*\mathsf{VFRAC}(\mathsf{INODE},2)$ 

- + +AREAY\*V(INODE,2)\*VFRAC(INODE,2)
- + +AREAZ\*W(INODE,2)\*VFRAC(INODE,2)

TOTALA = TOTALA+AREAX 302 CONTINUE 303 CONTINUE

IHEAD = IHEAD+ILEN

CALL IPREC('BLOCK-NUMBER-21', 'BLOCK', 'CENTRES', IPT, ILEN,

► JLEN,KLEN,CWORK,IWORK)

DO 403 K=1,KLEN DO 402 I=1,ILEN INODE = IP(I,8,K) IFACE = INODE+NCELL

С

C Record the variables

$$\begin{split} \mathsf{TMP}(\mathsf{KLEN+1-K},\mathsf{IHEAD+I},1) &= \mathsf{U}(\mathsf{INODE},1) \\ \mathsf{TMP}(\mathsf{KLEN+1-K},\mathsf{IHEAD+I},2) &= \mathsf{V}(\mathsf{INODE},1) \\ \mathsf{TMP}(\mathsf{KLEN+1-K},\mathsf{IHEAD+I},3) &= \mathsf{W}(\mathsf{INODE},1) \\ \mathsf{TMP}(\mathsf{KLEN+1-K},\mathsf{IHEAD+I},4) &= \mathsf{TE}(\mathsf{INODE},1) \\ \mathsf{TMP}(\mathsf{KLEN+1-K},\mathsf{IHEAD+I},5) &= \mathsf{ED}(\mathsf{INODE},1) \\ \mathsf{TMP}(\mathsf{KLEN+1-K},\mathsf{IHEAD+I},6) &= \mathsf{U}(\mathsf{INODE},2) \\ \mathsf{TMP}(\mathsf{KLEN+1-K},\mathsf{IHEAD+I},7) &= \mathsf{V}(\mathsf{INODE},2) \end{split}$$

$$\begin{split} \mathsf{TMP}(\mathsf{KLEN+1}\text{-}\mathsf{K},\mathsf{IHEAD}\text{+}\mathsf{I},\mathsf{8}) &= \mathsf{W}(\mathsf{INODE},2)\\ \mathsf{TMP}(\mathsf{KLEN}\text{+}\mathsf{1}\text{-}\mathsf{K},\mathsf{IHEAD}\text{+}\mathsf{I},9) &= \mathsf{VFRAC}(\mathsf{INODE},1)\\ \mathsf{TMP}(\mathsf{KLEN}\text{+}\mathsf{1}\text{-}\mathsf{K},\mathsf{IHEAD}\text{+}\mathsf{I},\mathsf{1}0) &= \mathsf{VFRAC}(\mathsf{INODE},2) \end{split}$$

- C End of recording variables
- C VEL is speed at current node
  - VEL = SQRT(U(INODE,1)\*U(INODE,1)+ + V(INODE,1)\*V(INODE,1)+W(INODE,1)\*W(INODE,1)) AREAX = AREA(IFACE,1) AREAY = AREA(IFACE,2) AREAZ = AREA(IFACE,3)
- C Following average variable is calculated according to Hua's Formula

TOTAL\_BOOT=TOTAL\_BOOT+VEL\*AREAX\*VFRAC(INODE,1)

TOT\_UP\_U=TOT\_UP\_U+U(INODE,1)\*VEL\*AREAX\*VFRAC(INODE,1) TOT\_UP\_V=TOT\_UP\_V+V(INODE,1)\*VEL\*AREAX\*VFRAC(INODE,1) TOT\_UP\_W=TOT\_UP\_W+W(INODE,1)\*VEL\*AREAX\*VFRAC(INODE,1) TOT\_UP\_K=TOT\_UP\_K+TE(INODE,1)\*VEL\*AREAX\*VFRAC(INODE,1) TOT\_UP\_E=TOT\_UP\_E+ED(INODE,1)\*VEL\*AREAX\*VFRAC(INODE,1)

- C End of Hua's Average variables
- C Followinf variables are based on area weighted average

 $TOT_AV_U = TOT_AV_U + AREAX^*U(INODE,1)^*VFRAC(INODE,1)$  $TOT_AV_V = TOT_AV_V + AREAX^*V(INODE,1)^*VFRAC(INODE,1)$  $TOT_AV_W = TOT_AV_W + AREAX^*W(INODE,1)^*VFRAC(INODE,1)$  $TOT_AV_K = TOT_AV_K + AREAX^*TE(INODE,1)^*VFRAC(INODE,1)$  $TOT_AV_E = TOT_AV_E + AREAX^*ED(INODE,1)^*VFRAC(INODE,1)$ 

C End of area weighted average

 $\begin{array}{l} \mathsf{GAS}\_\mathsf{AV}\_\mathsf{U} = \mathsf{GAS}\_\mathsf{AV}\_\mathsf{U} + \mathsf{AREAX}^*\mathsf{U}(\mathsf{INODE},2)^*\mathsf{VFRAC}(\mathsf{INODE},2) \\ \mathsf{GAS}\_\mathsf{AV}\_\mathsf{V} = \mathsf{GAS}\_\mathsf{AV}\_\mathsf{V} + \mathsf{AREAX}^*\mathsf{V}(\mathsf{INODE},2)^*\mathsf{VFRAC}(\mathsf{INODE},2) \\ \mathsf{GAS}\_\mathsf{AV}\_\mathsf{W} = \mathsf{GAS}\_\mathsf{AV}\_\mathsf{W} + \mathsf{AREAX}^*\mathsf{W}(\mathsf{INODE},2)^*\mathsf{VFRAC}(\mathsf{INODE},2) \\ \end{array}$ 

AREA\_FRAC1 = AREA\_FRAC1 + AREAX\*VFRAC(INODE,1) AREA\_FRAC2 = AREA\_FRAC2 + AREAX\*VFRAC(INODE,2)

C Liquid flowrate is calculate in three faces respectively. They are C added up in final stage

X\_LIQ\_OUT=X\_LIQ\_OUT+U(INODE,1)\*AREAX\*VFRAC(INODE,1) Y\_LIQ\_OUT=Y\_LIQ\_OUT+V(INODE,1)\*AREAY\*VFRAC(INODE,1) Z\_LIQ\_OUT=Z\_LIQ\_OUT+W(INODE,1)\*AREAZ\*VFRAC(INODE,1)

GAS\_FLOWRATE=GAS\_FLOWRATE+AREAX\*U(INODE,2)\*VFRAC(INODE,2) + +AREAY\*V(INODE,2)\*VFRAC(INODE,2)

+AREAZ\*W(INODE,2)\*VFRAC(INODE,2)

TOTALA = TOTALA+AREAX 402 CONTINUE 403 CONTINUE

IHEAD =IHEAD+ILEN

+

CALL IPREC('BLOCK-NUMBER-22', 'BLOCK', 'CENTRES', IPT, ILEN,

+ JLEN,KLEN,CWORK,IWORK)

```
DO 503 K=1,KLEN
DO 502 I=1,ILEN
INODE = IP(I,8,K)
IFACE = INODE+NCELL
```

C Record the variables

С

```
\begin{split} \mathsf{TMP}(\mathsf{KLEN+1-K},\mathsf{IHEAD+I},1) &= \mathsf{U}(\mathsf{INODE},1)\\ \mathsf{TMP}(\mathsf{KLEN+1-K},\mathsf{IHEAD+I},2) &= \mathsf{V}(\mathsf{INODE},1)\\ \mathsf{TMP}(\mathsf{KLEN+1-K},\mathsf{IHEAD+I},3) &= \mathsf{W}(\mathsf{INODE},1)\\ \mathsf{TMP}(\mathsf{KLEN+1-K},\mathsf{IHEAD+I},4) &= \mathsf{TE}(\mathsf{INODE},1)\\ \mathsf{TMP}(\mathsf{KLEN+1-K},\mathsf{IHEAD+I},5) &= \mathsf{ED}(\mathsf{INODE},2)\\ \mathsf{TMP}(\mathsf{KLEN+1-K},\mathsf{IHEAD+I},6) &= \mathsf{U}(\mathsf{INODE},2)\\ \mathsf{TMP}(\mathsf{KLEN+1-K},\mathsf{IHEAD+I},7) &= \mathsf{V}(\mathsf{INODE},2)\\ \mathsf{TMP}(\mathsf{KLEN+1-K},\mathsf{IHEAD+I},8) &= \mathsf{W}(\mathsf{INODE},2)\\ \mathsf{TMP}(\mathsf{KLEN+1-K},\mathsf{IHEAD+I},9) &= \mathsf{VFRAC}(\mathsf{INODE},1)\\ \mathsf{TMP}(\mathsf{KLEN+1-K},\mathsf{IHEAD+I},10) &= \mathsf{VFRAC}(\mathsf{INODE},2) \end{split}
```

C End of recording variables

C VEL is speed at current node

```
VEL = SQRT(U(INODE,1)*U(INODE,1)+
+ V(INODE,1)*V(INODE,1)+W(INODE,1)*W(INODE,1))
AREAX = AREA(IFACE,1)
AREAY = AREA(IFACE,2)
AREAZ = AREA(IFACE,3)
```

C Following average variable is calculated according to Hua's Formula

TOTAL\_BOOT=TOTAL\_BOOT+VEL\*AREAX\*VFRAC(INODE,1)

TOT\_UP\_U=TOT\_UP\_U+U(INODE,1)\*VEL\*AREAX\*VFRAC(INODE,1) TOT\_UP\_V=TOT\_UP\_V+V(INODE,1)\*VEL\*AREAX\*VFRAC(INODE,1) TOT\_UP\_W=TOT\_UP\_W+W(INODE,1)\*VEL\*AREAX\*VFRAC(INODE,1) TOT\_UP\_K=TOT\_UP\_K+TE(INODE,1)\*VEL\*AREAX\*VFRAC(INODE,1) TOT\_UP\_E=TOT\_UP\_E+ED(INODE,1)\*VEL\*AREAX\*VFRAC(INODE,1)

C End of Hua's Average variables

C Followinf variables are based on area weighted average TOT\_AV\_U = TOT\_AV\_U+AREAX\*U(INODE,1)\*VFRAC(INODE,1) TOT\_AV\_V = TOT\_AV\_V+AREAX\*V(INODE,1)\*VFRAC(INODE,1) TOT\_AV\_W = TOT\_AV\_W+AREAX\*W(INODE,1)\*VFRAC(INODE,1) TOT\_AV\_K = TOT\_AV\_K+AREAX\*TE(INODE,1)\*VFRAC(INODE,1) TOT\_AV\_E = TOT\_AV\_E+AREAX\*E0(INODE,1)\*VFRAC(INODE,1)

C End of area weighted average

 $\begin{array}{l} \mathsf{GAS\_AV\_U} = \; \mathsf{GAS\_AV\_U} + \mathsf{AREAX^*U}(\mathsf{INODE,2})^*\mathsf{VFRAC}(\mathsf{INODE,2}) \\ \mathsf{GAS\_AV\_V} = \; \mathsf{GAS\_AV\_V} + \mathsf{AREAX^*V}(\mathsf{INODE,2})^*\mathsf{VFRAC}(\mathsf{INODE,2}) \\ \mathsf{GAS\_AV\_W} = \; \mathsf{GAS\_AV\_W} + \mathsf{AREAX^*W}(\mathsf{INODE,2})^*\mathsf{VFRAC}(\mathsf{INODE,2}) \\ \end{array}$ 

AREA\_FRAC1 = AREA\_FRAC1 + AREAX\*VFRAC(INODE,1) AREA\_FRAC2 = AREA\_FRAC2 + AREAX\*VFRAC(INODE,2)

C Liquid flowrate is calculate in three faces respectively. They are C added up in final stage

X\_LIQ\_OUT=X\_LIQ\_OUT+U(INODE,1)\*AREAX\*VFRAC(INODE,1) Y\_LIQ\_OUT=Y\_LIQ\_OUT+V(INODE,1)\*AREAY\*VFRAC(INODE,1)

#### Z\_LIQ\_OUT=Z\_LIQ\_OUT+W(INODE,1)\*AREAZ\*VFRAC(INODE,1)

GAS\_FLOWRATE=GAS\_FLOWRATE+AREAX\*U(INODE,2)\*VFRAC(INODE,2)

- + +AREAY\*V(INODE,2)\*VFRAC(INODE,2)
  - +AREAZ\*W(INODE,2)\*VFRAC(INODE,2)

TOTALA = TOTALA+AREAX 502 CONTINUE 503 CONTINUE

> IHEAD = IHEAD+ILEN KHEAD = KLEN

С

+

AVERAGE\_U=TOT\_UP\_U/TOTAL\_BOOT AVERAGE\_V=TOT\_UP\_V/TOTAL\_BOOT AVERAGE\_W=TOT\_UP\_W/TOTAL\_BOOT AVERAGE\_K=TOT\_UP\_K/TOTAL\_BOOT AVERAGE\_E=TOT\_UP\_E/TOTAL\_BOOT

AREA\_AV\_FRAC1 = AREA\_FRAC1/TOTALA AREA\_AV\_FRAC2 = AREA\_FRAC2/TOTALA AREA\_AV\_U=TOT\_AV\_U/AREA\_FRAC1 AREA\_AV\_V=TOT\_AV\_V/AREA\_FRAC1 AREA\_AV\_W=TOT\_AV\_W/AREA\_FRAC1 AREA\_AV\_K=TOT\_AV\_K/AREA\_FRAC1 AREA\_AV\_E=TOT\_AV\_E/AREA\_FRAC1

AREA\_GAS\_U=GAS\_AV\_U/AREA\_FRAC2 AREA\_GAS\_V=GAS\_AV\_V/AREA\_FRAC2 AREA\_GAS\_W=GAS\_AV\_W/AREA\_FRAC2

```
FLOWRATE_LIQ = 2*(X_LIQ_OUT+Y_LIQ_OUT+Z_LIQ_OUT)
GAS_FLOWRATE = 2*GAS_FLOWRATE
GAS_FRAC=100*GAS_FLOWRATE/(GAS_FLOWRATE+FLOWRATE_LIQ)
```

Ver = AVERAGE\_W/AVERAGE\_U Hor = AVERAGE\_V/AVERAGE\_U

V\_JET\_ANGLE = ATAN(Ver)\*180/3.14159 H\_JET\_ANGLE = ATAN(Hor)\*180/3.14159

```
OPEN(6.FILE='DATA2.DAT'.STATUS='NEW')
WRITE(6,*) 'AVERAGE LIQUID U ON THE PORT', AVERAGE_U
WRITE(6,*) ' '
WRITE(6,*) 'AVERAGE LIQUID V ON THE PORT', AVERAGE_V
WRITE(6,*) ' '
WRITE(6,*) 'AVERAGE LIQUID W ON THE PORT', AVERAGE_W
WRITE(6,*) ' '
WRITE(6,*) 'AVERAGE K ON THE PORT
                                  ',AVERAGE_K
WRITE(6,*) ' '
WRITE(6,*) 'AVERAGE EPSILON ON THE PORT ', AVERAGE_E
WRITE(6,*) ' '
WRITE(6,*) 'VERTICAL JET ANGLE=
                                 ',V_JET_ANGLE
WRITE(6,*) ' '
WRITE(6,*) 'HORIZONTAL JET ANGLE=
                                  ',H_JET_ANGLE
WRITE(6,*) ' '
WRITE(6,*) 'TOTAL AREA OF PORT =',TOTALA
WRITE(6,*) 'LIQUID FLOWRATE AT X AXIS=',X_LIQ_OUT
```

```
WRITE(6,*) 'LIQUID FLOWRATE AT Y AXIS=',Y_LIQ_OUT
   WRITE(6,*) 'LIQUID FLOWRATE AT Z AXIS=',Z_LIQ_OUT
   WRITE(6,*) 'LIQUID FLOWRATE AT PORT =',FLOWRATE_LIQ
   WRITE(6,*) 'GAS FLOWRATE AT PORT =',GAS_FLOWRATE
   WRITE(6,*) 'GAS VOLUME FRACTION IN OUTFLOW AT PORT(%)=',GAS_FRAC
   WRITE(6,*) 'LIQUID AREA WEIGHTED AVERAGE U', AREA_AV_U
   WRITE(6,*) 'LIQUID AREA WEIGHTED AVERAGE V', AREA_AV_V
   WRITE(6,*) 'LIQUID AREA WEIGHTED AVERAGE W', AREA_AV_W
   WRITE(6,*) 'LIQUID AREA WEIGHTED AVERAGE VOL.FRAC', AREA_AV_FRAC1
   WRITE(6,*) 'LIQUID AREA WEIGHTED AVERAGE K', AREA_AV_K
   WRITE(6,*) 'LIQUID AREA WEIGHTED AVERAGE EPSILON', AREA_AV_E
   WRITE(6,*) 'GAS AREA WEIGHTED AVERAGE U', AREA_GAS_U
   WRITE(6,*) 'GAS AREA WEIGHTED AVERAGE V', AREA GAS V
   WRITE(6,*) 'GAS AREA WEIGHTED AVERAGE W', AREA GAS W
   WRITE(6,*) 'GAS AREA WEIGHTED AVERAGE VOL.FRAC', AREA_AV_FRAC2
   WRITE(6,*) ' '
   WRITE(6,*) 'K J
                     U
                            V
                                   W'
   DO 703 K=1,KHEAD
   DO 702 J=1,IHEAD
    WRITE(6,*) K,J,TMP(K,J,1),TMP(K,J,2),TMP(K,J,3)
702 CONTINUE
703 CONTINUE
  CLOSE(6)
       WRITE(5,*) KHEAD, IHEAD
  DO 603 K=1,KHEAD
   DO 602 J=1,IHEAD
    DO 601 I=1,10
        WRITE(5,*) TMP(K,J,I)
601 CONTINUE
602 CONTINUE
603 CONTINUE
  CLOSE(5)
С
С
С
  ENDIF
С
С
  RETURN
  END
  SUBROUTINE USRBCS(VARBCS, VARAMB, A, B, C, ACND, BCND, CCND
            ,IWGVEL,NDVWAL
  +
            ,FLOUT,NLABEL,NSTART,NEND,NCST,NCEN
  +
  +
            ,U,V,W,P,VFRAC,DEN,VIS,TE,ED,RS,T,H,RF,SCAL
            ,XP,YP,ZP,VOL,AREA,VPOR,ARPOR,WFACT,IPT
  +
            ,IBLK,IPVERT,IPNODN,IPFACN,IPNODF,IPNODB,IPFACB
  +
            ,WORK,IWORK,CWORK)
  +
С
       C****
С
```
```
C USER ROUTINE TO SET REALS AT BOUNDARIES.
С
С
 >>> IMPORTANT
                                           <<<
С
                                       <<<
  >>>
C >>> USERS MAY ONLY ADD OR ALTER PARTS OF THE SUBROUTINE WITHIN <<<
С
 >>> THE DESIGNATED USER AREAS
                                                 <<<
С
С
C THIS SUBROUTINE IS CALLED BY THE FOLLOWING SUBROUTINE
  CUSR SRLIST
С
С
C CREATED
С
     30/11/88 ADB
C MODIFIED
С
     08/09/90 ADB RESTRUCTURED FOR USER-FRIENDLINESS.
С
     10/08/91 IRH FURTHER RESTRUCTURING ADD ACND BCND CCND
С
     22/09/91 IRH CHANGE ICALL TO IUCALL + ADD /SPARM/
С
     10/03/92 PHA UPDATE CALLED BY COMMENT, ADD RF ARGUMENT,
С
           CHANGE LAST DIMENSION OF RS TO 6 AND IVERS TO 2
С
     03/06/92 PHA ADD PRECISION FLAG AND CHANGE IVERS TO 3
С
     30/06/92 NSW INCLUDE FLAG FOR CALLING BY ITERATION
С
           INSERT EXTRA COMMENTS
С
     03/08/92 NSW MODIFY DIMENSION STATEMENTS FOR VAX
С
     21/12/92 CSH INCREASE IVERS TO 4
С
     02/08/93 NSW INCORRECT AND MISLEADING COMMENT REMOVED
С
     05/11/93 NSW INDICATE USE OF FLOUT IN MULTIPHASE FLOWS
С
     23/11/93 CSH EXPLICITLY DIMENSION IPVERT ETC.
С
     01/02/94 NSW SET VARIABLE POINTERS IN WALL EXAMPLE.
С
           CHANGE FLOW3D TO CFDS-FLOW3D.
С
           MODIFY MULTIPHASE MASS FLOW BOUNDARY TREATMENT.
С
     03/03/94 FHW CORRECTION OF SPELLING MISTAKE
С
     02/07/94 BAS SLIDING GRIDS - ADD NEW ARGUMENT IWGVEL
С
           TO ALLOW VARIANTS OF TRANSIENT-GRID WALL BC
С
           CHANGE VERSION NUMBER TO 5
С
     09/08/94 NSW CORRECT SPELLING
С
           MOVE 'IF(IUSED.EQ.0) RETURN' OUT OF USER AREA
С
     19/12/94 NSW CHANGE FOR CFX-F3D
С
     02/02/95 NSW CHANGE COMMON /IMFBMP/
С
     02/06/97 NSW MAKE EXAMPLE MORE LOGICAL
С
     02/07/97 NSW UPDATE FOR CFX-4
С
     08/09/98 NSW CORRECT SIZE OF WALL ARRAY IN COMMENT
С
С
C SUBROUTINE ARGUMENTS
С
С
   VARBCS - REAL BOUNDARY CONDITIONS
С
   VARAMB - AMBIENT VALUE OF VARIABLES
С
       - COEFFICIENT IN WALL BOUNDARY CONDITION
   Α
С
       - COEFFICIENT IN WALL BOUNDARY CONDITION
   в
С
   С
      - COEFFICIENT IN WALL BOUNDARY CONDITION
   ACND - COEFFICIENT IN CONDUCTING WALL BOUNDARY CONDITION
С
С
   BCND - COEFFICIENT IN CONDUCTING WALL BOUNDARY CONDITION
С
   CCND - COEFFICIENT IN CONDUCTING WALL BOUNDARY CONDITION
   IWGVEL - USAGE OF INPUT VELOCITIES (0 = AS IS,1 = ADD GRID MOTION)
С
С
   NDVWAL - FIRST DIMENSION OF ARRAY IWGVEL
С
   FLOUT - MASS FLOW/FRACTIONAL MASS FLOW
С
   NLABEL - NUMBER OF DISTINCT OUTLETS
```

C NSTART - ARRAY POINTER С NEND - ARRAY POINTER NCST - ARRAY POINTER С С NCEN - ARRAY POINTER С U - U COMPONENT OF VELOCITY С V - V COMPONENT OF VELOCITY С W - W COMPONENT OF VELOCITY - PRESSURE С Р С **VFRAC - VOLUME FRACTION** С DEN - DENSITY OF FLUID С VIS - VISCOSITY OF FLUID С TE - TURBULENT KINETIC ENERGY С ED - EPSILON С RS - REYNOLD STRESSES С Т - TEMPERATURE С Н - ENTHALPY С RF - REYNOLD FLUXES С SCAL - SCALARS (THE FIRST 'NCONC' OF THESE ARE MASS FRACTIONS) С XP - X COORDINATES OF CELL CENTRES С YΡ - Y COORDINATES OF CELL CENTRES С ZP - Z COORDINATES OF CELL CENTRES С VOL - VOLUME OF CELLS С AREA - AREA OF CELLS С VPOR - POROUS VOLUME С ARPOR - POROUS AREA С WFACT - WEIGHT FACTORS С С IPT - 1D POINTER ARRAY С IBLK - BLOCK SIZE INFORMATION C IPVERT - POINTER FROM CELL CENTERS TO 8 NEIGHBOURING VERTICES С IPNODN - POINTER FROM CELL CENTERS TO 6 NEIGHBOURING CELLS С IPFACN - POINTER FROM CELL CENTERS TO 6 NEIGHBOURING FACES С IPNODF - POINTER FROM CELL FACES TO 2 NEIGHBOURING CELL CENTERS С IPNODB - POINTER FROM BOUNDARY CENTERS TO CELL CENTERS С IPFACB - POINTER TO NODES FROM BOUNDARY FACES С С WORK - REAL WORKSPACE ARRAY С IWORK - INTEGER WORKSPACE ARRAY С CWORK - CHARACTER WORKSPACE ARRAY С C SUBROUTINE ARGUMENTS PRECEDED WITH A '\*' ARE ARGUMENTS THAT MUST C BE SET BY THE USER IN THIS ROUTINE. С C NOTE THAT OTHER DATA MAY BE OBTAINED FROM CFX-4 USING THE C ROUTINE GETADD, FOR FURTHER DETAILS SEE THE VERSION 4 C USER MANUAL. С LOGICAL LDEN, LVIS, LTURB, LTEMP, LBUOY, LSCAL, LCOMP ,LRECT,LCYN,LAXIS,LPOROS,LTRANS + С CHARACTER\*(\*) CWORK С C---- AREA FOR USERS EXPLICITLY DECLARED VARIABLES С С COMMON + /ALL/ NBLOCK,NCELL,NBDRY,NNODE,NFACE,NVERT,NDIM

```
,NDVAR,NDPROP,NDXNN,NDGEOM,NDCOEF,NILIST,NRLIST,NTOPOL
  + /BCSOUT/ IFLOUT
  + /CHKUSR/ IVERS,IUCALL,IUSED
  + /DEVICE/ NREAD,NWRITE,NRDISK,NWDISK
  + /IDUM/ ILEN,JLEN
  + /IMFBMP/ IMFBMP,JMFBMP
  + /LOGIC/ LDEN,LVIS,LTURB,LTEMP,LBUOY,LSCAL,LCOMP
       ,LRECT,LCYN,LAXIS,LPOROS,LTRANS
  + /MLTGRD/ MLEVEL,NLEVEL,ILEVEL
  + /SGLDBL/ IFLGPR,ICHKPR
  + /SPARM/ SMALL,SORMAX,NITER,INDPRI,MAXIT,NODREF,NODMON
  + /TRANSI/ NSTEP,KSTEP,MF,INCORE
  + /TRANSR/ TIME.DT.DTINVF.TPARM
  + /UBCSFL/ IUBCSF
С
C---- AREA FOR USERS TO DECLARE THEIR OWN COMMON BLOCKS
C THESE SHOULD START WITH THE CHARACTERS 'UC' TO ENSURE
С
  NO CONFLICT WITH NON-USER COMMON BLOCKS
С
С
  DIMENSION
 + VARBCS(NVAR,NPHASE,NCELL+1:NNODE),VARAMB(NVAR,NPHASE)
  +,A(4+NSCAL,NPHASE,NSTART:*)
  +,B(4+NSCAL,NPHASE,NSTART:*),C(4+NSCAL,NPHASE,NSTART:*)
  +,FLOUT(*),ACND(NCST:*),BCND(NCST:*),CCND(NCST:*)
  +,IWGVEL(NDVWAL,NPHASE)
С
  DIMENSION
 + U(NNODE,NPHASE),V(NNODE,NPHASE),W(NNODE,NPHASE),P(NNODE,NPHASE)
 +,VFRAC(NNODE,NPHASE),DEN(NNODE,NPHASE),VIS(NNODE,NPHASE)
 +,TE(NNODE,NPHASE),ED(NNODE,NPHASE),RS(NNODE,NPHASE,6)
 +,T(NNODE,NPHASE),H(NNODE,NPHASE),RF(NNODE,NPHASE,4)
 +,SCAL(NNODE,NPHASE,NSCAL)
С
  DIMENSION
 + XP(NNODE), YP(NNODE), ZP(NNODE)
 +,VOL(NCELL),AREA(NFACE,3),VPOR(NCELL),ARPOR(NFACE,3),WFACT(NFACE)
  +.IPT(*).IBLK(5.NBLOCK)
 +,IPVERT(NCELL,8),IPNODN(NCELL,6),IPFACN(NCELL,6),IPNODF(NFACE,4)
 +, IPNODB(NBDRY, 4), IPFACB(NBDRY)
  +,IWORK(*),WORK(*),CWORK(*)
С
C---- AREA FOR USERS TO DIMENSION THEIR ARRAYS
С
C---- AREA FOR USERS TO DEFINE DATA STATEMENTS
С
С
C---- STATEMENT FUNCTION FOR ADDRESSING
   IP(I,J,K)=IPT((K-1)*ILEN*JLEN+(J-1)*ILEN+I)
С
C----VERSION NUMBER OF USER ROUTINE AND PRECISION FLAG
С
  IVERS=5
  ICHKPR = 1
```

+ /ALLWRK/ NRWS,NIWS,NCWS,IWRFRE,IWIFRE,IWCFRE

+ /ADDIMS/ NPHASE,NSCAL,NVAR,NPROP

С C---- TO USE THIS USER ROUTINE FIRST SET IUSED=1 C AND SET IUBCSF FLAG: C BOUNDARY CONDITIONS NOT CHANGING IUBCSF=0 C BOUNDARY CONDITIONS CHANGING WITH ITERATION IUBCSF=1 C BOUNDARY CONDITIONS CHANGING WITH TIME IUBCSF=2 С BOUNDARY CONDITIONS CHANGING WITH TIME AND ITERATION IUBCSF=3 С IUSED=1 IUBCSF=0 С С IF (IUSED.EQ.0) RETURN С C---- FRONTEND CHECKING OF USER ROUTINE IF (IUCALL.EQ.0) RETURN С С C---- AREA FOR SETTING VALUES AT INLETS, PRESSURE BOUNDARIES С AND OUTLETS. (NOTE THAT THE MASS FLOW AT OUTLETS IS С SPECIFIED IN USER AREA 7) С CALL GETVAR('USRBCS','P ',KP) С CALL IPREC('RIGHTPORT', 'PATCH', 'CENTRES', IPT, ILEN, JLEN, KLEN, CWORK, IWORK) + С ZMAX=0.0828 ZMIN=0.0048 PMAX=16717.43 PMIN=11351.34 C LOOP OVER PATCH DO 103 K = 1, KLEN DO 102 J = 1, JLEN DO 101 I = 1, ILEN C USE STATEMENT FUNCTION IP TO GET ADDRESSES INODE = IP(I,J,K)C SET VARBCS F=(ZMAX-ZP(INODE))/(ZMAX-ZMIN) VARBCS(KP,1,INODE) = F\*PMAX + (1.0-F)\*PMIN VARBCS(KP,2,INODE) = F\*PMAX + (1.0-F)\*PMIN 101 CONTINUE 102 CONTINUE 103 CONTINUE С C----END OF EXAMPLE С С RETURN END

## Appendix F. Subroutine for multiphase MUSIG model using output of nozzle simulation as input conditions

SUBROUTINE USRSRC(IEQN,ICALL,CNAME,CALIAS,AM,SP,SU,CONV ,U,V,W,P,VFRAC,DEN,VIS,TE,ED,RS,T,H,RF,SCAL ,XP,YP,ZP,VOL,AREA,VPOR,ARPOR,WFACT,IPT + ,IBLK,IPVERT,IPNODN,IPFACN,IPNODF,IPNODB,IPFACB + ,WORK,IWORK,CWORK) С LOGICAL LDEN, LVIS, LTURB, LTEMP, LBUOY, LSCAL, LCOMP ,LRECT,LCYN,LAXIS,LPOROS,LTRANS С CHARACTER\*(\*) CWORK CHARACTER CNAME\*6, CALIAS\*24 С С COMMON + /ALL/ NBLOCK,NCELL,NBDRY,NNODE,NFACE,NVERT,NDIM + /ALLWRK/ NRWS,NIWS,NCWS,IWRFRE,IWIFRE,IWCFRE + /ADDIMS/ NPHASE,NSCAL,NVAR,NPROP ,NDVAR,NDPROP,NDXNN,NDGEOM,NDCOEF,NILIST,NRLIST,NTOPOL + /CHKUSR/ IVERS,IUCALL,IUSED + /DEVICE/ NREAD,NWRITE,NRDISK,NWDISK + /IDUM/ ILEN,JLEN + /LOGIC/ LDEN,LVIS,LTURB,LTEMP,LBUOY,LSCAL,LCOMP ,LRECT,LCYN,LAXIS,LPOROS,LTRANS + /MLTGRD/ MLEVEL, NLEVEL, ILEVEL + /SGLDBL/ IFLGPR,ICHKPR + /SPARM/ SMALL,SORMAX,NITER,INDPRI,MAXIT,NODREF,NODMON + /TRANSI/ NSTEP,KSTEP,MF,INCORE + /TRANSR/ TIME, DT, DTINVF, TPARM С С COMMON /UCSURF/ JISURF,NISURF,JDSURF,NDSURF COMMON /UCINJR/ FINJ1, FINJ2, UGASIN, VGASIN, WGASIN COMMON /UCFLUX/ GASIN, GASOUT С С DIMENSION AM(NCELL,6,NPHASE),SP(NCELL,NPHASE),SU(NCELL,NPHASE) +,CONV(NFACE,NPHASE) С DIMENSION + U(NNODE,NPHASE),V(NNODE,NPHASE),W(NNODE,NPHASE),P(NNODE,NPHASE) +,VFRAC(NNODE,NPHASE),DEN(NNODE,NPHASE),VIS(NNODE,NPHASE) +,TE(NNODE,NPHASE),ED(NNODE,NPHASE),RS(NNODE,NPHASE,6) +,T(NNODE,NPHASE),H(NNODE,NPHASE),RF(NNODE,NPHASE,4) +,SCAL(NNODE,NPHASE,NSCAL) С DIMENSION + XP(NNODE), YP(NNODE), ZP(NNODE) +,VOL(NCELL),AREA(NFACE,3),VPOR(NCELL),ARPOR(NFACE,3) +,WFACT(NFACE) +,IPT(\*),IBLK(5,NBLOCK) +,IPVERT(NCELL,8),IPNODN(NCELL,6),IPFACN(NCELL,6),IPNODF(NFACE,4) +,IPNODB(NBDRY,4),IPFACB(NBDRY) +,IWORK(\*),WORK(\*),CWORK(\*) С

```
C---- STATEMENT FUNCTION FOR ADRESSING
   IP(I,J,K)=IPT((K-1)*ILEN*JLEN+(J-1)*ILEN+I)
С
C----VERSION NUMBER OF USER ROUTINE AND PRECISION FLAG
С
  IVERS=4
  ICHKPR = 1
  IUSED=1
  IF (IUSED.EQ.0) RETURN
С
C---- FRONTEND CHECKING OF USER ROUTINE
   IF (IUCALL.EQ.0) RETURN
С
C---- ADD TO SOURCE TERMS
  IF (ICALL.EQ.1) THEN
С
С
C-----
C VOLUME FRACTION [M]/[T]
С
  IF(CALIAS.EQ.'VOLUME FRACTION') THEN
С
C.....SINK AT TOP
   CALL USRDEGAS(CALIAS,SP,W,VFRAC,DEN,AREA,IWORK(JDSURF),NDSURF
         ,IPNODB,IPFACB,TOTAL)
  +
   WRITE(NWRITE,*) ' VOLUME GASOUT=:',TOTAL
   WRITE(NWRITE,*) ' VF :',TOTAL
С
С
  ENDIF
С
C-----
C PRESSURE [M]/[T]
С
  IF(CALIAS.EQ.'PRESSURE') THEN
С
   GASIN = 0.0
С
C.....SINK AT TOP
   CALL USRDEGAS(CALIAS,SU,W,VFRAC,DEN,AREA,IWORK(JDSURF),NDSURF
  +
         ,IPNODB,IPFACB,TOTAL)
   GASOUT = TOTAL
   WRITE(NWRITE,*) 'PRESSURE GASOUT=:',TOTAL
С
  ENDIF
С
C-----
C UVELOCITY [M][L]/[T]^2
С
  IF(CALIAS.EQ.'U VELOCITY') THEN
С
С
C.....SINK AT TOP
   CALL USRDEGAS(CALIAS,SP,W,VFRAC,DEN,AREA,IWORK(JDSURF),NDSURF
         ,IPNODB,IPFACB,TOTAL)
  +
C.....SET MINIMUM VOID FRACTION
    VFMIN=1.0E-8
```

```
VFMIN=1.0E-8
DO 101 INODE=1,NNODE
VFRAC(INODE,2)=MAX(VFRAC(INODE,2),VFMIN)
```

```
101 CONTINUE
С
  ENDIF
С
C-
C V VELOCITY [M][L]/[T]^2
С
  IF(CALIAS.EQ.'V VELOCITY') THEN
С
С
C.....SINK AT TOP
   CALL USRDEGAS(CALIAS,SP,W,VFRAC,DEN,AREA,IWORK(JDSURF),NDSURF
         ,IPNODB,IPFACB,TOTAL)
  +
C WRITE(NWRITE,*) ' V :',TOTAL
C.....SET MINIMUM VOID FRACTION
   VFMIN=1.0E-8
   DO 201 INODE=1,NNODE
   VFRAC(INODE,2)=MAX(VFRAC(INODE,2),VFMIN)
201 CONTINUE
С
  ENDIF
С
C-----
C W VELOCITY [M][L]/[T]^2
С
  IF(CALIAS.EQ.'W VELOCITY') THEN
С
C.....SINK AT TOP
   CALL USRDEGAS(CALIAS,SP,W,VFRAC,DEN,AREA,IWORK(JDSURF),NDSURF
        ,IPNODB,IPFACB,TOTAL)
С
    WRITE(NWRITE,*) ' W :',TOTAL
С
C.....SET MINIMUM VOID FRACTION
   VFMIN=1.0E-8
   DO 301 INODE=1,NNODE
   VFRAC(INODE,2)=MAX(VFRAC(INODE,2),VFMIN)
301 CONTINUE
С
  ENDIF
С
ENDIF
C---- OVERWRITE SOURCE TERMS
  IF (ICALL.EQ.2) THEN
С
  ENDIF
  RETURN
  END
С
С
  SUBROUTINE USRDEGAS(CALIAS,SUP,W,VFRAC,DEN,AREA,IPTSRF,NDSURF
          ,IPNODB,IPFACB,TOTAL)
  +
С
  *****
C,
С
C UTILITY SUBROUTINE FOR USER-SUPPLIED FREE SURFACE DEGASSING
С
```

С C THIS SUBROUTINE IS CALLED BY THE FOLLOWING SUBROUTINES С USRSRC С CHARACTER CALIAS\*24 С C---- AREA FOR USERS EXPLICITLY DECLARED VARIABLES С С COMMON + /ALL/ NBLOCK,NCELL,NBDRY,NNODE,NFACE,NVERT,NDIM + /ALLWRK/ NRWS,NIWS,NCWS,IWRFRE,IWIFRE,IWCFRE + /ADDIMS/ NPHASE,NSCAL,NVAR,NPROP ,NDVAR,NDPROP,NDXNN,NDGEOM,NDCOEF,NILIST,NRLIST,NTOPOL + + /DEVICE/ NREAD,NWRITE,NRDISK,NWDISK С С DIMENSION SUP(NCELL,NPHASE) +,W(NNODE,NPHASE),VFRAC(NNODE,NPHASE),DEN(NNODE,NPHASE) +,AREA(NFACE,3),IPTSRF(NDSURF),IPNODB(NBDRY,4),IPFACB(NBDRY) С С C---- STATEMENT FUNCTION FOR ADRESSING С C-----DEGASSING PHASE INDEX TOTAL=0.0 C-----C VOLUME FRACTION С IF(CALIAS.EQ.'VOLUME FRACTION') THEN С DO 101 I=1,NDSURF INODE = IPTSRF(I)IBDRY = INODE-NCELL INODE1= IPNODB(IBDRY,1) IFACE = IPFACB(IBDRY)AREAM = SQRT ( AREA(IFACE,1)\*AREA(IFACE,1) +AREA(IFACE,2)\*AREA(IFACE,2) + +AREA(IFACE,3)\*AREA(IFACE,3)) + FLUX = DEN(INODE1,2)\* AREAM\*MAX(W(INODE1,2),0.0) + SUP(INODE1,2)=SUP(INODE1,2)-FLUX С **FNDIF** TOTAL = TOTAL+FLUX С WRITE(NWRITE,\*) ' FLUX=',FLUX 101 CONTINUE С C-----C OTHER VARIABLES С ELSE С DO 201 I=1,NDSURF

INODE = IPTSRF(I) IBDRY = INODE-NCELL INODE1= IPNODB(IBDRY,1)

```
IFACE = IPFACB(IBDRY)
   AREAM = SQRT ( AREA(IFACE,1)*AREA(IFACE,1)
          +AREA(IFACE,2)*AREA(IFACE,2)
  +
          +AREA(IFACE,3)*AREA(IFACE,3) )
  +
   FLUX = VFRAC(INODE1,2)*DEN(INODE1,2)
        *AREAM*MAX(W(INODE1,2),0.0)
  +
   SUP(INODE1,2)=SUP(INODE1,2)-FLUX
      TOTAL = TOTAL+FLUX
С
    WRITE(NWRITE,*) ' FLUX=',FLUX
201 CONTINUE
С
  ENDIF
401 CONTINUE
С
  RETURN
  END
С
С
  SUBROUTINE USRTRN(U,V,W,P,VFRAC,DEN,VIS,TE,ED,RS,T,H,RF,SCAL,
           XP,YP,ZP,VOL,AREA,VPOR,ARPOR,WFACT,CONV,IPT,
  +
           IBLK, IPVERT, IPNODN, IPFACN, IPNODF, IPNODB, IPFACB,
  +
           WORK, IWORK, CWORK)
  +
С
       *****
C
С
C USER SUBROUTINE TO ALLOW USERS TO MODIFY OR MONITOR THE SOLUTION AT
C THE END OF EACH TIME STEP
C THIS SUBROUTINE IS CALLED BEFORE THE START OF THE RUN AS WELL AS AT
C THE END OF EACH TIME STEP
С
  ******
C*
С
C THIS SUBROUTINE IS CALLED BY THE FOLLOWING SUBROUTINES
С
  CUSR TRNMOD
С
           *****
C***
С
C*
     *****
С
  LOGICAL LDEN, LVIS, LTURB, LTEMP, LBUOY, LSCAL, LCOMP
      ,LRECT,LCYN,LAXIS,LPOROS,LTRANS
  +
С
  CHARACTER*(*) CWORK
С
C---- AREA FOR USERS EXPLICITLY DECLARED VARIABLES
С
С
  COMMON
  + /ALL/ NBLOCK,NCELL,NBDRY,NNODE,NFACE,NVERT,NDIM
  + /ALLWRK/ NRWS,NIWS,NCWS,IWRFRE,IWIFRE,IWCFRE
  + /ADDIMS/ NPHASE,NSCAL,NVAR,NPROP
       ,NDVAR,NDPROP,NDXNN,NDGEOM,NDCOEF,NILIST,NRLIST,NTOPOL
  + /CHKUSR/ IVERS,IUCALL,IUSED
  + /CONC/ NCONC
  + /DEVICE/ NREAD,NWRITE,NRDISK,NWDISK
  + /IDUM/ ILEN,JLEN
  + /LOGIC/ LDEN,LVIS,LTURB,LTEMP,LBUOY,LSCAL,LCOMP
       ,LRECT,LCYN,LAXIS,LPOROS,LTRANS
```

```
+ /MLTGRD/ MLEVEL,NLEVEL,ILEVEL
  + /SGLDBL/ IFLGPR,ICHKPR
  + /SPARM/ SMALL,SORMAX,NITER,INDPRI,MAXIT,NODREF,NODMON
  + /TIMUSR/ DTUSR
  + /TRANSI/ NSTEP,KSTEP,MF,INCORE
  + /TRANSR/ TIME, DT, DTINVF, TPARM
С
C---- AREA FOR USERS TO DECLARE THEIR OWN COMMON BLOCKS
C THESE SHOULD START WITH THE CHARACTERS 'UC' TO ENSURE
  NO CONFLICT WITH NON-USER COMMON BLOCKS
С
С
  COMMON /UCSURF/ JISURF, NISURF, JDSURF, NDSURF
  COMMON /UCINJR/ FINJ1, FINJ2, UGASIN, VGASIN, WGASIN
  COMMON /UCFLUX/ GASIN.GASOUT
С
С
  DIMENSION
 + U(NNODE,NPHASE),V(NNODE,NPHASE),W(NNODE,NPHASE),P(NNODE,NPHASE)
 +,VFRAC(NNODE,NPHASE),DEN(NNODE,NPHASE),VIS(NNODE,NPHASE)
 +,TE(NNODE,NPHASE),ED(NNODE,NPHASE),RS(NNODE,NPHASE,6)
 +,T(NNODE,NPHASE),H(NNODE,NPHASE),RF(NNODE,NPHASE,4)
 +,SCAL(NNODE,NPHASE,NSCAL)
  DIMENSION
 + XP(NNODE), YP(NNODE), ZP(NNODE)
 +,VOL(NCELL),AREA(NFACE,3),VPOR(NCELL),ARPOR(NFACE,3)
 +,WFACT(NFACE),CONV(NFACE,NPHASE)
  +,IPT(*),IBLK(5,NBLOCK)
 +,IPVERT(NCELL,8),IPNODN(NCELL,6),IPFACN(NCELL,6),IPNODF(NFACE,4)
 +,IPNODB(NBDRY,4),IPFACB(NBDRY)
  +,IWORK(*),WORK(*),CWORK(*)
С
С
C---- STATEMENT FUNCTION FOR ADDRESSING
  IP(I,J,K)=IPT((K-1)*ILEN*JLEN+(J-1)*ILEN+I)
С
C----VERSION NUMBER OF USER ROUTINE AND PRECISION FLAG
С
  IVERS=3
  ICHKPR = 1
С
C---- TO USE THIS USER ROUTINE FIRST SET IUSED=1
С
  IUSED=1
  IF (IUSED.EQ.0) RETURN
С
C---- FRONTEND CHECKING OF USER ROUTINE
  IF (IUCALL.EQ.0) RETURN
С
С
  IF (KSTEP.EQ.0) THEN
С
C-----INJECTION BOUNDARY LOCATIONS
С
  CALL IPALL('NOZZLE INLET','*','PATCH','CENTRES',IPT,NPT,
```

+

CWORK, IWORK)

```
175
```

```
С
  NISURF=NPT
C SET INJECTION BOUNDARY LIST INTO INTERGER WORK SPACE
   CALL SETPER('USRTRN','IWORK ','ISURF ',NISURF,JISURF)
С
   DO 101 I=1,NISURF
   IWORK(JISURF+I-1)=IPT(I)
101 CONTINUE
  WRITE(NWRITE,*) ' '
  WRITE(NWRITE,*) '******* USRTRN - INLET *********
  WRITE(NWRITE,*) ' NISURF=',NISURF,' JISURF=',JISURF
  WRITE(NWRITE,*) ' IWORK(JISURF)=',IWORK(JISURF)
   WRITE(NWRITE,*) ' IWORK(JISURF+NISURF-1)=',IWORK(JISURF+NISURF-1)
С
C-----INJECTION RATES: MASS FLUX IN kg/s, VELOCITY IN m/s
   FGASIN = 0
   UGASIN = 0
   VGASIN = 0
   WGASIN = 0
С
C-----DETERMINE TOTAL AREA OF INLET BOUNDARY AND MASS FLUX PER AREA
   TOTALA = 0.0
   DO 201 I=1,NISURF
    INODE = IWORK(JISURF+I-1)
    IBDRY = INODE-NCELL
    IFACE = IPFACB(IBDRY)
    AREAM = SQRT ( AREA(IFACE,1)*AREA(IFACE,1)
   +
             +AREA(IFACE,2)*AREA(IFACE,2)
             +AREA(IFACE,3)*AREA(IFACE,3))
    TOTALA = TOTALA+AREAM
201 CONTINUE
С
  WRITE(NWRITE,*) ' TOTAL AREA OF INLET =',TOTALA
   FINJ2 = FGASIN/TOTALA
С
C-----DEGASSING BOUNDARY LOCATIONS
С
   CALL IPALL('TOP SURFACE','*','PATCH','CENTRES',IPT,NPT,
          CWORK, IWORK)
   +
С
  NDSURF=NPT
C SET DEGASSING BOUNDARY LIST INTO INTERGER WORK SPACE
   CALL SETPER('USRTRN','IWORK ','DSURF ',NDSURF, JDSURF)
С
   DO 401 I=1,NDSURF
   IWORK(JDSURF+I-1)=IPT(I)
401 CONTINUE
  WRITE(NWRITE,*) ' '
  WRITE(NWRITE,*) '******** USRTRN - OUTLET *********
  WRITE(NWRITE,*) 'NDSURF=',NDSURF,' JDSURF=',JDSURF
  WRITE(NWRITE,*) ' IWORK(JDSURF)=',IWORK(JDSURF)
  WRITE(NWRITE,*) ' IWORK(JDSURF+NDSURF-1)=',IWORK(JDSURF+NDSURF-1)
С
   ENDIF
С
C-
               -----
   END OPERATION
С
С
  IF(KSTEP.EQ.NSTEP) THEN
С
```

С

C.....GAS FLOW RATE FROM SOURCE IS SET IN USRSRC

INFLOE\_SOURCE=GASIN

GASOUT = GASOUT\_2

C....INLET

CALL IPALL('NOZZLE INLET','\*','PATCH','CENTRES',IPT,NPT, + CWORK, IWORK) С GASIN=0.0 GASIN 2=0.0 WATIN = 0.0C.....GASIN INITIALISED IN USRSRC DO 501 I=1,NPT INODE = IPT(I)IBDRY = INODE-NCELL IFACE = IPFACB(IBDRY) WATIN = WATIN+CONV(IFACE,1) GASIN\_2 = GASIN\_2+CONV(IFACE,2) 501 CONTINUE GASIN=GASIN\_2 SUM=0.0 SUM=INFLOE\_SOURCE+GASIN WRITE(NWRITE,\*) ' ' WRITE(NWRITE,\*) '&&&&&& SHOWING OF OUTPUT OF INFLOW DATA &&&&& WRITE(NWRITE,\*) ' ' WRITE(NWRITE,\*) 'GAS INFLOW RATE FROM SOURCE=',INFLOE\_SOURCE WRITE(NWRITE,\*) 'GAS INFLOW RATE FROM INLET=',GASIN WRITE(NWRITE,\*) 'THE SUM OF GAS INFLOW RATE FROM SOURCE AND +INLET=',SUM WRITE(NWRITE,\*) ' ' WRITE(NWRITE,\*) ' ' С C....OUTLET GASOUT SINK=GASOUT WRITE(NWRITE,\*) '&&&&&& SHOWING OF OUTPUT OF OUTFLOW DATA &&&&& WRITE(NWRITE,\*) ' ' WRITE(NWRITE,\*) 'GAS OUTFLOW RATE FROM TOP SURFACE=', GASOUT\_SINK WRITE(NWRITE,\*) ' ' CALL IPALL('OUTLET','\*','PATCH','CENTRES',IPT,NPT, CWORK, IWORK) + С GASOUT = 0.0 $GASOUT_2 = 0.0$ WATOUT = 0.0C.....GASOUT INITIALISED IN USRSRC DO 601 I=1,NPT INODE = IPT(I)IBDRY = INODE-NCELL IFACE = IPFACB(IBDRY) WATOUT = WATOUT+CONV(IFACE,1) GASOUT\_2 = GASOUT\_2+CONV(IFACE,2) 601 CONTINUE

```
SUM1=0.0
   SUM1= GASOUT_SINK+GASOUT
   GFLOWRATE_IN=SUM/0.559
   GFLOWRATE_OUT=SUM1/0.559
   WFLOWRATE_IN=WATIN/7020
  WFLOWRATE_OUT=WATOUT/7020
  VOLFRAC = GFLOWRATE_IN/(GFLOWRATE_IN+WFLOWRATE_IN)
  WRITE(NWRITE,*) 'GAS OUTFLOW RATE FROM BOTTUM=', GASOUT
  WRITE(NWRITE,*) 'THE SUM OF GAS OUTFLOW RATE FROM TOP SURFACE AND
  + BOTTOM=',SUM1
  WRITE(NWRITE,*) ' '
  WRITE(NWRITE,*) ' '
С
  WRITE(NWRITE,*) ' *** MASS BALANCE ***'
  WRITE(NWRITE,*) ' '
  WRITE(NWRITE,*) ' AIR'
  WRITE(NWRITE,*) ' TOTAL INFLOW (m^3/s) =',GFLOWRATE_IN
  WRITE(NWRITE,*) ' TOTAL OUTFLOW (m^3/s) =',GFLOWRATE_OUT
  GASERR=(SUM1-SUM)*100.0/SUM
  WRITE(NWRITE,*) ' % ERROR
                               =',GASERR
  WRITE(NWRITE,*) ' '
  WRITE(NWRITE,*) 'WATER'
  WRITE(NWRITE,*) ' INFLOW (m^3/s) =',WFLOWRATE_IN
  WRITE(NWRITE,*) ' OUTFLOW (m^3/s) =',WFLOWRATE_OUT
   WATERR=(WATOUT+WATIN)*100.0/WATIN
  WRITE(NWRITE,*) ' % ERROR
                               =',WATERR
  WRITE(NWRITE,*) ' '
С
  WRITE(NWRITE,*) 'Volume Fraction of gas =',VOLFRAC
С
C.....SET IPLT=1 TO CLEAN UP VELOCITY FOR PLOTTING
C*****NOTE THAT THIS OPERATION WILL SLIGHTLY ALTER THE VELOCITY FIELD
C HENCE THE DUMP FILE GENERATED SHOULD BE USED FOR PLOTTING ONLY
С
  NOT FOR RESTART.
С
    IPLT=0
С
C.....FILTER OUT VELOCITY WITH SMALL VOLUME FRACTION
   IF(IPLT.EQ.1) THEN
     DO 701 INODE=1,NNODE
      FILTER=1.0-EXP(-3.0*VFRAC(INODE,2)/1.E-4)
     U(INODE,2)=U(INODE,2)*FILTER
      V(INODE,2)=V(INODE,2)*FILTER
     W(INODE,2)=W(INODE,2)*FILTER
      CONTINUE
701
С
   ENDIF
С
  ENDIF
С
С
  RETURN
  END
  SUBROUTINE USRBCS(VARBCS, VARAMB, A, B, C, ACND, BCND, CCND
  +
             ,IWGVEL,NDVWAL
```

,FLOUT,NLABEL,NSTART,NEND,NCST,NCEN + ,U,V,W,P,VFRAC,DEN,VIS,TE,ED,RS,T,H,RF,SCAL ,XP,YP,ZP,VOL,AREA,VPOR,ARPOR,WFACT,IPT + ,IBLK,IPVERT,IPNODN,IPFACN,IPNODF,IPNODB,IPFACB ,WORK,IWORK,CWORK) С C\* С C USER ROUTINE TO SET REALS AT BOUNDARIES. С C >>> IMPORTANT <<< C >>> <<< C >>> USERS MAY ONLY ADD OR ALTER PARTS OF THE SUBROUTINE WITHIN <<< C >>> THE DESIGNATED USER AREAS <<< С С C THIS SUBROUTINE IS CALLED BY THE FOLLOWING SUBROUTINE C CUSR SRLIST С C CREATED С 30/11/88 ADB C MODIFIED С 08/09/90 ADB RESTRUCTURED FOR USER-FRIENDLINESS. С 10/08/91 IRH FURTHER RESTRUCTURING ADD ACND BCND CCND С 22/09/91 IRH CHANGE ICALL TO IUCALL + ADD /SPARM/ С 10/03/92 PHA UPDATE CALLED BY COMMENT, ADD RF ARGUMENT, С CHANGE LAST DIMENSION OF RS TO 6 AND IVERS TO 2 С 03/06/92 PHA ADD PRECISION FLAG AND CHANGE IVERS TO 3 С 30/06/92 NSW INCLUDE FLAG FOR CALLING BY ITERATION С INSERT EXTRA COMMENTS С 03/08/92 NSW MODIFY DIMENSION STATEMENTS FOR VAX С 21/12/92 CSH INCREASE IVERS TO 4 С 02/08/93 NSW INCORRECT AND MISLEADING COMMENT REMOVED С 05/11/93 NSW INDICATE USE OF FLOUT IN MULTIPHASE FLOWS С 23/11/93 CSH EXPLICITLY DIMENSION IPVERT ETC. С 01/02/94 NSW SET VARIABLE POINTERS IN WALL EXAMPLE. С CHANGE FLOW3D TO CFDS-FLOW3D. С MODIFY MULTIPHASE MASS FLOW BOUNDARY TREATMENT. С 03/03/94 FHW CORRECTION OF SPELLING MISTAKE С 02/07/94 BAS SLIDING GRIDS - ADD NEW ARGUMENT IWGVEL С TO ALLOW VARIANTS OF TRANSIENT-GRID WALL BC С CHANGE VERSION NUMBER TO 5 С 09/08/94 NSW CORRECT SPELLING MOVE 'IF(IUSED.EQ.0) RETURN' OUT OF USER AREA С С 19/12/94 NSW CHANGE FOR CFX-F3D С 02/02/95 NSW CHANGE COMMON /IMFBMP/ С 02/06/97 NSW MAKE EXAMPLE MORE LOGICAL 02/07/97 NSW UPDATE FOR CFX-4 С 08/09/98 NSW CORRECT SIZE OF WALL ARRAY IN COMMENT С С \*\*\*\*\* C\* С C SUBROUTINE ARGUMENTS С С VARBCS - REAL BOUNDARY CONDITIONS С VARAMB - AMBIENT VALUE OF VARIABLES С Α - COEFFICIENT IN WALL BOUNDARY CONDITION

C B - COEFFICIENT IN WALL BOUNDARY CONDITION

- C ACND COEFFICIENT IN CONDUCTING WALL BOUNDARY CONDITION
- C BCND COEFFICIENT IN CONDUCTING WALL BOUNDARY CONDITION
- C CCND COEFFICIENT IN CONDUCTING WALL BOUNDARY CONDITION
- C IWGVEL USAGE OF INPUT VELOCITIES (0 = AS IS,1 = ADD GRID MOTION)
- C NDVWAL FIRST DIMENSION OF ARRAY IWGVEL
- C FLOUT MASS FLOW/FRACTIONAL MASS FLOW
- C NLABEL NUMBER OF DISTINCT OUTLETS
- C NSTART ARRAY POINTER
- C NEND ARRAY POINTER
- C NCST ARRAY POINTER
- C NCEN ARRAY POINTER
- C U U COMPONENT OF VELOCITY
- C V V COMPONENT OF VELOCITY
- C W W COMPONENT OF VELOCITY
- C P PRESSURE
- C VFRAC VOLUME FRACTION
- C DEN DENSITY OF FLUID
- C VIS VISCOSITY OF FLUID
- C TE TURBULENT KINETIC ENERGY
- C ED EPSILON
- C RS REYNOLD STRESSES
- C T TEMPERATURE
- C H ENTHALPY
- C RF REYNOLD FLUXES
- C SCAL SCALARS (THE FIRST 'NCONC' OF THESE ARE MASS FRACTIONS)
- C XP X COORDINATES OF CELL CENTRES
- C YP Y COORDINATES OF CELL CENTRES
- C ZP Z COORDINATES OF CELL CENTRES
- C VOL VOLUME OF CELLS
- C AREA AREA OF CELLS
- C VPOR POROUS VOLUME
- C ARPOR POROUS AREA
- C WFACT WEIGHT FACTORS
- С
- C IPT 1D POINTER ARRAY C IBLK - BLOCK SIZE INFORMA
- C IBLK BLOCK SIZE INFORMATION
- C IPVERT POINTER FROM CELL CENTERS TO 8 NEIGHBOURING VERTICES
- C IPNODN POINTER FROM CELL CENTERS TO 6 NEIGHBOURING CELLS
- C IPFACN POINTER FROM CELL CENTERS TO 6 NEIGHBOURING FACES
- C IPNODF POINTER FROM CELL FACES TO 2 NEIGHBOURING CELL CENTERS
- C IPNODB POINTER FROM BOUNDARY CENTERS TO CELL CENTERS
- C IPFACB POINTER TO NODES FROM BOUNDARY FACES
- C
- C WORK REAL WORKSPACE ARRAY
- C IWORK INTEGER WORKSPACE ARRAY
- C CWORK CHARACTER WORKSPACE ARRAY
- С
- C SUBROUTINE ARGUMENTS PRECEDED WITH A '\*' ARE ARGUMENTS THAT MUST C BE SET BY THE USER IN THIS ROUTINE.
- С
- C NOTE THAT OTHER DATA MAY BE OBTAINED FROM CFX-4 USING THE
- C ROUTINE GETADD, FOR FURTHER DETAILS SEE THE VERSION 4
- C USER MANUAL.
- LOGICAL LDEN,LVIS,LTURB,LTEMP,LBUOY,LSCAL,LCOMP
  - + ,LRECT,LCYN,LAXIS,LPOROS,LTRANS
- С

CHARACTER\*(\*) CWORK

```
С
C---- AREA FOR USERS EXPLICITLY DECLARED VARIABLES
С
С
  COMMON
  + /ALL/ NBLOCK,NCELL,NBDRY,NNODE,NFACE,NVERT,NDIM
  + /ALLWRK/ NRWS,NIWS,NCWS,IWRFRE,IWIFRE,IWCFRE
  + /ADDIMS/ NPHASE,NSCAL,NVAR,NPROP
       ,NDVAR,NDPROP,NDXNN,NDGEOM,NDCOEF,NILIST,NRLIST,NTOPOL
  + /BCSOUT/ IFLOUT
  + /CHKUSR/ IVERS,IUCALL,IUSED
  + /DEVICE/ NREAD,NWRITE,NRDISK,NWDISK
  + /IDUM/ ILEN,JLEN
  + /IMFBMP/ IMFBMP,JMFBMP
  + /LOGIC/ LDEN,LVIS,LTURB,LTEMP,LBUOY,LSCAL,LCOMP
       ,LRECT,LCYN,LAXIS,LPOROS,LTRANS
  +
  + /MLTGRD/ MLEVEL,NLEVEL,ILEVEL
  + /SGLDBL/ IFLGPR,ICHKPR
  + /SPARM/ SMALL,SORMAX,NITER,INDPRI,MAXIT,NODREF,NODMON
  + /TRANSI/ NSTEP,KSTEP,MF,INCORE
  + /TRANSR/ TIME, DT, DTINVF, TPARM
  + /UBCSFL/ IUBCSF
C
C---- AREA FOR USERS TO DECLARE THEIR OWN COMMON BLOCKS
C THESE SHOULD START WITH THE CHARACTERS 'UC' TO ENSURE
С
  NO CONFLICT WITH NON-USER COMMON BLOCKS
С
С
  DIMENSION
 + VARBCS(NVAR,NPHASE,NCELL+1:NNODE),VARAMB(NVAR,NPHASE)
 +,A(4+NSCAL,NPHASE,NSTART:*)
 +,B(4+NSCAL,NPHASE,NSTART:*),C(4+NSCAL,NPHASE,NSTART:*)
 +,FLOUT(*),ACND(NCST:*),BCND(NCST:*),CCND(NCST:*)
 +,IWGVEL(NDVWAL,NPHASE)
С
  DIMENSION
 + U(NNODE,NPHASE),V(NNODE,NPHASE),W(NNODE,NPHASE),P(NNODE,NPHASE)
 +,VFRAC(NNODE,NPHASE),DEN(NNODE,NPHASE),VIS(NNODE,NPHASE)
 +,TE(NNODE,NPHASE),ED(NNODE,NPHASE),RS(NNODE,NPHASE,6)
 +,T(NNODE,NPHASE),H(NNODE,NPHASE),RF(NNODE,NPHASE,4)
 +,SCAL(NNODE,NPHASE,NSCAL)
С
  DIMENSION
 + XP(NNODE), YP(NNODE), ZP(NNODE)
 +,VOL(NCELL),AREA(NFACE,3),VPOR(NCELL),ARPOR(NFACE,3),WFACT(NFACE)
  +,IPT(*),IBLK(5,NBLOCK)
 +,IPVERT(NCELL,8),IPNODN(NCELL,6),IPFACN(NCELL,6),IPNODF(NFACE,4)
 +,IPNODB(NBDRY,4),IPFACB(NBDRY)
  +,IWORK(*),WORK(*),CWORK(*),TMP(20,20,11)
С
C---- AREA FOR USERS TO DIMENSION THEIR ARRAYS
С
C---- AREA FOR USERS TO DEFINE DATA STATEMENTS
С
```

С C---- STATEMENT FUNCTION FOR ADDRESSING IP(I,J,K)=IPT((K-1)\*ILEN\*JLEN+(J-1)\*ILEN+I) С C----VERSION NUMBER OF USER ROUTINE AND PRECISION FLAG С IVERS=5 ICHKPR = 1 С C---- TO USE THIS USER ROUTINE FIRST SET IUSED=1 C AND SET IUBCSF FLAG: BOUNDARY CONDITIONS NOT CHANGING IUBCSF=0 С С BOUNDARY CONDITIONS CHANGING WITH ITERATION IUBCSF=1 С BOUNDARY CONDITIONS CHANGING WITH TIME IUBCSF=2 С BOUNDARY CONDITIONS CHANGING WITH TIME AND ITERATION IUBCSF=3 С IUBCSF=0 IUSED=1 С С IF (IUSED.EQ.0) RETURN С C---- FRONTEND CHECKING OF USER ROUTINE IF (IUCALL.EQ.0) RETURN С С C---- AREA FOR SETTING VALUES AT INLETS, PRESSURE BOUNDARIES AND OUTLETS. (NOTE THAT THE MASS FLOW AT OUTLETS IS С C SPECIFIED IN USER AREA 7) С OPEN(5,FILE='DATA1.DAT',STATUS='OLD',FORM='FORMATTED') READ(5,\*) KHEAD, JHEAD DO 13 K = 1, KHEAD DO 12 J = 1, JHEAD DO 11 I = 1, 10 READ(5,\*) TMP(K,J,I) CONTINUE 11 12 CONTINUE 13 CONTINUE CALL GETVAR('USRBCS','U ',IU) CALL GETVAR('USRBCS','V ',IV) CALL GETVAR('USRBCS','W ',IW) ',IVFRAC) CALL GETVAR('USRBCS','VFRAC CALL GETVAR('USRBCS','TE ',ITE) CALL GETVAR('USRBCS','ED ',IED) С C USE IPREC TO FIND ADDRESSES CALL IPREC('NOZZLE INLET', 'PATCH', 'CENTRES', IPT, ILEN, JLEN, KLEN, CWORK, IWORK) + С C LOOP OVER PATCH

DO 103 K = 1, KLEN

182

```
DO 102 J = 1, JLEN

DO 101 I = 1, ILEN

INODE = IP(I,J,K)

JSCALE = JHEAD*J/JLEN

KSCALE = KHEAD*K/KLEN

VARBCS(IU,1,INODE) = TMP(KSCALE,JSCALE,1)

VARBCS(IV,1,INODE) = TMP(KSCALE,JSCALE,2)

VARBCS(IV,1,INODE) = TMP(KSCALE,JSCALE,3)

VARBCS(ITE,1,INODE) = TMP(KSCALE,JSCALE,4)

VARBCS(IED,1,INODE) = TMP(KSCALE,JSCALE,5)

VARBCS(IU,2,INODE) = TMP(KSCALE,JSCALE,5)

VARBCS(IV,2,INODE) = TMP(KSCALE,JSCALE,6)

VARBCS(IV,2,INODE) = TMP(KSCALE,JSCALE,7)

VARBCS(IV,2,INODE) = TMP(KSCALE,JSCALE,8)

XFRAC=TMP(KSCALE,JSCALE,9)

IF(XFRAC.GE.1.0) THEN
```

```
VARBCS(IVFRAC,1,INODE) = 1.0
ENDIF
IF(XFRAC.LT.1.0) THEN
VARBCS(IVFRAC,1,INODE) = XFRAC
ENDIF
```

VARBCS(IVFRAC,2,INODE) = 1-VARBCS(IVFRAC,1,INODE)

101 CONTINUE 102 CONTINUE 103 CONTINUE C+++++++++++ END OF USER AREA 5 C RETURN END

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